



Spin excitations and spin dynamics in nanoscopic systems

Antonio T. Costa

National Nanotechnology Laboratory¹
National Center for Research in Energy and Materials
Campinas - SP - Brazil

FAPESP, 23/06/2017

¹On sabbatical leave from Instituto de Física, Universidade Federal Fluminense  

Collaborators

Theory

- Adalberto Fazzio (LNNano)
- Marcio Costa (LNNano)
- Mauro Ferreira (TCD)
- Roberto Bechara Muniz (UFF)
- Filipe Guimarães (FZ Jülich)
- Samir Lounis (FZ Jülich)

Experiments

- Harald Ibach (FZ Jülich)
- Alex Khajetoorians (Radboud University, Nijmegen)
- Jens Wiebe (Hamburg University)

Why spin dynamics (excitations)?

“Static” (ground state) properties provide only a small fraction of the available information

Dynamic properties are related to how the system

- behaves at different temperatures,
- relax towards equilibrium,
- absorbs and emits energy,
- interacts with external probes,
- converts charge currents into spin currents and vice-versa,
- etc.

Experimental techniques

- Ferromagnetic resonance (FMR) – bulk; $\lambda \rightarrow \infty$.
- Neutron scattering – bulk; resolves λ .
- Electron energy loss spectroscopy (EELS) – surface sensitive; resolves λ .
- Inelastic scanning tunneling spectroscopy – localized excitations.
- Spin-Hall effects (direct and inverse) – transport experiments.

Typical Systems – metallic substrates

Ultrathin films



Trilayers

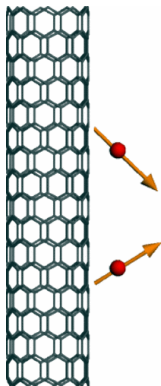


Adatoms

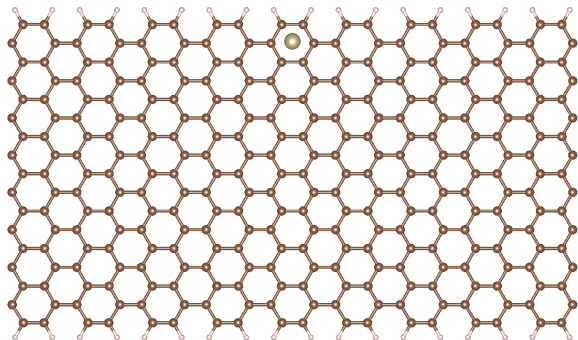


Typical Systems – low-dimensional substrates

Adatoms on carbon nanotubes



Adatoms on graphene nanoribbons and topological insulators



- LCAO description of the electronic structure.

$$H = H_0 + H_U + H_{\text{SO}}$$

$$H_0 = \sum_{l,l';\mu,\nu;\sigma} T_{ll'}^{\mu\nu} a_{l\mu\sigma}^\dagger a_{l'\nu\sigma}$$

$$H_U = \sum_l \sum_{\mu,\nu,\mu'\nu'} \sum_{\sigma,\sigma'} U_l^{\mu\nu\mu'\nu'} a_{l\mu\sigma}^\dagger a_{l\nu\sigma'}^\dagger a_{l\nu'\sigma'} a_{l\mu'\sigma}$$

$$H_{\text{SO}} = \sum_l \sum_{\mu\nu} \frac{\xi_l}{2} \left[L_{\mu\nu}^z (c_{l\mu\uparrow}^\dagger c_{l\nu\uparrow} - c_{l\mu\downarrow}^\dagger c_{l\nu\downarrow}) + L_{\mu\nu}^+ c_{l\mu\downarrow}^\dagger c_{l\nu\uparrow} + L_{\mu\nu}^- c_{l\mu\uparrow}^\dagger c_{l\nu\downarrow} \right]$$

- $T_{ll'}^{\mu\nu}$ obtained from DFT calculations [PRB **88**, 165127 (2013)].

