### A Proof of Theorem 1

Here, we give a full proof of **Theorem 1**. Recall that an equivalent construction of the PAGA graphs is obtained by collapsing the graph  $G_1^{mt}$  to  $G_m^t$  leading to t super nodes. Accordingly, based on  $G_m^t$ , we can directly generate  $G_m^{t+1}$  by (temporarily) adding m more nodes, each linking with a single edge, and then collapsing these m nodes to the (t+1)th super node. Let  $d_{t,i}(v)$  denote the degree of the super-node v after mt+i-1 edges have been introduced, the probability that the mt+ith node (or equivalently its corresponding super-node t+1) connects to v is

$$\Pr\{(t+1,v)\} = \begin{cases} \frac{\mathsf{d}_{t,i}(v) + m(f(v,t+1)-1)}{(\overline{f}(\cdot,t+1)+1)(mt+i)-1} & 1 \le v \le t\\ \frac{\mathsf{d}_{t,i}(t+1) + i \cdot (\overline{f}(\cdot,t+1)-1) + 1}{(\overline{f}(\cdot,t+1)+1)(mt+i)-1} & v = t+1. \end{cases}$$
(15)

Given these definitions, we begin with proving statement 1. With Eq. (9) and starting conditions

$$N(k,1) = [k = 2m], \quad N(k,t) = N(k,t,1)$$
 (16)

we get a complete description of N(k,t) as Eq. (8) of statement 1

$$N(k, t+1) = N(k, t, m+1) + \Pr(d_{t,m+1}(t+1) = k)$$

Next we prove statement 2 which is straightforward from the definition. First, note that  $k \in \{m, m+1, ..., 2m\}$ , i.e.:

$$\Pr(\mathsf{d}_{t,k+1}(t+1)k) = 0, \quad k < m \text{ or } k > 2m.$$
 (17)

The minimal value  $d_{t,m+1}(t+1) = m$  is obtained when no one of the k edges is a loop. In this case,  $d_{t,i}(t+1) = i-1$  for all i, so:

$$\Pr(\mathsf{d}_{t,m+1}(t+1) = m) = \prod_{i=1}^{m} \left( 1 - \frac{i \cdot \overline{f}(\cdot, t+1)}{(\overline{f}(\cdot, t+1) + 1)(mt+i) - 1} \right)$$
$$= 1 + O_m \left( \frac{1}{t\overline{f}(\cdot, t+1)} \right).$$

From  $\sum_{k=m}^{2m} \Pr(d_{t,m+1} = k) = 1$  and  $\Pr(d_{t,m+1} = k) \ge 0$ :

$$\Pr(\mathsf{d}_{t,m+1}(t+1) = k) = O_m\left(\frac{1}{t\overline{f}(\cdot, t+1)}\right), m < d \le 2m.$$
(18)

Finally the claim follows from the definition of [k = m].

Before proving statement 3, we first derive some properties of c(k, t, i) and prove statement 4 en-route. Here, for brevity, we denote c(k, t, i) as c(k). Starting from c(m) we study the step-wise change of c(k) as:

$$c(m) = \frac{B(m \cdot \overline{f}(m,t,i), \overline{f}(\cdot,t) + 2)}{B(m \cdot \overline{f}(m,t,i), \overline{f}(\cdot,t) + 1)}$$

$$= \frac{\Gamma(\overline{f}(\cdot,t) + 2)}{\Gamma(\overline{f}(\cdot,t) + 1)} \times \frac{\Gamma(m \cdot \overline{f}(m,t,i) + \overline{f}(\cdot,t) + 1)}{\Gamma(m \cdot \overline{f}(m,t,i) + \overline{f}(\cdot,t) + 2)}$$

$$= \frac{\overline{f}(\cdot,t) + 1}{m \cdot \overline{f}(m,t,i) + \overline{f}(\cdot,t) + 1}$$
(19)

