## CL-Chemy Transforms Fibonacci-Type Sequences to Arrays

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#### Abstract

The *CL*, or *coefficient list* concept is a natural adjunct to many recursive formulas and algorithms. The basic idea is that coefficients may vary as iteration proceeds. This article explores the effect of CLs on a fascinating family of sequences; namely the close relatives of the eponymous numbers brought to our attention by Leonardo da Pisa (aka *Fibonacci*) circa 1200 CE.

Terms of the Fibonacci sequence will derive, by Binet's formula, from powers of the zeros of  $x^2 - x - 1$ . By a mathematical rule-of-thumb, patterns spawned by  $2^{nd}$ -degree equations have a 2-dimensional representation. The exposition that follows provides a context in which the familiar Fibonacci sequence is properly thought of as a single row/column of a (degenerate) 2D array.

# CL-Chemy Transforms Fibonacci-Type Sequences to Arrays

Openings to a New Dimension in the Exploration of the Remarkable  $\varphi$  Sequence

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A Derivation of the Limit Ratios of a Generalized Fibonacci Sequence We begin with a generalized Fibonacci formula

$$(c)F_n + (b)F_{n+1} = F_{n+2}$$
  $F_0 = 0, F_1 = 1$  (1.1)

Henceforth (1.1) will be called a  $\varphi$ -sequence or  $\varphi_1$ .

*Theorem*: For *b*,*c* in  $\mathbb{N}$  ( $\mathbb{N} = 1,2,3...\infty$ ), the ratio of adjacent terms in  $\varphi_1$  converges to a real number, *x*. I.e.,

$$\lim_{n \to \infty} \frac{F_{n+1}}{F_n} = x$$

*Proof*: In the limit:  $x^2 = \frac{F_{n+2}}{F_n} = \frac{(c)F_n + (b)F_{n+1}}{F_n} = c + \frac{(b)F_{n+1}}{F_n} = bx + c$ , and  $x^2 - bx - c = 0$ .

Since the discriminant  $b^2 - 4ac$  is always positive (see the quadratic formula in (1.2)), x is real.

QED

E.g., if b = 2 and c = 3, then  $\varphi_1 = 0, 1, 2, 7, 20, 61, 182, 547, 1640, 4921, 14762...$ 

To what ratio does this sequence converge? Let  $x = r_{\pm}$ . Then by the *quadratic formula* 

$$r_{\pm} = \frac{-b \pm \sqrt{b^2 - 4ac}}{2a}$$
(1.2)

the roots of  $x^2 - 2x - 3 = 0$  are  $r_+ = 3$  and  $r_- = -1$ .

Then, in the limit, the ratio of  $F_n$  to  $F_{n-1}$  is (1.2), the quadratic formula itself.

Note that rearranging (1.1) to  $F_n = \frac{F_{n+2} - (b)F_{n+1}}{c}$  generates terms leftward of  $F_0$ .

In the above example, terms to the left of  $F_0$  are (moving to the left from zero) 0.333..., -0.222..., 0.259..., -0.247... ... This part of the sequence converges to  $r_-$ .

For convenience, let  $Q = x^2 - bx - c$ . Now, by way of *Binet's formula*,

$$\frac{r_{+}^{n} - r_{-}^{n}}{r_{+} - r_{-}} = F_{n}$$
(1.3)

the zeros (roots) of Q generate the *n*th term of  $\varphi_1$ .

E.g., take the sequence from above, 0, 1, 2, 7, 20, 61, 182, 547, 1640, 4921, 14762.... The roots of  $x^2 - 2x - 3$  are  $r_+ = 3$  and  $r_- = -1$ . Let's say n = 7. Then

$$F_7 = \frac{3^7 - (-1)^7}{3 - (-1)} = 547$$

Thus, the characteristics of a  $\varphi_1$  sequence are encoded in the associated quadratic equation, Q. The patterns associated with a 2<sup>nd</sup>-degree equation are normally, naturally two-dimensional, yet (1.3) generates only a linear sequence. So these  $\varphi_1$  sequences are missing a dimension, it seems...

#### **Dynamic Coefficients**

Indeed, this  $2^{nd}$  dimension has usually evaded observation. It will be opened to exploration by the use of *coefficient lists* (*CLs*). The idea behind such lists is that the coefficients of variables  $F_n$  and  $F_{n+1}$  are themselves a sequence, and terms in a list will apply sequentially as iteration proceeds.

A more general version of (1.1) employs two coefficient lists,  $\beta$  and  $\gamma$ .

$$(\gamma)F_n + (\beta)F_{n+1} = F_{n+2}$$
  $F_0 = 0, F_1 = 1$  (1.4)

Where  $\beta$  and  $\gamma$  are defined as

$$\beta = [b_1, b_2 \dots b_i] \text{ and } \gamma = [c_1, c_2 \dots c_i]$$
(1.5)

Upon the first iteration of (1.4),  $b_1$  and  $c_1$  apply; on the second,  $b_2$  and  $c_2$  and so forth...

Let the indices *i* and  $j = \lambda$ , so  $\beta$  and  $\gamma$  each contain  $\lambda$  terms, where  $1 \le \lambda < \infty$ . Then  $\lambda$  = the *period* of the sequence  $\varphi_{\lambda}$ . Again, as (1.4) is iterated and expanded recursively, the terms in  $\beta$  and  $\gamma$  apply sequentially, and, because  $\lambda < \infty$ , cyclically as iteration goes on. E.g., for  $\lambda \ge 4$ :

$$\varphi_4 = 0, 1, b_1, c_2 + b_1 b_2, c_3 b_1 + c_2 b_3 + b_1 b_2 b_3, c_2 c_4 + c_4 b_1 b_2 + c_3 b_1 b_4 + c_2 b_3 b_4 + b_1 b_2 b_3 b_4 \dots$$
(1.6)

Note that  $\beta$  and  $\gamma$ , as defined in (1.5) may have different lengths; i.e., *i* and *j* respectively. So, given  $i \neq j$ , we say that  $\lambda = \text{LCM}(i,j)$ . This is for economy; else given, say, a  $\varphi_6$  sequence such as  $\beta = [1,2,3], \gamma = [1,2]$ , it would be necessary to write out  $\beta = [1,2,3,1,2,3]$  and  $\gamma = [1,2,1,2,1,2]$ .

Now, for an example of numerical coefficient lists in action, take  $\beta$  [1,2,3] and  $\gamma$  [1] (where, for brevity, the = sign is implied). At the first iteration, the first (and only) term in  $\gamma$  applies to  $F_0$ , and the first term in  $\beta$  to  $F_1$ : (1)0 + (1)1 = 1 = F\_2. On the second iteration, c = 1 and  $b_2 = 2$  apply to  $F_1$  and  $F_2$  respectively: (1)1 + (2)1 = 3 = F\_3. For the third iteration, (1)1 + (3)3 = 10 = F\_4, and on the fourth iteration, the cycle starts over. This process generates the sequence below:

 $\varphi_3 = 0, 1, 1, 3, 10, 13, 36, 121, 157, 435, 1462, 1897, 5256, 17665, 22921, 63507, 213442...$ 

#### **Convergence to Multiple Limit Ratios**

Where the example sequence on page 1,  $\varphi_1 = 0, 1, 2, 7, 20, 61, 182, 547, 1640, 4921...$  converges to the roots of a single equation, the positive (to the rightward of zero) section of  $\varphi_3$  above converges simultaneously to three ratios. These are, in the limit, the roots  $r_{1+} \approx 3.3609$ ,  $r_{2+} \approx 1.2975$  and  $r_{3+} \approx 2.7707$ , i.e., the positive roots of the quadratic equations  $Q_1, Q_2$  and  $Q_3$ . The next step is to fashion a procedure that generates  $Q_i$  coefficients in terms of elements in  $\beta$  and  $\gamma$ .

It will be seen that, in general, a sequence  $\varphi_{\lambda}$  converges to the roots of  $\lambda$  quadratics,  $Q_j$ . The approach to finding the coefficients of these equations will be an extension of the strategy used for  $\varphi_1$ . To this end, (1.4) is expanded below in an indeterminate form.