What I talk about when I talk about spin dynamics

- Spin excitation spectra from the electronic structure alone:

$$\chi_{ll'}^{+-}(t) = -i\theta(t) \left\langle \left[S_l^+(t), S_{l'}^-(0) \right] \right\rangle$$

$$S_l^- \equiv a_{l\downarrow}^\dagger a_{l\uparrow}$$

- All damping mechanisms (except for magnon-magnon scattering) are inherent to this approach:
 - ▶ spin-orbit coupling damping,
 - ▶ spin pumping damping,
 - ▶ damping via Stoner excitations.
- SOC \Rightarrow Magnetocrystalline anisotropy \Rightarrow gap in the spin wave spectrum.

Spin dynamics in the presence of SOC

SOC couples transverse spin excitations, given by

$$\chi_{ll'}^{+-}(t) = -i\theta(t) \left\langle \left[S_l^+(t), S_{l'}^-(0) \right] \right\rangle$$

to longitudinal spin excitations and charge excitations,

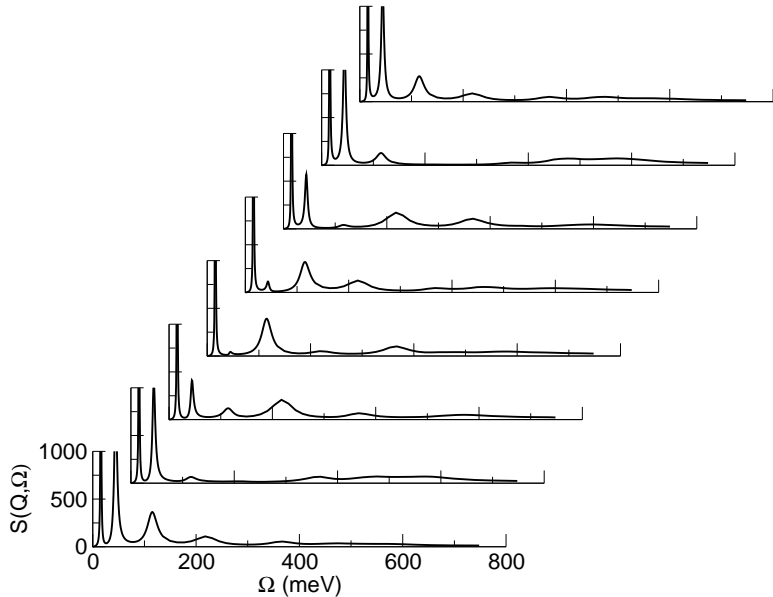
$$\chi_{ll'}^{\uparrow-}(t) = -i\theta(t) \left\langle \left[n_l^{\uparrow}(t), S_{l'}^-(0) \right] \right\rangle$$

$$\chi_{ll'}^{\downarrow-}(t) = -i\theta(t) \left\langle \left[n_l^{\downarrow}(t), S_{l'}^-(0) \right] \right\rangle$$

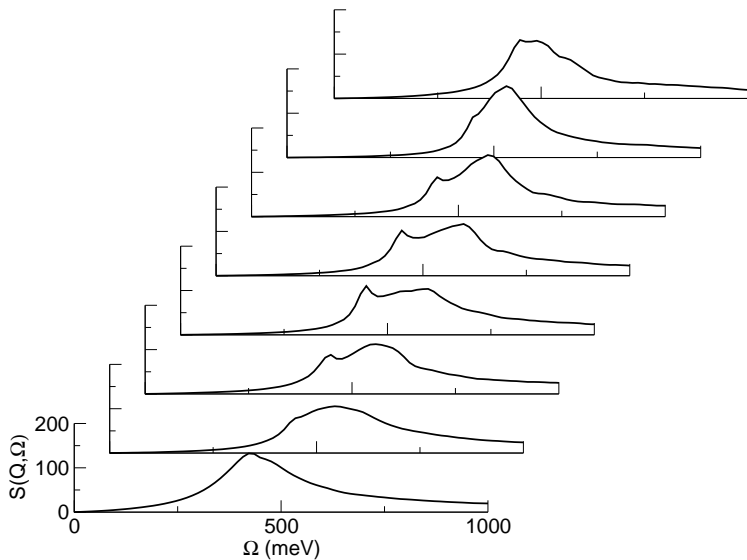
$$\chi_{ll'}^{\bar{-}}(t) = -i\theta(t) \left\langle \left[S_l^-(t), S_{l'}^-(0) \right] \right\rangle$$