For k > m, the ratio of c(k, i) to c(k - 1, i) can be simplified as:

$$\frac{c(k-1)}{c(k)} = \frac{B\left(k-1+m(\overline{f}(k-1,t,i)-1)\right), \overline{f}(\cdot,t)+2\right)}{B\left(k+m(\overline{f}(k,t,i)-1)\right), \overline{f}(\cdot,t)+2)} 
= \frac{\Gamma\left(k-1+m(\overline{f}(k-1,t,i)-1)\right)}{\Gamma\left(k+m(\overline{f}(k,t,i)-1)\right)} \times \frac{\Gamma\left(k+m(\overline{f}(k,t,i)-1)\right)+\overline{f}(\cdot,t)+2\right)}{\Gamma\left(k-1+m(\overline{f}(k-1,t,i)-1)+\overline{f}(\cdot,t)+2\right)} 
= \frac{k+m(\overline{f}(k,t)-1)+1+\overline{f}(\cdot,t)}{k-1+m(\overline{f}(k-1,t)-1)},$$
(20)

where we make the approximation  $\overline{f}(k,t,i) \simeq \overline{f}(k-1,t,i)$  with the assumption that the average similarity of nodes with degree k to node t is similar to that of nodes with degree k+1. In particular, c(k-1) > c(k), so c(k) < c(m) < 1 for all  $k \ge m$ .

Now to analyze the asymptotic behavior of c(k) we begin by taking logarithm as:

$$\ln c(k,i) = C_0 + \ln \Gamma(k+C_1) - \ln \Gamma(k+C_1 + \overline{f}(\cdot,t) + 2)$$

$$C_0 = \ln \frac{\Gamma(\overline{f}(\cdot,t) + 2)}{B(m \cdot \overline{f}(m,t,i), \overline{f}(\cdot,t) + 1)}, C_1 = m(\overline{f}(k,t,i) - 1)$$
(21)

Exploiting the fact that asymptotically

 $\Gamma(x+1) \sim \sqrt{2\pi x} \left(\frac{x}{e}\right)^x \xrightarrow{\ln} \ln \sqrt{2\pi} + x(\ln x - 1) + \frac{1}{2} \ln x$ , Eq. (21) becomes:

$$C_{0} + (k + C_{1}) \left( \ln(k + C_{1}) - 1 \right) + O\left(\frac{1}{k}\right) - \left(k + C_{1} + \overline{f}(\cdot, t) + 2\right) \left( \ln(k + C_{1} + \overline{f}(\cdot, t) + 2) - 1 \right)$$

$$= C_{0} + (k + C_{1}) \left( \ln k + \frac{C_{1}}{k} - 1 \right) + O\left(\frac{1}{k}\right) - \left(k + C_{1} + \overline{f}(\cdot, t) + 2\right) \left( \ln k + \frac{C_{1} + \overline{f}(\cdot, t) + 2}{k} - 1 \right)$$

$$= C_{0} - \left(2 + \overline{f}(\cdot, t)\right) \ln k + O\left(\frac{1}{k}\right)$$

Rewriting the  $C_0$  term, the full expression of asymptotic value of c(k) as k grows becomes

$$c(k) = \frac{\Gamma(\overline{f}(\cdot,t)+2)}{B(m(\overline{f}(m,t,i)),\overline{f}(\cdot,t)+1)} k^{-2-\overline{f}(\cdot,t)} \bigg(1 + O\bigg(\frac{1}{k}\bigg)\bigg).$$

To prove the statement 3, we show

$$\begin{split} &N(k,t,i+1) \\ &= \sum_{v=1}^{t} \Pr(\mathsf{d}_{t,i+1}(v) = k) \\ &= \sum_{v=1}^{t} (\Pr(\mathsf{d}_{t,i+1}(v) = k, \gamma = v) + \Pr(\mathsf{d}_{t,i+1}(v) = k, \gamma \neq v)) \\ &= \sum_{v=1}^{t} (\Pr(\mathsf{d}_{t,i}(v) = k - 1, \gamma = v) + \Pr(\mathsf{d}_{t,i}(v) = k, \gamma \neq v)) \\ &= \sum_{v=1}^{t} \left( \Pr(\mathsf{d}_{t,i}(v) = k - 1) \frac{k - 1 + m(f(v,t+1) - 1)}{(\overline{f}(\cdot,t+1) + 1)(mt+i) - 1} + \Pr(\mathsf{d}_{t,i}(v) = k) \left( 1 - \frac{k + m(f(v,t+1) - 1)}{(\overline{f}(\cdot,t+1) + 1)(mt+i) - 1} \right) \right) \\ &= N(k - 1, t, i) \frac{k - 1 + m(\overline{f}(k - 1, t + 1, i) - 1)}{(\overline{f}(\cdot,t+1) + 1) \cdot (mt+i) - 1} + N(k, t, i) \left( 1 - \frac{k + m(\overline{f}(k, t + 1, i) - 1)}{(\overline{f}(\cdot,t+1) + 1) \cdot (mt+i) - 1} \right) \end{split}$$