As just explicated, a sequence  $\varphi_{\lambda}$  is generated by applying the terms of  $\beta$  and  $\gamma$  in order as iterations are performed. At the first iteration, the initial terms  $F_0$  and  $F_1$  are multiplied by  $c_1$  and  $b_1$  respectively; at the second,  $F_1$  and  $F_2$  are multiplied by  $c_2$  and  $b_2$ ; for the third iteration  $c_3$  and  $b_3$  apply and so on. After  $\lambda$  iterations, the cycle repeats, and for  $\lambda \ge 2$  the sequence begins like this:

$$\varphi_{\lambda \ge 2} = F_0, F_1, (c_1)F_0 + (b_1)F_1 = F_2, (c_2)F_1 + (b_2)F_2 = F_3 = (b_2c_1)F_0 + (b_1b_2 + c_2)F_1.$$

It so happens that there is sufficient information in any such single sequence to provide the coefficients for all of the associated  $Q_j$ . This approach, however, spawns huge, unwieldy equations for even relatively small  $\lambda$ . As an alternative, note that any of the  $\varphi_{\lambda}$  has  $\lambda-1$  'siblings', and such information as they hold is most readily accessed when they are all utilized in tandem to form an *array*. The key to generating the constituent sequences of an array is the cyclical permutation of elements in the CLs  $\beta$  and  $\gamma$ . E.g., take  $\beta' [b_2, b_3... b_{\lambda}, b_1]$  and  $\gamma' [c_2, c_3... c_{\lambda}, c_1]$ , which generates the sequence  $\varphi_{\lambda}'$ :

$$\varphi_{\lambda}' = F_0, F_1, (c_2)F_0 + (b_2)F_1 = F_2, (c_3)F_1 + (b_3)F_2 = F_3 = (b_3c_2)F_0 + (b_2 b_3 + c_3)F_1...$$

Then  $\beta''$  [ $b_3, b_4...b_{\lambda}, b_1, b_2$ ],  $\gamma''$  [ $c_3, c_4...c_{\lambda}, c_1, c_2$ ] generates the sequence  $\varphi_{\lambda}''$ ;  $\beta'''$  [ $b_4...b_{\lambda}, b_1, b_2, b_3$ ],  $\gamma'''$  [ $c_4...c_{\lambda}, c_1, c_2, c_3$ ] generates  $\varphi_{\lambda}'''$  and so on.

Such  $\varphi_{\lambda}$  sequences aligned as an array will be represented as  $\Phi_{\lambda} [b_1, b_2 \dots b_i] [c_1, c_2 \dots c_j]$ . The familial  $\varphi_{\lambda}$  sequences such arrays comprise will be notated as  $S_j$ . On route to this, though, is a look at ways in which a single  $\varphi_{\lambda}$  sequence may be mined for the coefficients of all its associated  $Q_s$ .

#### A Single-Sequence Method that Solves for Q<sub>j</sub>'s Coefficients

Take for example a  $\varphi_2$  sequence,  $\beta$  [ $b_1$ ,  $b_2$ ],  $\gamma$  [ $c_1$ ,  $c_2$ ]. For  $\beta$ ,  $\gamma \in \mathbb{N}$ , the ratios of adjacent terms rightward of  $F_0 = 0$  converge alternatively to  $r_{1+}$  and  $r_{2+}$ , the positive roots of the equations  $Q_1$  and  $Q_2$  respectively. To find the coefficients of  $Q_1$ , we take four consecutive terms of  $\varphi_2$ :

$$F_n, F_{n+1}, F_{n+2} = (c_1)F_n + (b_1)F_{n+1}, F_{n+3} = (b_2c_1)F_n + (b_1b_2 + c_2)F_{n+1}$$

Then in the limit as  $n \to \infty$ :  $\frac{F_{n+1}}{F_n} = x_1 = \frac{F_{n+3}}{F_{n+2}}$  and  $x_1 \cdot F_{n+2} = F_{n+3}$ Multiplication by  $x_1$  gives  $x_1^2 \cdot F_{n+2} = x_1 \cdot F_{n+3}$  Or, in terms of  $F_n$  and  $F_{n+1}$ :  $(c_1)F_n \cdot x_1^2 + (b_1)F_{n+1} \cdot x_1^2 = (b_2c_1)F_n \cdot x_1 + (b_1b_2 + c_2)F_{n+1} \cdot x_1$ But  $F_n = \frac{F_{n+1}}{x_1}$ , so  $(c_1)F_{n+1} \cdot x_1 + (b_1)F_{n+1} \cdot x_1^2 = (b_2c_1)F_{n+1} + (b_1b_2 + c_2)F_{n+1} \cdot x_1$ 

Then dividing by  $F_{n+1}$  and collecting terms;

$$Q_1 = b_1 x_1^2 - (b_1 b_2 - c_1 + c_2) x_1 - b_2 c_2 = 0$$
(1.7)

For a numerical example, take  $\Phi_2$  [1,2][1]. Iterating per the CL definition gives the sequence:

$$S_1 = 0, 1, 1, 3, 4, 11, 15, 41, 56, 153, 209, 571, 780, 2131..$$

Solve  $F_{11}/F_{10} = 571/209$  for the ratio 2.7321... Now, substituting terms from  $\beta$  [1,2],  $\gamma$ [1] into (1.7) gives  $Q_1 = 1x_1^2 - (1 \cdot 2 - 1 + 1)x_1 - 2$ , which equals  $x_1^2 - 2x_1 - 2$ , and  $r_{1+} \approx 2.7321$ .

The coefficients of  $Q_2$  may then be found by cyclically permuting the indices of the coefficients of  $Q_1$  in (1.7). Another alternative is to start from the beginning, i.e., perform the above derivation of the coefficients with the terms in  $\beta$  and  $\gamma$  permuted, as shown below. Obviously this last method takes more work:

$$F_n$$

$$F_{n+1}$$

$$F_{n+2} = (c_2)F_n + (b_2)F_{n+1}$$

$$F_{n+3} = (b_1c_2)F_n + (b_1b_2 + c_1)F_{n+1}$$

 $Q_2 = b_2 x_2^2 - (b_1 b_2 - c_2 + c_1) x_2 - b_1 c_2$  results from the methods outlined above. Substituting from  $\beta$ [1,2],  $\gamma$ [1] gives  $2x_2^2 - 2x_2 - 1$ , and  $r_{2+} \approx F_{12}/F_{11} = 780/571 \approx 1.3660$ .

In the above, an equation that solves for the roots of  $Q_1$  derives from setting  $x_1$  equal to two fractions expressed indeterminately in terms of  $\varphi_2$ . Cyclical permutation of  $Q_1$ 's indices then provides the coefficients for  $Q_2$ .

Next, for any  $\lambda$ , a general method for expressing the coefficients of  $Q_1$  in terms of the elements in  $\beta$  and  $\gamma$  is described in the following steps:

- Beginning at  $F_n$ , take  $\lambda + 2$  consecutive terms of  $\varphi_{\lambda}$ :  $F_n$ ,  $F_{n+1}$ ...  $F_{n+\lambda}$ ,  $F_{n+\lambda+1}$
- Then set  $\frac{F_{n+1}}{F_n} = x_1 = \frac{F_{n+\lambda+1}}{F_{n+\lambda}}$
- Clear the second fraction and express  $F_{n+\lambda}$  and  $F_{n+\lambda+1}$  in terms of  $F_n$  and  $F_{n+1}$
- Multiply both sides by  $x_1$ . Now, because  $F_n = \frac{F_{n+1}}{x_1}$ ,  $F_n$  can be transformed to  $F_{n+1}$
- Dividing out  $F_{n+1}$  then leaves a 2<sup>nd</sup> degree equation in  $x_1$ , with coefficients that are expressed as terms from  $\beta$  and  $\gamma$

A derivation of  $Q_1$  for a  $\varphi_3$  sequence  $\beta$  [ $b_1$ ,  $b_2$ ,  $b_3$ ],  $\gamma$  [ $c_1$ ,  $c_2$ ,  $c_3$ ] by this regime is illustrative. First, the sequence is expanded out to five terms:

$$F_n$$

$$F_{n+1}$$

$$F_{n+2} = (c_1)F_n + (b_1)F_{n+1}$$

$$F_{n+3} = (b_2c_1)F_n + (b_1b_2 + c_2)F_{n+1}$$

$$F_{n+4} = (b_2b_3c_1 + c_1c_3)F_n + (b_1b_2b_3 + b_1c_3 + b_3c_2)F_{n+1}$$

Next the fraction in  $x_1 = \frac{F_{n+4}}{F_{n+3}}$  is cleared;  $x_1 \cdot F_{n+3} = F_{n+4}$  is then expressed in terms of  $F_n$  and  $F_{n+1}$ , and both sides are multiplied by  $x_1$ :

and both sides are multiplied by  $x_1$ :

$$(b_2c_1)F_n \cdot x_1^2 + (b_1b_2 + c_2)F_{n+1} \cdot x_1^2 = (b_2b_3c_1 + c_1c_3)F_n \cdot x_1 + (b_1b_2b_3 + b_1c_3 + b_3c_2)F_{n+1} \cdot x_1 \quad (1.8)$$

Substitute  $F_n = \frac{F_{n+1}}{x_1}$  in the above to eliminate  $F_n$ , divide by  $F_{n+1}$  and collect the terms:

$$Q_{1} = (b_{1}b_{2} + c_{2})x_{1}^{2} - (b_{1}b_{2}b_{3} + b_{1}c_{3} - b_{2}c_{1} + b_{3}c_{2})x_{1} - (b_{2}b_{3}c_{1} + c_{1}c_{3}) = 0$$
(1.9)

Then for  $\beta$ [1,2,3] and  $\gamma$ [1,2,3],  $Q_1 = 4x_1^2 - 13x_1 - 9$ . Cyclical permutation of the terms in  $\beta$  and  $\gamma$ , (or indices of  $Q_1$  in (1.9)) gives  $Q_2 = 9x_1^2 - 5x_1 - 8$  and  $Q_3 = 4x_1^2 - 11x_1 - 12$ . So, the sequence  $S_1 = 0, 1, 1, 4, 15, 19, 68, 261...$  converges to the roots  $r_{1+} \approx 3.8364, r_{2+} \approx 1.2607$  and  $r_{3+} \approx 3.5865$ .