PRB **82**, 014428 (2010)

SW spectrum: 8Co/Cu(001), long wavelength

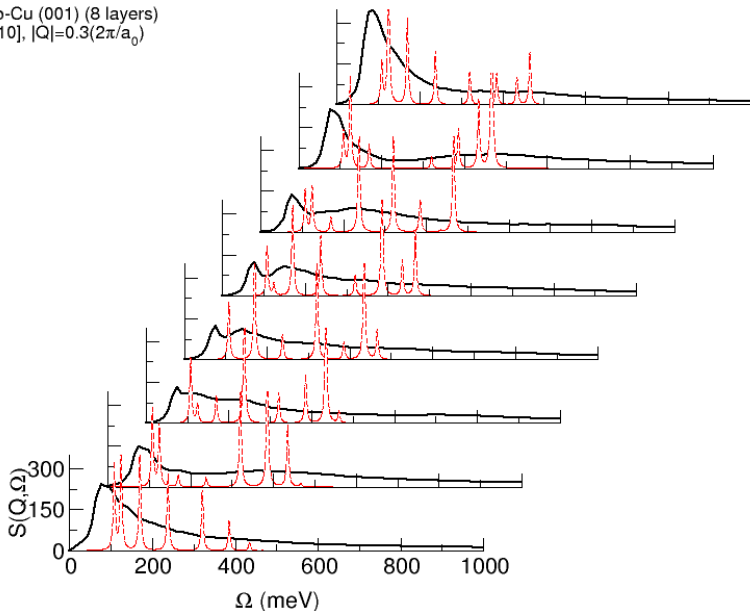


SW spectrum: 8Co/Cu(001), short wavelength



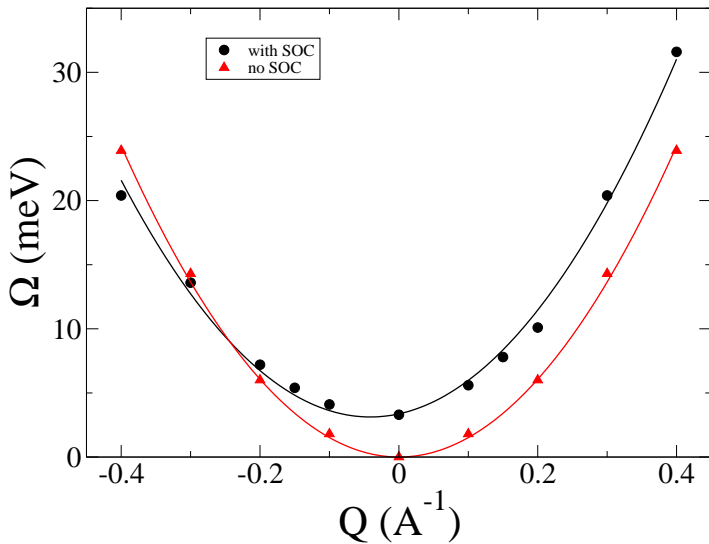
SW spectrum: 8Co/Cu(001), intermediate wavelength

Co-Cu (001) (8 layers)
[110], $|Q|=0.3(2\pi/a_0)$



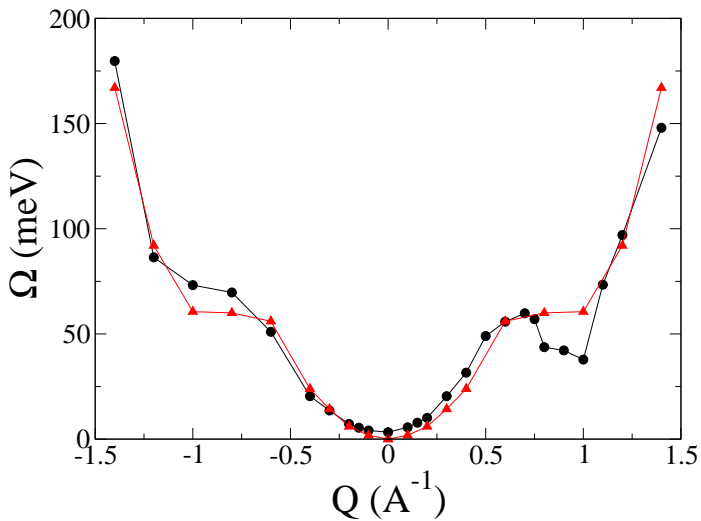
SW dispersion relation - Fe/W(110)

Magnetocrystalline anisotropy gap.