By proving statement 4, we complete our proof of **Theorem 1**. For the ease of proof, let us define N(k,t,i):

$$\tilde{N}(k,t,i) = N(k,t,i) - c(k,t,i) \left( t + \frac{i}{m} - \frac{1}{m(1+\overline{f}(\cdot,t))} \right)$$
(23)

The lhs of (11) can be expressed in a recursive fashion as:

$$\begin{split} \tilde{N}(k,t,i+1) + \frac{i}{m}[k=m] \\ &= \tilde{N}(k,t,i+1) + \frac{1}{m}[k=m] + \frac{i-1}{m}[k=m] \\ &= \tilde{N}(k,t,i+1) + [k=m]c(m,t,i) \frac{m(\overline{f}(m,t,i)) + \overline{f}(\cdot,t) + 1}{m(1+\overline{f}(\cdot,t))} + \frac{i-1}{m}[k=m] \end{split}$$

$$\begin{split} &= N(k,t,i+1) - c(k,t,i) \frac{(\overline{f}(\cdot,t)+1)(mt+i+1)-1}{m(1+\overline{f}(\cdot,t))} - c(m,t,i) \frac{[k=m](m(\overline{f}(m,t,i))+\overline{f}(\cdot,t)+1)}{m(1+\overline{f}(\cdot,t))} + \frac{i-1}{m} [k=m] \\ &= N(k,t,i+1) - \left( \frac{(\overline{f}(\cdot,t)+1)(mt+i)-1}{m(\overline{f}(\cdot,t)+1)} \right) \\ &\quad \times c(k,t,i) \frac{(\overline{f}(\cdot,t)+1)(mt+i+1)-1-[k=m]\left(k+m(\overline{f}(k,t)-1)+\overline{f}(\cdot,t)+1\right)}{(1+\overline{f}(\cdot,t))\cdot(mt+i)-1} + \frac{i-1}{m} [k=m] \\ &= N(k,t,i+1) + \frac{i-1}{m} [k=m] - \left(t+\frac{i}{m} - \frac{1}{m(1+\overline{f}(\cdot,t))}\right) \frac{c(k,i)}{(1+\overline{f}(\cdot,t))(mt+i)-1} \\ &\quad \times \left\{ (1-[k=m])\left(k+m(\overline{f}(k,t,i)-1)+1+\overline{f}(\cdot,t)\right) + \left((\overline{f}(\cdot,t)+1)(mt+i)-1\right) - \left(k+m(\overline{f}(k,t,i)-1)\right) \right\} \\ &= N(k,t,i+1) + \frac{i-1}{m} [k=m] - \left(t+\frac{i}{m} - \frac{1}{m(1+\overline{f}(\cdot,t))}\right) \\ &\quad \times \left\{ c(k-1,t,i) \frac{k-1+m(\overline{f}(k-1,t,i)-1)}{(1+\overline{f}(\cdot,t))\cdot(mt+i)-1} + c(k,i) \left(1-\frac{k+m(\overline{f}(k,t,i)-1)}{(1+\overline{f}(\cdot,t))(mt+i)-1}\right) \right\} \\ &= \tilde{N}(k-1,t,i) \frac{k-1+m(\overline{f}(k-1,t,i)-1)}{(1+\overline{f}(\cdot,t))\cdot(mt+i)-1} + \tilde{N}(k,t,i) \left(1-\frac{k+m(\overline{f}(k,t,i)-1)}{(1+\overline{f}(\cdot,t))(mt+i)-1}\right) + \frac{i-1}{m} [k=m] \end{split}$$

Note that the last line comes from Eq. (22) and Eq. (23). Now, in a fashion similar to [18] involving manipulations like the ones done above, we have  $\tilde{N}(k,t,i)+[k=m]\frac{i-1}{m}=O_m\left(\frac{1}{k}\right)$  finishing the proof of the theorem.

#### B Proof of Theorem 3

The proof of the bound on the diameter for PAGA graphs can by obtained by extending the original proof for the standard preferential attachment model from a uniform measure to a non-uniform measure, described by the affinity function f. Here, we ignore the case of having disconnected components in a graph as the graph will be one connected component with high probability.<sup>2</sup> As self loops do not affect the eccentricity of a node and hence the diameter, we simply ignore them in the generation process itself – for the purpose of bounding the diameter.