As a practical matter, this method is impossibly cumbersome for large  $\lambda$ : e.g., a  $\lambda = 10$  equation has about 1600  $\beta$  and  $\gamma$  terms in its coefficients. This difficulty however is to a great extent mitigated in the context of numerical arrays, in which burdensome symbolic expressions are reduced to more compact and manageable terms. The symbols-to-numbers transition is described below.

#### A General Quadratic in $x_i$ Solves for All Limit Ratios in $\Phi_{\lambda}$

To show how the process of deriving coefficients may be simplified, a set of  $\varphi_3$  sequences will be aligned as an array. Where the symbolic version of the  $\varphi_3$  sequence in (1.8) provided components for the coefficients of  $Q_1$  in (1.9), now set  $F_n = F_0$  to construct a more specific array.

Table I

Once the sequence  $S_1$  is established, cycling the indices of the coefficients gives  $S_2$  and  $S_3$ .

|              | -  |  |       |
|--------------|--|--|-------|
| $\Phi_3 [b]$ | $[b_1, b_2, b_3][c_1, c_2, c_3]$                             |  |       |
|              | $S_1$  | $S_2$  | $S_3$ |
| $F_0 =$      | $F_0$  | $F_0$  | $F_0$ |
| $F_1 =$      | $F_1$  | $F_1$  | $F_1$ |
| $F_2 =$      | $(c_1)F_0 + (b_1)F_1$  | $(c_2)F_0 + (b_2)F_1$  |       |
| $F_3 =$      | $(b_2c_1)F_0 + (b_1b_2 + c_2)F_1$                            | $(b_3c_2)F_0 + (b_2b_3 + c_3)F_1$                            |       |
| $F_4 =$      | $(b_2b_3c_1 + c_1c_3)F_0 + (b_1b_2b_3 + b_1c_3 + b_3c_2)F_1$ | $(b_1b_3c_2 + c_1c_2)F_0 + (b_1b_2b_3 + b_1c_3 + b_2c_1)F_1$ |       |
|              |  |  |       |

Multiplying by  $F_0 = 0$  and  $F_1 = 1$  simplifies the situation considerably, but at the cost of causing some important information to disappear.

Table 2
$$S_1$$
 $S_2$  $S_3$  $F_1$ 11 $F_2$  $b_1$  $b_2$  $F_3$  $b_1b_2+c_2$  $b_2b_3+c_3$  $b_1b_3+c_1$  $F_4$  $b_1b_2b_3+b_1c_3+b_3c_2$  $b_1b_2b_3+b_2c_1+b_1c_3$  $b_1b_2b_3+b_3c_2+b_2c_1$ 

A comparison of the two tables shows that the zeroing out process has resulted in a substantial loss of information, the importance of which will be emphasized below. As a first step, recall that coefficients in (1.9) were found by setting  $F_4$  over  $F_3$ , which can be written like so:

$$x_{1} = \frac{(b_{2}b_{3}c_{1} + c_{1}c_{3})F_{n} + (b_{1}b_{2}b_{3} + b_{1}c_{3} + b_{3}c_{2})F_{n+1}}{(b_{2}c_{1})F_{n} + (b_{1}b_{2} + c_{2})F_{n+1}}$$
(1.10)

The information in (1.10) is that which was manipulated to produce the formula in (1.9). However, in numerical expressions of  $\varphi_{\lambda}$ , the initial terms as defined in (1.4) are always  $F_0 = 0$  and  $F_1 = 1$ . Thus when  $F_n = F_0$ , the terms  $(b_2c_1)$  and  $(b_2b_3c_1 + c_1c_3)$  in the equation above will vanish, and so (1.9) would now appear as below:

$$Q_1 = (b_1 b_2 + c_2) x_1^2 - (b_1 b_2 b_3 + b_1 c_3 + b_3 c_2) x_1 = 0$$
(1.11)

(1.11) solves for  $F_4/F_3$ , but not for the numbers to which  $S_1$  in  $\Phi_3$  will eventually converge. Yet reference to table 2 shows that, serendipitously, terms zeroed from  $S_1$  are available in the adjacent column  $S_2$ , although each lacks their common factor  $c_1$ . The strategy then is to construct a formula (expressed as an equation) that restores (1.11) to (1.9).

The self-similarity inherent in these arrays allows a generalized formula to apply to every column in  $\Phi_{\lambda}$ . We'll see later that this generality extends to rows (periodically) as well. For now, we'll reference the array in table 2 and derive a general formula from that.

Some additional notation facilitates expression of the formula, so let the initial index of  $F_{i,j}$  (where  $F_{i,j}$  is an element of an array  $\Phi_{\lambda}$ ) represent its row and the second its column. Then let  $F_{\lambda}$  denote the 'baseline' of  $\Phi_{\lambda}$ ; that is, the row of number coincident with the period length  $\lambda$ . The coefficients *a*, *b* and *c* of  $Q_j$  are now expressed in terms from  $S_j$  and  $S_{j+1}$  in the following manner:

$$a = F_{\lambda,j}; \quad b = (F_{\lambda+1,j} - c_j \cdot F_{\lambda-1,j+1}); \quad c = -c_j \cdot F_{\lambda,j+1}$$
(1.12)

Referencing (1.11) vis-à-vis (1.9), nothing was lost from the *a* coefficient, so in (1.12) it remains unchanged. The *c* coefficient of (1.9), missing entirely in (1.11), is retrieved from the next column  $(S_{j+1})$  on the same row. Then still in that same column, but going back a row, we find the fragment that completes *b*. The new equation looks like this:

$$Q_{j} = F_{\lambda,j} \cdot x_{j}^{2} - (F_{\lambda+1,j} - c_{j} \cdot F_{\lambda-1,j+1}) x_{j} - c_{j} \cdot F_{\lambda,j+1}$$
(1.13)

For verification, (1.13) applied to  $S_1$  and  $S_2$  in table 2 returns the equation in (1.9):

$$Q_{1} = (b_{1}b_{2} + c_{2})x_{1}^{2} - (b_{1}b_{2}b_{3} + b_{1}c_{3} - b_{2}c_{1} + b_{3}c_{2})x_{1} - (b_{2}b_{3}c_{1} + c_{1}c_{3})$$

#### Extraction of Q<sub>i</sub> Coefficients from Numerical Arrays

If  $\Phi_3$  is generated in numerical form, though, then this clutter of symbols disappears. E.g., take the array  $\Phi_3$  [1,2,3][1,2,3] in table 3 below.

| Table 3                 |       |                       |       |       |       |       |       |       |  |
|-------------------------|-------|-----------------------|-------|-------|-------|-------|-------|-------|--|
| $\Phi_3$ [1,2,3][1,2,3] | $S_1$ | $S_2$                 | $S_3$ | $S_1$ | $S_2$ | $S_3$ | $S_1$ | $S_2$ |  |
| $F_0$                   | 0     | 0                     | 0     | 0     | 0     | 0     | 0     | 0     |  |
| $F_1$                   | 1     | 1                     | 1     | 1     | 1     | 1     | 1     | 1     |  |
| $F_2$                   | 1     | <b>≁</b> <sup>2</sup> | 3     | 1     | 2     | 3     | 1     | 2     |  |
| $F_{\lambda}$           | 4     | <b>→</b> 9            | 4     | 4     | 9     | 4     | 4     | 9     |  |
| $F_4$                   | 15    | 11                    | 14    | 15    | 11    | 14    | 15    | 11    |  |
| $F_5$                   | 19    | 40                    | 54    | 19    | 40    | 54    | 19    | 40    |  |
| $F_{2\lambda}$          | 68    | 153                   | 68    | 68    | 153   | 68    | 68    | 153   |  |
| $F_7$                   | 261   | 193                   | 244   | 261   | 193   | 244   | 261   | 193   |  |
| $F_8$                   | 329   | 692                   | 936   | 329   | 692   | 936   | 329   | 692   |  |