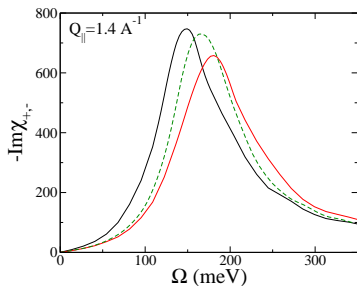
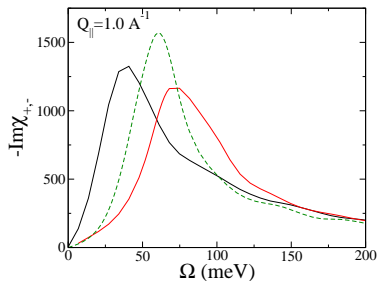
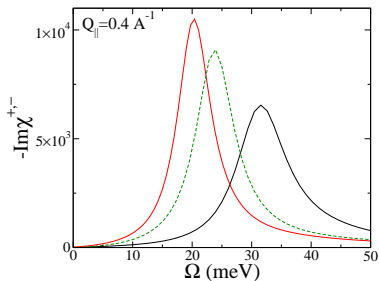


SW dispersion relation - Fe/W(110)

Dzyaloshiskii-Moriya coupling - left-right asymmetry.

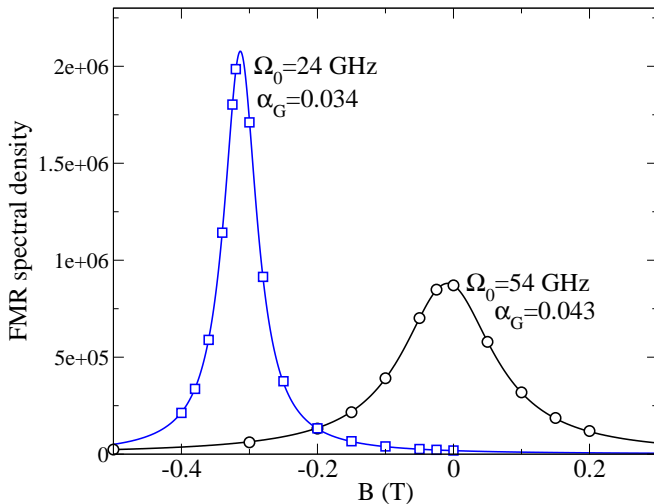


Magnon spectral density



Frequency-dependent Gilbert damping - 2Co/2Pt(001)

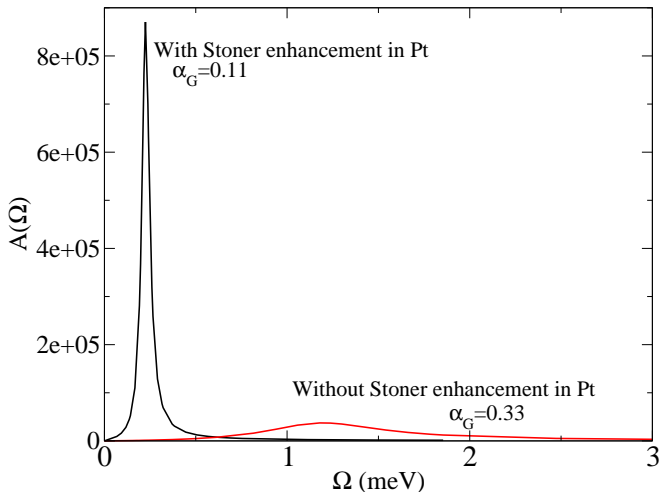
FMR sweeping Zeeman field.



PRB **92**, 014419 (2015).