We begin the proof by noting that for m = 1 the graph is a tree and for the cases  $m \ge 2$ , which are formed by collapsing the graph  $G_1^{mt}$  graph, the diameter can only shrink. Next, the diameter of the graph can not be larger than twice the height of this tree, which is equal to the maximal graph distance between vertex 1 and any of the other vertices. So, it is sufficient to bound the height of the tree.

For bounding the tree height, we follow the steps of [48] and outline here the differences.<sup>3</sup> We start with a continuous time branching process, where the rate is given by  $\lambda^t(j) = d(j) + f(j,t)$ . Therefore, the overall transition rate after t vertices are present (i.e. after t-1 births) is given by:

$$\sum_{i=1}^{t} (\mathsf{d}(j) + f(t,j)) = 2t + t\bar{f}(\cdot, t+1)$$
 (24)

Now we can decompose the time  $\tau_t$  as a sum of independent variables, exponentially distributed with parameter  $2t + t\bar{f}(\cdot, t+1)$ , i.e.

$$\tau_t = \sum_{i=1}^t \mathfrak{t}_j$$

where  $t_j \sim \text{Exp}(2j + j\bar{f}(\cdot, j+1))$ . It follows that the mean and variance are bounded by

$$\mathbf{E}[\tau_t] = \sum_{j=1}^{t} \frac{1}{2j + j\bar{f}(\cdot, j+1)} = O(\log t)$$

<sup>&</sup>lt;sup>2</sup>In order for a PAGA graph to have a disconnected component, a new super-node t+1 has to make m self loops. This probability asymptotically goes to 0. Note that even in the uniform case this probability is  $\frac{1}{t}^m$ , while the self-loop probability in PAGA is typically smaller than 1/t.

<sup>&</sup>lt;sup>3</sup>To avoid clash of notations, we also redefine the birth times as  $\tau_1,...,\tau_t$  instead of  $t_i$ .

$$var(\tau_t) = \sum_{i=1}^{t} \frac{1}{2j + j\bar{f}(\cdot, j+1)} = O(1)$$

These conditions match the ones required in [48]. Accordingly, the small world property holds for the PAGA model.

# C Proof of Theorem 4

Note that at every timestep, there can be at most  $\frac{m(m-1)}{2}$  new triangles added. So the number of triangles  $T(G_m^t)$  is bounded by O(n). On the other hand, the number of triplets in  $G_m^t$  follows  $\sum_{k=1}^{d_{max}} N(k,t) {k \choose 2} \propto \sum_{k=1}^{d_{max}} t \cdot k^{\bar{f}(\cdot,t)}$ . As the sum  $\sum_{k=1}^{d_{max}} k = mt$ , it is straightforward that  $C(G_m^t) \to 0$  unless  $\bar{f}(\cdot,t) > 1$ . We omit the analysis for the cases where  $\bar{f}(\cdot,t) > 1$  as real graphs follow power law exponent around -2.1 to -2.5  $(0.1 \le \bar{f}(\cdot,t) \le 0.5)$ 

## D Spyplot of W

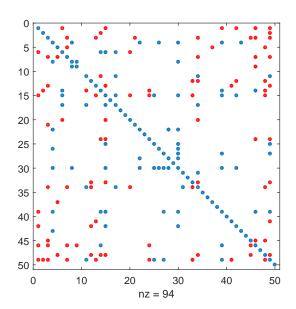


Figure 6: Spyplot of the matrix W. Blue entries correspond to positive elements; red entries to negative elements

# E Topics and keywords learned by LDA

We used 50 topics (i.e. W is a  $50 \times 50$  matrix) for learning the affinity function f. Here, we represent the top 20 keywords for top 10 topics we obtained from LDA.