The arrows connect the terms of  $\Phi_3$  that are needed to construct the coefficients of  $Q_1$  according to (1.12). Obviously there is information in table 2 well beyond that required to find the coefficients of  $Q_1$ ,  $Q_2$  and  $Q_3$ . This is due partially to redundancy, as columns repeat in a cycle of  $\lambda$ . Since the second index of  $F_{i,j}$  cannot exceed  $\lambda$ , we can forego this needless repetition. As for the rows, the equation in (1.13) requires that  $F_{\lambda+1,j}$  be available, but nothing beyond it. So, in this  $\lambda_3$  example, we need only nine numbers beyond the initial zeros and ones, as seen below:

| $\Phi_{3}[1,$ | 2,3][1,2 | ,3]   |       | $Q_{j} = F_{\lambda, j} \cdot x_{j}^{2} - (F_{\lambda+1, j} - c_{j} \cdot F_{\lambda-1, j+1})x_{j} - c_{j} \cdot F_{\lambda, j+1}$ |
|---------------|----------|-------|-------|--|
|               | $S_1$    | $S_2$ | $S_3$ |  |
|               |          | 2     |       | $Q_1 = 4x_1^2 - (15 - 1 \cdot 2)x_1 - 1 \cdot 9 = 4x_1^2 - 13x_1 - 9$  |
| $F_{\lambda}$ | 4        | 9     | 4     | $Q_2 = 9x_2^2 - (11 - 2 \cdot 3)x_2 - 2 \cdot 4 = 9x_2^2 - 5x_2 - 8$   |
| $F_4$         | 15       | 11    | 14    | $Q_3 = 4x_3^2 - (14 - 3 \cdot 1)x_3 - 3 \cdot 4 = 4x_3^2 - 11x_3 - 12$   |

Applying (1.13) to  $\Phi_3$  [1,2,3][1,2,3] above gives the equations  $Q_1$ ,  $Q_2$  and  $Q_3$ . The sequence 0, 1, 1, 4, 15, 19, 68, 261... therefore converges to the positive roots  $r_{1+} \approx 3.8364$ ,  $r_{2+} \approx 1.2607$  and  $r_{3+} \approx 3.5865$ , as do sequences  $S_2$  and  $S_3$ . These  $Q_j$  are the same equations that were stated earlier, just below the equation in (1.9), verifying that the single sequence and numerical array methods provide identical results.

Note that equations generated in this way share a single discriminant (*D*), where  $D = b^2 - 4ac$ . In the case just above,  $D = 13^2 - 4 \cdot -9 = 5^2 - 9 \cdot -8 = 11^2 - 4 \cdot -12 = 313$ . This will be of interest later on.

Writer's note:

Steps involved in generalizing  $\varphi$ -sequences to arrays have, so far, been mostly empirical; picking out patterns by inspection and trying to describe them in terms of algorithms and equations. Up to now, it's been obvious and easy, and a rigorous approach to the material would probably not require much more effort or time. But it's about to become more complicated, and it may present some challenges to describe what follows in a formal style of theorems and proofs. Or maybe not. This article will continue with the 'heuristic' approach to identifying patterns. Anyone who discovers something of interest here is, of course, welcome to try to tighten it up...

#### Generalizations of $Q_i$

The equation for  $Q_j$  in the form of (1.13) is a useful tool for delving into patterns that emerge as Fibonacci-type sequences with  $\lambda > 1$  are expanded into two dimensions. Yet there are expressions of (1.13) that have greater generality and power. These versions have qualities that are requisite for certain applications, such as the 2D version of Binet's formula that appears in (1.17). Derivations of these generalized formulas are presented below.

First up is to enhance our notation. Let  $Q = ax^2 + bx + c = 0$  and we seek to define, for k an integer, the symbol  $Q^k$ . To this end, consider  $r_{\pm}$ , the roots of Q, in their quadratic formula representation.

$$r_{+} = \frac{-b + \sqrt{b^2 - 4ac}}{2a}$$
 and  $r_{-} = \frac{-b - \sqrt{b^2 - 4ac}}{2a}$ 

Then, since  $Q = x^2 - (r_+ + r_-)x + (r_+ \cdot r_-)$ , we say  $Q^{-1} = x^2 - (r_+^{-1} + r_-^{-1})x + (r_+^{-1} \cdot r_-^{-1})$ 

Thus if 
$$r_{+}^{-1} = \frac{2a}{-b + \sqrt{b^2 - 4ac}}$$
 and  $r_{-}^{-1} = \frac{2a}{-b - \sqrt{b^2 - 4ac}}$   
then  $r_{+}^{-1} + r_{-}^{-1} = \frac{-4ab}{4ac}$  and  $r_{+}^{-1} \cdot r_{-}^{-1} = \frac{4a^2}{4ac}$ . Hence  
 $Q^{-1} = x^2 - (\frac{-4ab}{4ac})x + \frac{4a^2}{4ac} = 4acx^2 + 4abx + 4a^2 = cx^2 + bx + a$ 

Also, 
$$Q \cdot Q^{-1} = Q^0 = x^2 - (r_+ \cdot r_+^{-1} + r_- \cdot r_-^{-1})x + (r_+ \cdot r_+^{-1} \cdot r_- \cdot r_-^{-1}) = x^2 - 2x + 1$$

In general,  $Q^k = x^2 - (r_+^k + r_-^k)x + (r_+^k \cdot r_-^k)$ . We'll walk through  $Q^2$  for an example.

$$r_{+}^{2} = \frac{b^{2} - 2ac - b\sqrt{b^{2} - 4ac}}{2a^{2}} \text{ and } r_{-}^{2} = \frac{b^{2} - 2ac + b\sqrt{b^{2} - 4ac}}{2a^{2}}, \text{ so}$$
$$Q^{2} = x^{2} - (r_{+}^{2} + r_{-}^{2})x + (r_{+}^{2} \cdot r_{-}^{2}) = x^{2} - \frac{(b^{2} - 2ac)}{a^{2}}x + \frac{c^{2}}{a^{2}} = a^{2}x^{2} - (b^{2} - 2ac)x + c^{2}$$

For instance, if  $Q = x^2 - 2x - 3$ , then  $Q^2 = x^2 - 10x + 9$ .

Note then, that for any k the radical vanishes, and the coefficients of  $Q^k$  will be integers. Now consider two distinct equations, Q and Q'. We will soon see, in numerous examples, that it is possible for the equation below to also have integer coefficients.

$$Q'' = Q \cdot Q' = x^2 - (r_+ \cdot r_+' + r_- \cdot r_-')x + (r_+ \cdot r_+' \cdot r_- \cdot r_-')$$

A seemingly necessary, but apparently not sufficient condition on this is that the discriminants of Q and Q' are of equal value (or one is an  $n^2$  multiple of the other, n = 1, 2, 3...). That said, take the formula in (1.1) and set b = c = 1. Iteration in both directions yields the familiar Fibonacci sequence.

This is of course, <u>the</u> Fibonacci sequence, (call it  $\phi$ ), which converges, leftward and rightward to the roots of  $Q_{\phi} = x^2 - x - 1 = 0$ . The positive root of this equation (or its inverse) is often called the divine proportion or the golden mean:  $r_{\phi+} = 1.618034..., r_{\phi-} = -0.618034...$  The relationship of the roots of  $Q_{\phi}$  to its *b* and *c* coefficients is  $r_{\phi+} + r_{\phi-} = -b$  and  $r_{\phi+} \cdot r_{\phi-} = c$ .