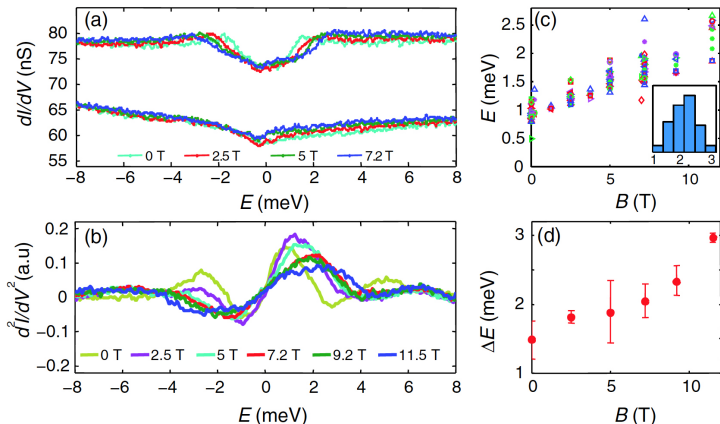


Gilbert damping and Stoner enhancement - 2Co/2Pt(001)



PRB **92**, 014419 (2015).

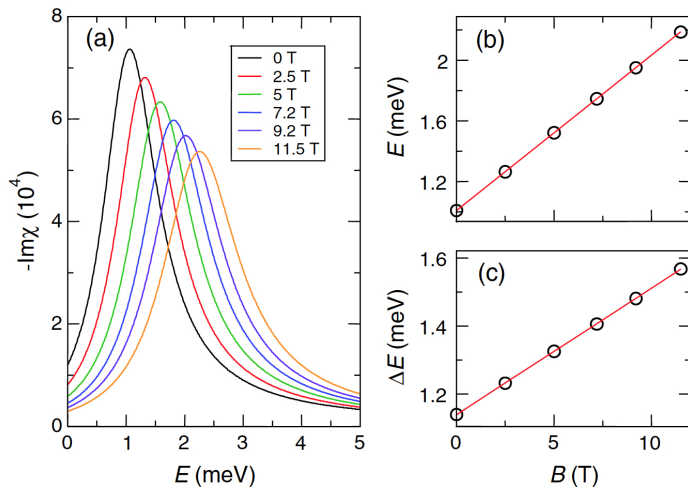
Fe adatom on Cu(111) surface



$$\bar{g} \approx 2$$

PRL **106**, 037205 (2011).

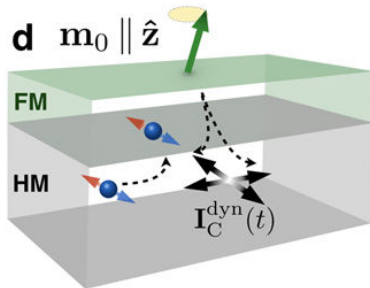
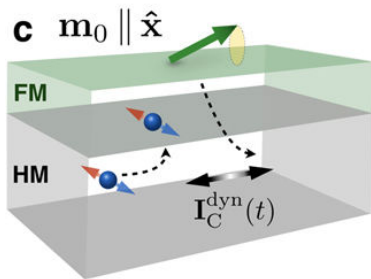
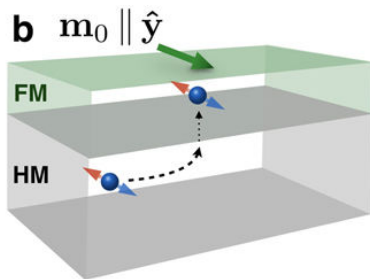
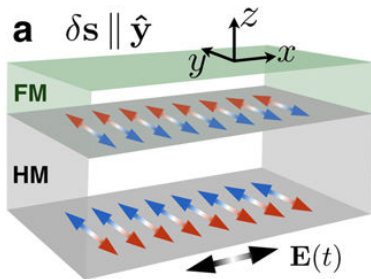
Fe adatom on Cu(111) surface - Calculations