Topic 0	Topic 1	Topic 2	Topic 3	Topic 4
method	quantization	gauge	theory	gauge
integral	hamiltonian	invariant	field	magnetic
path	dirac	invariance	approach	nonabelian
using	zero	local	based	abelian
functional	modes	brst	discuss	chernsimons
regularization	canonical	lagrangian	nonperturbative	electric
approach	physical	covariant	presented	charge
formula	lightcone	constraints	properties	theory
use	coordinates	formulation	new	monopole
integrals	shown	formalism	methods	su
expansion	quantized	class	may	monopoles
used	formalism	first	present	flux
series	variables	ghost	talk	charged
procedure	mode	cohomology	version	higgs
technique	constraint	extended	interpretation	dual
applied	leads	fields	developed	term
obtained	operator	lorentz	techniques	vortex
integration	approach	fixing	provide	selfdual
expression	constraints	transformations	used	vortices
formalism	formulation	auxiliary	analysis	yangmills
Topic 5	Topic 6	Topic 7	Topic 8	Topic 9
n	matrix	coupling	loop	form
n supersymmetric		coupling point	loop theory	
n supersymmetric supersymmetry	matrix	coupling	loop theory perturbation	form general parameters
n supersymmetric supersymmetry supergravity	matrix model matrices ansatz	coupling point limit fixed	loop theory	form general parameters arbitrary
n supersymmetric supersymmetry	matrix model matrices ansatz spectral	coupling point limit	loop theory perturbation	form general parameters arbitrary given
n supersymmetric supersymmetry supergravity	matrix model matrices ansatz	coupling point limit fixed	loop theory perturbation renormalization wilson finite	form general parameters arbitrary
n supersymmetric supersymmetry supergravity superconformal super superspace	matrix model matrices ansatz spectral models relation	coupling point limit fixed large	loop theory perturbation renormalization wilson	form general parameters arbitrary given
n supersymmetric supersymmetry supergravity superconformal super superspace yangmills	matrix model matrices ansatz spectral models relation limit	coupling point limit fixed large strong	loop theory perturbation renormalization wilson finite	form general parameters arbitrary given explicit parameter q
n supersymmetric supersymmetry supergravity superconformal super superspace yangmills multiplet	matrix model matrices ansatz spectral models relation limit integrable	coupling point limit fixed large strong flow	loop theory perturbation renormalization wilson finite perturbative expansion one	form general parameters arbitrary given explicit parameter
n supersymmetric supersymmetry supergravity superconformal super superspace yangmills multiplet superfield	matrix model matrices ansatz spectral models relation limit	coupling point limit fixed large strong flow points constant group	loop theory perturbation renormalization wilson finite perturbative expansion	form general parameters arbitrary given explicit parameter q
n supersymmetric supersymmetry supergravity superconformal super superspace yangmills multiplet	matrix model matrices ansatz spectral models relation limit integrable	coupling point limit fixed large strong flow points constant	loop theory perturbation renormalization wilson finite perturbative expansion one	form general parameters arbitrary given explicit parameter q values
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n supersymmetric supersymmetry supergravity superconformal super superspace yangmills multiplet superfield multiplets cal bps theories harmonic susy	matrix model matrices ansatz spectral models relation limit integrable spin smatrix scattering elements eigenvalues chain bethe	coupling point limit fixed large strong flow points constant group renormalization weak theory infrared constants rg	loop theory perturbation renormalization wilson finite perturbative expansion one order loops non result lambda cutoff diagrams	form general parameters arbitrary given explicit parameter q values set independent can terms one
n supersymmetric supersymmetry supergravity superconformal super superspace yangmills multiplet superfield multiplets cal bps theories harmonic	matrix model matrices ansatz spectral models relation limit integrable spin smatrix scattering elements eigenvalues chain	coupling point limit fixed large strong flow points constant group renormalization weak theory infrared constants	loop theory perturbation renormalization wilson finite perturbative expansion one order loops non result lambda cutoff diagrams divergences	form general parameters arbitrary given explicit parameter q values set independent can terms one structure
n supersymmetric supersymmetry supergravity superconformal super superspace yangmills multiplet superfield multiplets cal bps theories harmonic susy	matrix model matrices ansatz spectral models relation limit integrable spin smatrix scattering elements eigenvalues chain bethe	coupling point limit fixed large strong flow points constant group renormalization weak theory infrared constants rg	loop theory perturbation renormalization wilson finite perturbative expansion one order loops non result lambda cutoff diagrams divergences expectation	form general parameters arbitrary given explicit parameter q values set independent can terms one structure forms present real
n supersymmetric supersymmetry supergravity superconformal super superspace yangmills multiplet superfield multiplets cal bps theories harmonic susy superfields	matrix model matrices ansatz spectral models relation limit integrable spin smatrix scattering elements eigenvalues chain bethe corresponding	coupling point limit fixed large strong flow points constant group renormalization weak theory infrared constants rg behavior	loop theory perturbation renormalization wilson finite perturbative expansion one order loops non result lambda cutoff diagrams divergences	form general parameters arbitrary given explicit parameter q values set independent can terms one structure forms present

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