Consequently, equations  $Q_{\phi}^{k}$  with integral coefficients are formed on the roots  $r_{\phi\pm}^{k}$  by

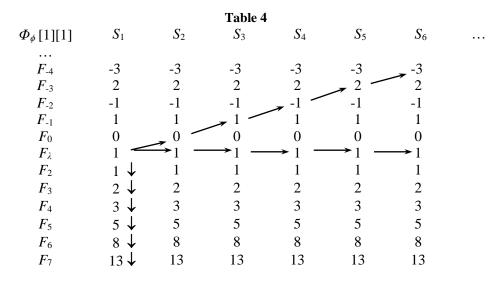
$$Q_{\phi}^{k} = x^{2} - (r_{\phi+}^{k} + r_{\phi-}^{k})x + r_{\phi+}^{k} \cdot r_{\phi-}^{k}$$

But  $r_{\phi_+}^n + r_{\phi_-}^n$  is also a formula for  $L_n$ , which is the *n*<sup>th</sup> term of the Lucas sequence, where

$$(c)L_n + (b)L_{n+1} = L_{n+2}; \quad L_0 = 2, L_1 = 1$$

Therefore, for b = c = 1, the <u>Lucas numbers</u> 2, 1, 3, 4, 7, 11, 18, 29, 47... are the *b* coefficients of successive  $Q_{\phi}^{k}$ ,  $k \ge 0$ . But the Lucas sequence can also be stated in terms of Fibonacci numbers as  $L_{n} = F_{n-1} + F_{n+1}$ . With that in mind, the array  $\Phi_{\phi}[1][1]$  is constructed on  $\phi$  below.

Note: the patterns that associated adjacent table 3 elements to form the coefficients of  $Q_j$  in (1.13) are, as indicated by the arrows, extended throughout this array.



The equalities in (1.12) are now modified to form coefficients for  $Q_{\phi}^{k}$  from terms in  $\Phi_{\phi}$ :

$$a = F_{1,1}; \ b = -(F_{1+k,1} + F_{1-k,1+k} \cdot (-1)^k); \ c = F_{1,1+k} \cdot (-1)^k$$
$$Q_{\phi}^k = F_{1,1}x^2 - (F_{1+k,1} + F_{1-k,1+k} \cdot (-1)^k)x + F_{1,1+k} \cdot (-1)^k$$
(1.14)

According to the patterns in  $\Phi_{\phi}$ , the *a* coefficient is fixed. As *k* increases, the first term of *b* moves downward through the first column as the second term moves up and rightward on a diagonal, while

*c* slides across the columns on the first row. From (1.14), which embodies these patterns, the coefficients of  $Q_{\phi}^{k}$  for the first few values of *k* are given below.

$$\begin{aligned} Q_{\phi}^{0} &= 1x^{2} - (1 + 1(-1)^{0})x + 1(-1)^{0} = x^{2} - 2x + 1\\ Q_{\phi}^{1} &= 1x^{2} - (1 + 0(-1)^{1})x + 1(-1)^{1} = x^{2} - x - 1\\ Q_{\phi}^{2} &= 1x^{2} - (2 + 1(-1)^{2})x + 1(-1)^{2} = x^{2} - 3x + 1\\ Q_{\phi}^{3} &= 1x^{2} - (3 - 1(-1)^{3})x + 1(-1)^{3} = x^{2} - 4x - 1\\ Q_{\phi}^{4} &= 1x^{2} - (5 + 2(-1)^{4})x + 1(-1)^{4} = x^{2} - 7x + 1\\ Q_{\phi}^{5} &= 1x^{2} - (8 - 3(-1)^{5})x + 1(-1)^{5} = x^{2} - 11x - 1\end{aligned}$$

#### Further Generalizations of $Q_J$

This shows that the array  $\Phi_{\phi}$  contains the information necessary to provide the coefficients for the equations with roots that are powers of  $r_{\phi+}$ . This extended pattern also applies to higher order arrays, but to navigate those realms requires modifying the notation again. To this end, the symbol for the exponent is now enhanced.

By convention,  $B^k$  represents *B* to the power of *k*. For *k* a positive integer, this is a simple and convenient product notation; e.g.,  $B^3 = B \cdot B \cdot B$ . We will say in this case that the base *B* is *monotonic*. Now an underline to  $\underline{k}$  in  $B^{\underline{k}}$  will be taken to signify that *B* is *polytonic*, i.e., it comprises two or more elements (not necessarily all distinct), say,  $B = [t_1, t_2...t_{\lambda}]$ .

A property of *B* is that any term may be the initial, as in  $t_j^k$ . For example, take  $B = \beta = [b_1, b_2, b_3]$ ; then  $b_2^5 = b_2 \cdot b_3 \cdot b_1 \cdot b_2 \cdot b_3$ . (A potential problem with the underline arises in this example, as the superscript and subscript in  $b_2^5$  may seem to combine to form a fraction. Bear in mind that fractional exponents will not appear in this context; but if they were necessary, the somewhat clumsy but still intelligible  $b_{\frac{j_1}{2}}^{\frac{j_2}{2}}$  could be used.)

For further examples, if  $\beta$  [1,2,3] then,  $b_2^5 = 2 \cdot 3 \cdot 1 \cdot 2 \cdot 3 = 36$ . For  $\Phi_3$  [1,2,3][2] then  $c_2^7 = 2^7$ . It could be said that in general any base is polytonic, but when each  $t_i = t_j$ , then  $\underline{k} = k$ . Note that for a base *B* of  $\lambda$  elements, the product  $t_j^{\underline{\lambda}}$  is the same regardless of which element is chosen as the initial; i.e.,  $t_1^{\underline{\lambda}} = t_2^{\underline{\lambda}} = t_3^{\underline{\lambda}} \dots = t_j^{\underline{\lambda}}$ .

*Problem*: Define the exponent  $\underline{k}$  for negative integer values

Having derived (1.14) empirically from table 4, we go on to devise a more general expression for a period- $\lambda$  pattern,  $\Phi_{\lambda}$  [ $b_1, b_2 \dots b_{\lambda}$ ][1]. Beyond underlining exponents, it's necessary to ensure that the coefficients are taken from the proper rows. This entails that *F*'s initial subscript be equal to the period  $\lambda$ ; thus every coefficient in the new formula will preface as  $F_{\lambda}$ , as in, say,  $F_{\lambda-k,j+k}$ .

With this established, the pattern used in  $\Phi_{\phi}$  (table 4) applies now to the array  $\Phi_3$  [1,2,3][1], in table 5 below:

|                     |       |                       | Та         | ble 5      |            |                      |       |            |
|---------------------|-------|-----------------------|------------|------------|------------|----------------------|-------|------------|
| $\Phi_3$ [1,2,3][1] | $S_1$ | $S_2$                 | $S_3$      | $S_4$      | $S_5$      | $S_6$                | $S_7$ | $S_8$      |
| <br>F <sub>-4</sub> | -11   | -10                   | -9         | -11        | -10        | -9                   | -11   | 10         |
| $F_{-\lambda}$      | 3     | -10                   | 4          | 3          | -10        | 4                    | 3     | <b>≁</b> 7 |
| $F_{-2}$            | -2    | -3                    | -1         | -2         | -3         | <b>_</b> -1 <b>/</b> | -2    | -3         |
| $F_{-1}$            | 1     | 1                     | 1          | 1          | <b>y</b> 1 | 1                    | 1     | 1          |
| $F_0$               | 0     | 0                     | 0          | <u>→</u> 0 | 0          | 0                    | 0     | 0          |
| $F_1$               | 1     | 1                     | <b>,</b> 1 | 1          | 1          | 1                    | 1     | 1          |
| $F_2$               | 1     | <b>x</b> <sup>2</sup> | 3          | 1          | 2          | 3                    | 1     | 2          |
| $F_\lambda$         | 3 🧹   | → 7 —                 | → 4 —      | → 3 —      | → 7 —      | → 4 —                | → 3 — | <b>→</b> 7 |
| $F_4$               | 10↓   | 9                     | 11         | 10         | 9          | 11                   | 10    | 9          |
| $F_5$               | 13↓   | 25                    | 37         | 13         | 25         | 37                   | 13    | 25         |
| $F_{2\lambda}$      | 36↓   | 84                    | 48         | 36         | 84         | 48                   | 36    | 84         |
| $F_7$               | 121↓  | 109                   | 133        | 121        | 109        | 133                  | 121   | 109        |
| $F_8$               | 157↓  | 302                   | 447        | 157        | 302        | 447                  | 157   | 302        |

Define  $Q_j^{\underline{k}}$  as the equation that has its *a* coefficient in the column  $S_j$ , and the roots of which are the product  $r_j \cdot r_{j+1} \cdot \ldots \cdot r_{j+k-1} = r_j^{\underline{k}}$  for  $r_{j+}$  and  $r_{j-}$ .