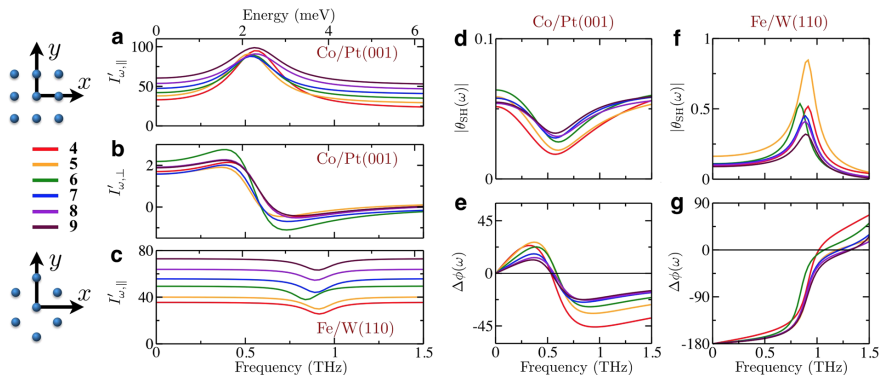
$g = 1.8$; magnetic anisotropy ~ 1 meV.

PRL **106**, 037205 (2011).

Dynamical (THz) Spin Hall Effects



Dynamical (THz) Spin Hall Effects



Scientific Reports **7**, 3686 (2017).

Ongoing project: adatoms on topological insulators

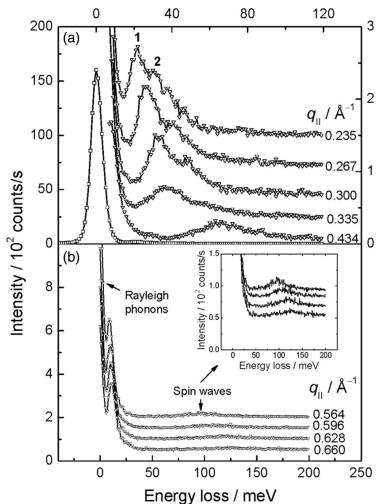
Do spin excitations couple to “topological states”?

- How does the coupling affects the spin excitations?
- How does the coupling affect the “topological character” of the edge states?
- Is it possible to modulate the edge spin current via the coupling?
- Do adatoms couple to each other through the edge states?
- ...

Concluding remarks

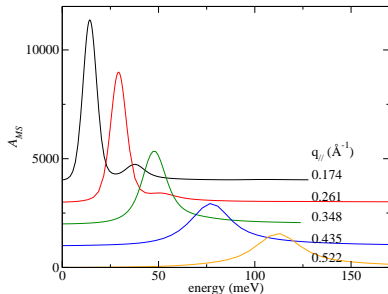
- We have a method for calculating spin excitation spectra that requires nothing more than tight-binding parameters provided by a DFT-based calculation.
- All damping mechanisms are intrinsic to our approach, except those coming from magnon-magnon scattering.
- We can address results from many different experimental techniques with essentially the same formalism.
- Dynamical Spin Hall Effects are the latest addition to our capabilities.

Spin wave spectra of ultrathin films - 8Co/Cu(001)



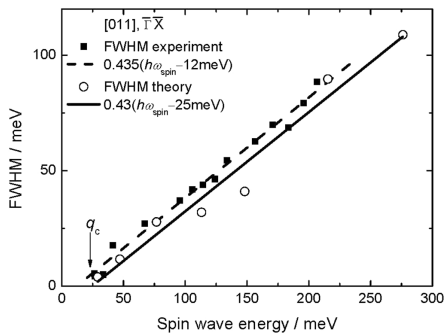
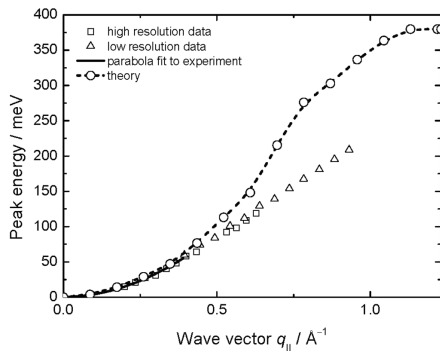
First experimental observation of optical spin wave modes on ultrathin metallic films.

Theory



PRB 86, 165436 (2012)

Spin waves on 8Co/Cu(001) - dispersion and linewidths



PRB **86**, 165436 (2012)