The equation in (1.14) is generalized so as to derive the coefficients of  $Q_j^{\underline{k}}$  from  $\Phi_{\lambda}$ :

$$Q_{j}^{k} = F_{\lambda,j} \cdot x^{2} - (F_{\lambda+k,j} + F_{\lambda-k,j+k} \cdot (-1)^{k})x + F_{\lambda,j+k} \cdot (-1)^{k}$$
(1.15)

The formula in (1.15) is now applied to the  $\Phi_3$  array above, and the coefficients of  $Q_1^{\underline{k}}$  for a few sequential values of *k* are shown below.

$$Q_{1}^{0} = 3x_{1}^{2} - (3 + 3(-1)^{0})x_{1} + 3(-1)^{0} = x_{1}^{2} - 2x_{1} + 1$$

$$Q_{1}^{1} = 3x_{1}^{2} - (10 + 2(-1)^{1})x_{1} + 7(-1)^{1} = 3x_{1}^{2} - 8x_{1} - 7$$

$$Q_{1}^{2} = 3x_{1}^{2} - (13 + 1(-1)^{2})x_{1} + 4(-1)^{2} = 3x_{1}^{2} - 14x_{1} + 4$$

$$Q_{1}^{\underline{\lambda}} = 3x_{1}^{2} - (36 + 0(-1)^{3})x_{1} + 3(-1)^{3} = x_{1}^{2} - 12x_{1} - 1$$

$$Q_{1}^{4} = 3x_{1}^{2} - (121 + 1(-1)^{4})x_{1} + 7(-1)^{4} = 3x_{1}^{2} - 122x_{1} + 7$$

$$Q_{1}^{5} = 3x_{1}^{2} - (157 - 1(-1)^{5})x_{1} + 4(-1)^{5} = 3x_{1}^{2} - 158x_{1} - 4$$

Moving over to the column  $S_2$  gives the coefficients for  $Q_2^k$ .

$$\begin{aligned} Q_2^0 &= 7x_2^2 - (7+7(-1)^0)x_2 + 7(-1)^0 &= x_2^2 - 2x_2 + 1\\ Q_2^1 &= 7x_2^2 - (9+3(-1)^1)x_2 + 4(-1)^1 &= 7x_2^2 - 6x_2 - 4\\ Q_2^2 &= 7x_2^2 - (25+1(-1)^2)x_2 + 3(-1)^2 &= 7x_2^2 - 26x_2 + 3\\ Q_2^{\underline{\lambda}} &= 7x_2^2 - (84+0(-1)^3)x_2 + 7(-1)^3 &= x_2^2 - 12x_2 - 1\\ Q_2^{\underline{\lambda}} &= 7x_2^2 - (109+1(-1)^4)x_2 + 4(-1)^4 &= 7x_2^2 - 110x_2 + 4\\ Q_2^{\underline{5}} &= 7x_2^2 - (302-2(-1)^5)x_2 + 3(-1)^5 &= 7x_2^2 - 304x_2 - 3 \end{aligned}$$

A shift to  $S_3$  gives the coefficients for  $Q_3^{\underline{k}}$ .

$$Q_{3}^{0} = 4x_{3}^{2} - (4 + 4(-1)^{0})x_{3} + 4(-1)^{0} = x_{3}^{2} - 2x_{3} + 1$$

$$Q_{3}^{1} = 4x_{3}^{2} - (11 + 1(-1)^{1})x_{3} + 3(-1)^{1} = 4x_{3}^{2} - 10x_{3} - 3$$

$$Q_{3}^{2} = 4x_{3}^{2} - (37 + 1(-1)^{2})x_{3} + 7(-1)^{2} = 4x_{3}^{2} - 38x_{3} + 7$$

$$Q_{3}^{\underline{\lambda}} = 4x_{3}^{2} - (48 + 0(-1)^{3})x_{3} + 4(-1)^{3} = x_{3}^{2} - 12x_{3} - 1$$

$$Q_{3}^{\underline{\lambda}} = 4x_{3}^{2} - (133 + 1(-1)^{4})x_{3} + 3(-1)^{4} = 4x_{3}^{2} - 134x_{3} + 3$$

$$Q_{3}^{\underline{5}} = 4x_{3}^{2} - (447 - 3(-1)^{5})x_{3} + 7(-1)^{5} = 4x_{3}^{2} - 450x_{3} - 7$$

To reiterate, where the roots of  $Q_{\phi}^{k}$  were powers of  $r_{\phi\pm}$ , the roots of  $Q_{j}^{k}$  here are products of some or all of three limit ratios in  $\Phi_{3}$  [1,2,3][1]. More examples will help to clarify this idea.

Given 
$$S_j$$
 in  $\Phi_{\lambda}$ , then  $\frac{F_{n\lambda+1,j}}{F_{n\lambda,j}} \to r_j$ ;  $\frac{F_{n\lambda+2,j}}{F_{n\lambda+1,j}} \to r_{j+1}$ ;  $\frac{F_{n\lambda+\lambda,j}}{F_{n\lambda+\lambda-1,j}} \to r_{j+1}$  as  $n \to \infty$ .

In  $\Phi_3$  [1,2,3][1],  $S_1 = 0, 1, 1, 3, 10, 13, 36, 121, 157, 435, 1462, 1897, 5256, 17665...$ 

One expression of 
$$\frac{F_{2\lambda+1,1}}{F_{2\lambda,1}}$$
 is  $\frac{17665}{5256}$ , which  $\approx r_{1+}$  of  $Q_1 = 3x_1^2 - 8x_1 - 7 \approx 3.36092$   
$$\frac{F_{2\lambda+2,1}}{F_{2\lambda+1,1}} = \frac{1897}{1462} \approx r_{2+}$$
 of  $Q_2 = 7x_2^2 - 6x_2 - 4 \approx 1.29754$ 
$$\frac{F_{2\lambda+3,1}}{F_{2\lambda+2,1}} = \frac{5256}{1897} \approx r_{3+}$$
 of  $Q_3 = 4x_3^2 - 10x_3 - 3 \approx 2.77069$ 

And  $r_{1+} \cdot r_{2+} \cdot r_{3+} = r_{j+}^{\lambda} \approx 12.08276...$ , a root of  $Q_j^{\lambda} = x^2 - 12x - 1$ .

A look at some discriminants of equations from the array  $\Phi_3$  [1,2,3][1]:

| $Q_{j}^{\underline{k}}$       | D =                   | $Q_{j}^{\underline{k}}$       | D =                   |
|-------------------------------|-----------------------|-------------------------------|-----------------------|
| $Q_1^2 = 3x^2 - 14x + 4$      | 148                   | $Q_2^4 = 7x_2^2 - 110x_2 + 4$ | 148 * 9 <sup>2</sup>  |
| $Q_1^3 = x_1^2 - 12x_1 - 1$   | 148                   | $Q_2^5 = 7x_2^2 - 304x_2 - 3$ | 148 * 25 <sup>2</sup> |
| $Q_1^4 = 3x^2 - 122x + 7$     | $148 * 10^2$          | $Q_3^2 = 4x_3^2 - 38x_3 + 7$  | 148 * 3 <sup>2</sup>  |
| $Q_1^5 = 3x_1^2 - 158x_1 - 4$ | 148 * 13 <sup>2</sup> | $Q_3^3 = x_3^2 - 12x_3 - 1$   | 148                   |
| $Q_2^2 = 7x_2^2 - 26x_2 + 3$  | $148 * 2^2$           | $Q_3^4 = 4x^2 - 134x + 3$     | $148 * 25^2$          |
| $Q_2^3 = x_2^2 - 12x_2 - 1$   | 148                   | $Q_3^{5} = 4x^2 - 450x - 7$   | $148 * 37^2$          |

The basic structure of these  $\Phi_{\lambda}$  arrays is established; yet they have intriguing properties that remain to be explored. More on this is to follow in other articles. At this point, after all the effort to forge

our way into this new territory, suppose we just look around for a while, to try to get a feel for the lay of the land.

#### Fibonacci Identities and Binet's Formula in 2D

First up, we'll examine some of the *Fibonacci identities*. These are formulas that equate various terms of the  $\phi$ -sequence in diverse and often unexpected ways. What happens to such an identity in  $\Phi_{\lambda}$ ? Let's check out some examples...

(i) 
$$F_n^2 + F_{n+1}^2 = F_{2n+1} \rightarrow F_{n\lambda, j} \cdot F_{n\lambda, j+1} \cdot c_j + F_{n\lambda+1, j}^2 = F_{2n\lambda+1, j}$$
 (where  $\rightarrow$  is 'maps to')

(ii) 
$$F_n \cdot (F_{n+1} + F_{n-1}) = F_{2n} \to F_{n\lambda, j} \cdot (F_{n\lambda+1, j} + F_{n\lambda-1, j+1} \cdot c_j) = F_{2n\lambda, j}$$

(iii) 
$$F_n \cdot F_{n+1} - F_{n-1} \cdot F_{n-2} = F_{2n-1} \rightarrow \frac{F_{n\lambda, j} \cdot F_{n\lambda+1, j+1} - F_{n\lambda-1, j+1} \cdot F_{n\lambda-2, j+2} \cdot c_j^2}{b_j} = F_{2n\lambda-1, j+1}$$

These formulas for  $F_{n>0}$  are verified using the array in table 6 below.

| Table 6                 |        |        |        |        |  |  |  |  |
|-------------------------|--------|--------|--------|--------|--|--|--|--|
| $\Phi_3$ [1,1,2,3][1,2] | $S_1$  | $S_2$  | $S_3$  | $S_4$  |  |  |  |  |
| $F_0$                   | 0      | 0      | 0      | 0      |  |  |  |  |
|                         | 1      | 1      | 1      | 1      |  |  |  |  |
|                         | 1      | 1      | 2      | 3      |  |  |  |  |
|                         | 3      | 3      | 8      | 4      |  |  |  |  |
| $F_{\lambda}$           | 7      | 11     | 10     | 10     |  |  |  |  |
|                         | 27     | 14     | 26     | 24     |  |  |  |  |
|                         | 34     | 36     | 62     | 92     |  |  |  |  |
|                         | 88     | 86     | 238    | 116    |  |  |  |  |
| $F_{2\lambda}$          | 210    | 330    | 300    | 300    |  |  |  |  |
|                         | 806    | 416    | 776    | 716    |  |  |  |  |
|                         | 1016   | 1076   | 1852   | 2748   |  |  |  |  |
|                         | 2628   | 2568   | 7108   | 3464   |  |  |  |  |
| $F_{3\lambda}$          | 6272   | 9856   | 8960   | 8960   |  |  |  |  |
|                         | 24072  | 12424  | 23176  | 21384  |  |  |  |  |
|                         | 30344  | 32136  | 55312  | 82072  |  |  |  |  |
|                         | 78488  | 76696  | 212288 | 103456 |  |  |  |  |
| $F_{4\lambda}$          | 187320 | 294360 | 267600 | 267600 |  |  |  |  |
|                         | 718936 | 371056 | 692176 | 638656 |  |  |  |  |
|                         |        |        |        |        |  |  |  |  |

| Beginning with (i): If $F_{n\lambda,j} =$ | $F_{\lambda,1}$ ; then $7 \cdot 11 \cdot 1 + 27^2 = 806 = F_{2\lambda+1,1}$  |
|---|--|
|   | $F_{\lambda,2}$ ; then $11 \cdot 10 \cdot 2 + 14^2 = 416 = F_{2\lambda+1,2}$   |
|   | $F_{2\lambda,3}$ ; then $300 \cdot 300 \cdot 1 + 776^2 = 692176 = F_{4\lambda+1,3}$  |
|   | $F_{2\lambda,4}$ ; then $300 \cdot 210 \cdot 2 + 716^2 = 638656 = F_{4\lambda+1,4}$  |
| For identity (ii): If $F_{n\lambda,j} =$  | $F_{\lambda,1}$ ; then $7 \cdot 27 + 7 \cdot 3 \cdot 1 = 210 = F_{2\lambda,1}$<br>$F_{\lambda,2}$ ; then $11 \cdot 14 + 11 \cdot 8 \cdot 2 = 330 = F_{2\lambda,2}$ |
|   | $E_{1}$ , then 200 776 + 200 116 1 267600 E  |

Then, for (iii): If 
$$F_{n\lambda,j} = F_{\lambda,1}$$
; then  $(7 \cdot 14 - 3 \cdot 2 \cdot (1 \cdot 2))/1 = 86 = F_{2\lambda-1,2}$   
 $F_{\lambda,2}$ ; then  $(11 \cdot 26 - 8 \cdot 3 \cdot (2 \cdot 1))/1 = 238 = F_{2\lambda-1,3}$   
 $F_{2\lambda,4}$ ; then  $(300 \cdot 806 - 88 \cdot 36 \cdot (2 \cdot 1))/3 = 78488 = F_{4\lambda-1,1}$ 

(What happens if these  $\Phi_{\lambda}$  formulas are applied over  $F_{n < 0}$ ? A worksheet is provided on page 20 for those who wish to investigate.)

There is another way to verify these formulas that is rather interesting. Take this version of the sequence  $\varphi_1$ :

$$(-1)F_n + (2)F_{n+1} = F_{n+2} F_0 = y, F_1 = y+1 (1.16)$$

For y an integer, working in both directions generates  $\mathbb{Z}$ , the ring of integers; i.e.:

If we construct the array  $\Phi_{(z)}$  [2][-1] on (1.16) sequences, then the generalized identities i, ii, and iii above reduce to ordinary algebra:

(i) 
$$\rightarrow (-1)y^2 + (y+1)^2 = 2y + 1$$
  
(ii)  $\rightarrow y((y+1) + (-1)(y-1)) = 2y$   
(iii)  $\rightarrow (y(y+1) - (-1)^2(y-1)(y-2))/2 = 2y - 1$ 

Finally, consider a version of Binet's  $\varphi_1$  formula (1.3) generalized for  $\varphi_{\lambda}$ :

$$F_{\lambda,j} \cdot \frac{r_{j+}^n - r_{j-}^n}{r_{j+}^\lambda - r_{j-}^\lambda} = F_{n,j}$$
(1.17)

Recall that  $F_{\lambda,j}$  is the *a* coefficient of  $Q_j^{\underline{n}}$ . We'll verify this formula with an example. Suppose that  $F_{7,1}$  (i.e., the 7<sup>th</sup> term of  $S_1$ ) in  $\Phi_3$  [1,2,3][1,2,3] is to be determined. Then, since  $7 = 2\lambda + 1$ , only the roots of  $Q_1$  and  $Q_j^{\underline{\lambda}}$  are needed. The truncated array

supplies the numbers to plug into the generalized formula

$$Q_{j}^{k} = F_{\lambda, j} \cdot x_{j}^{2} - (F_{\lambda+k, j} + F_{\lambda-k, j+k} \cdot c_{j}^{k} \cdot (-1)^{k}) x_{j} + F_{\lambda, j+k} \cdot c_{j}^{k} \cdot (-1)^{k} = 0$$
(1.18)

Whence  $Q_1 = 4x_1^2 - 13x_1 - 9$ . We can also use this  $\Phi_3$  fragment to find roots of  $Q_j^{\frac{\lambda}{2}}$  by the formula

$$Q_{j}^{\lambda} = x_{j}^{2} - (F_{\lambda+1,j} + c_{j} \cdot F_{\lambda-1,j+1})x_{j} + c_{j}^{\lambda} \cdot (-1)^{\lambda} = 0$$
(1.19)

from which  $Q_j^{\lambda} = x^2 - 17x - 6$ .

The equations in (1.18) and (1.19) above have roots:  $r_{1+} = 3.8365$ ;  $r_{1-} = -.5865$ ;  $r_{j+}^{\lambda} = 17.3459$ ;  $r_{j-}^{\lambda} = -.3459$ . Then  $r_{1+} \cdot r_{j+}^{2\lambda} - r_{1-} \cdot r_{j-}^{2\lambda} = 1154.3202 + .0702 = 1154.3904$ , and  $r_{j+}^{\lambda} - r_{j-}^{\lambda} = 17.6918$ . Then 1154.3904/17.6918 = 65.25, and 65.25  $\cdot 4 = 261$ .  $S_1 = 0, 1, 1, 4, 15, 19, 68, 261$ ... voila!

In closing, here's an array with a monotonic  $\beta$ .

| Table 7             |                                |                               |                                |                                |                               |                                |                                |     |
|---------------------|--------------------------------|-------------------------------|--------------------------------|--------------------------------|-------------------------------|--------------------------------|--------------------------------|-----|
| $\Phi_3$ [1][1,2,3] | $S_1$                          | $S_2$                         | $S_3$                          | $S_4$                          | $S_5$                         | $S_6$                          | $S_7$                          | ••• |
| $F_{-4}$            | - <sup>4</sup> / <sub>18</sub> | - <sup>6</sup> / <sub>6</sub> | - <sup>5</sup> / <sub>12</sub> | - <sup>4</sup> / <sub>18</sub> | - <sup>6</sup> / <sub>6</sub> | - <sup>5</sup> / <sub>12</sub> | - <sup>4</sup> / <sub>18</sub> |     |
| $F_{-\lambda}$      | $\frac{3}{6}$                  | $\frac{4}{6}$                 | $\frac{2}{6}$                  | $\frac{-718}{3}_{6}$           | $\frac{4}{6}$                 | $\frac{2}{6}$                  | $\frac{3}{6}$                  |     |
| $F_{-2}$            | - <sup>1</sup> / <sub>6</sub>  | -1/3                          | -1/2                           | -1/6                           | - <sup>1</sup> / <sub>3</sub> | -1/2                           | $-^{1}/_{6}$                   |     |
| $F_{\text{-1}}$     | $^{1}/_{3}$                    | 1/1                           | $^{1}/_{2}$                    | $^{1}/_{3}$                    | $\rightarrow 1/1$             | $^{1}/_{2}$                    | <sup>1</sup> / <sub>3</sub>    |     |
| $F_0$               | 0←                             | $\overline{}$                 | -0-                            | $\longrightarrow_0$            | 0                             | 0                              | 0                              |     |
| $F_1$               | 1-                             | T                             | $\searrow_1$                   | 1                              | 1                             | 1                              | 1                              |     |
| $F_2$               | 1                              | $\mathbf{N}_1$                | 1                              | 1                              | 1                             | 1                              | 1                              |     |
| $F_\lambda$         | 3                              | 4                             | 2                              | 3                              | 4                             | 2                              | 3                              |     |
| $F_4$               | 6                              | 5                             | 4                              | 6                              | 5                             | 4                              | 6                              |     |
| $F_5$               | 9                              | 13                            | 10                             | 9                              | 13                            | 10                             | 9                              |     |
| $F_{2\lambda}$      | 21                             | 28                            | 14                             | 21                             | 28                            | 14                             | 21                             |     |
| $F_7$               | 48                             | 41                            | 34                             | 48                             | 41                            | 34                             | 48                             |     |

Recall the array  $\Phi_{\phi}$  in table 4, where each column is the  $\phi$ -sequence, 0,1,1,2,3.... The formula for the second part of the *b* coefficient of  $Q_{\phi}^k$  in (1.14) is  $F_{1-k,1+k} \cdot (-1)^k$ . Indices 1-k,1+k entail moving across the array in a diagonal direction; but these same numbers are available in every column, including the first, so in this case, diagonal motion isn't necessary. We could find those same numbers by 1-k,1; i.e., moving straight 'down'.

But in arrays of higher order, adjacent columns (sequences) differ. Thus the indices  $\lambda - k, \lambda + k$  from  $Q_j^k$  in (1.18) will take a term from each sequence in turn. Although this diagonal motion seems, for  $\lambda > 1$ , a necessity, it turns out that if the second element in the *b* coefficient is given a different form, there is apparently a way to construct it from the first column alone. To show how this idea has developed so far, the terms connected by arrows in table 7 above are equated by the formula:

$$F_{\lambda-k,\,j+k}\cdot c_j^{\underline{k}}\cdot (-1)^k = F_{-\lambda+k,\,j}\cdot c_j^{\underline{\lambda}}\cdot (-1)^{\lambda-1}$$

Thus for  $\underline{k} \ge 0$ , if this equivalence is correct, the coefficients of (1.18) can also be stated as:

$$Q_j^{\underline{k}} = F_{\lambda,j} \cdot x_j^2 - (F_{\lambda+k,j} - F_{-\lambda+k,j} \cdot c_j^{\underline{\lambda}} \cdot (-1)^{\lambda}) x_j + F_{\lambda,j+k} \cdot c_j^{\underline{k}} \cdot (-1)^{\lambda}$$

*Problem(s)*: Provide a formal proof for any of the formulas/equations/identities stated above.

#### Addenda

Recall the formula in (1.16) which, for y = an integer generates the ring of integers. This sequence converges to  $F_{n+1}/F_n = 1$  in both directions. It would seem that no integer sequence could converge more slowly than ...-5, -4, -3, -2, -1, 0, 1, 2, 3, 4, 5... But other sequences, found in arrays such as  $\Phi_4$  [1, -2,3, -4][1] and  $\Phi_4$  [-1,2, -3,4][1], do converge at the same rate. The discriminant D = 0 ( $D_0$ ) seems likely be the key here.

How many  $D_0$  arrays are there? An infinite number perhaps? What sort of patterns will emerge in  $\beta$  and  $\gamma$  if all of the  $D_0$  arrays are themselves considered as a set?

The empirical formula below gives, in test cases, coefficients for the equation with roots  $r_{i+}^2$ :

$$Q_{j}^{2} = F_{\lambda, j}^{2} \cdot x_{j}^{2} - (F_{\lambda+1, j}^{2} + F_{\lambda-1, j+1}^{2} - 2c_{j}^{\lambda} \cdot (-1)^{\lambda}) \cdot x_{j} + F_{\lambda, j+1}^{2} \cdot c_{j}^{2}$$

Should this be valid for all arrays, perhaps a pattern is incipient here, and an algorithm can be found to generate coefficients for  $Q_i^k$  when k is an ordinary exponent.

A 3<sup>rd</sup> degree sequence,  $(c)F_n + (b)F_{n+1} + (a)F_{n+2} = F_{n+3}$ , converges to a root (or roots) of a cubic equation. Is there something analogous to the arrays  $\Phi_{\lambda}$  for use in 3D?

For all of the aesthetic expressions of the Fibonacci sequence and the golden mean that abound in geometry, art and nature, are expressions of any  $\Phi_{\lambda>1}$  sequences and/or ratios to be found in these realms?

Using  $\mathbb{N} = 1, 2, 3... \infty$  as a coefficient of the sequence in (1.4), then  $F_n + (\mathbb{N})F_{n+1} = \underline{OEIS} \underline{A058307}$ 

0, 1, 3, 13, 68, 421, 3015, 24541, 223884, 2263381, 25121075, 303716281, 3973432728...

 $(\mathbb{N})F_n + (\mathbb{N})F_{n+1} = \underline{A002467}$ 

0, 1, 1, 4, 15, 76, 455, 3186, 25487, 229384, 2293839, 25232230, 302786759, 3936227868...

 $(\mathbb{N})F_n + F_{n+1} = \underline{A000932}$ 

0, 1, 1, 3, 6, 18, 48, 156, 492, 1740, 6168, 23568, 91416, 374232, 1562640, 6801888...

#### References:

N.J.A. Sloane's The Online Encyclopedia of Integer Sequences http://oeis.org/

Ron Knott's Fibonacci Numbers and the Golden Section http://www.maths.surrey.ac.uk/hosted-sites/R.Knott/Fibonacci/

#### Appendix

A survey of the arrays that are the theme of this paper is not complete without mentioning continued fractions, as a close correlation exists between the two. In what follows, certain periodic continued fractions ( $\theta_{\lambda}$ ) can be seen as a  $\varphi_{\lambda}$ -sequence turned inside out.

Recall that at the beginning of the paper, the formula

$$(c)F_n + (b)F_{n+1} = F_{n+2}$$
;  $F_0 = 0, F_1 = 1$ 

was the basis for a derivation of  $Q = x^2 - bx - c$ .

By substitution, Q expands to an infinite, cyclical continued fraction of period length  $\lambda_1$ . I.e., since  $x^2 - bx - c = 0$ , then  $x = b + \frac{c}{x}$ ; hence, continuous substitution of  $b + \frac{c}{x}$  for x on the right-hand

side yields the continued fraction  $x = b + \frac{c}{b + \frac{c$ 

To restore this fraction to a quadratic form, let the iterative algorithm

$$\frac{c}{x} + b \to x \tag{2.1}$$

represent its solution. Let I designate the number of iterations of (2.1). Then;

$$\lim_{t \to \infty} \left( \frac{c}{x} + b \to x \right) = \left( \frac{c}{x} + b = x \right)$$

and  $x^2 - bx - c = 0$ .

Thus for  $\lambda_1$  and *b*,*c* positive integers, both the infinite continued fraction and the ratio of adjacent terms of a  $\varphi_1$  series converge to the same number.

For  $\lambda > 1$ , we said that in the limit as  $n \to \infty$ ,  $\frac{F_{n+1}}{F_n} = x_1 = \frac{F_{n+\lambda+1}}{F_{n+\lambda}}$ . By use of the procedures that were then devised, such fractions reduce to  $\lambda$  quadratics. So, for  $S_1$  and  $S_2$  in  $\Phi_2$ , we find

$$Q_1 = b_1 x_1^2 - (b_1 b_2 - c_1 + c_2) x_1 - b_2 c_1 = 0$$
$$Q_2 = b_2 x_2^2 - (b_1 b_2 + c_1 - c_2) x_2 - b_1 c_2 = 0$$

Now for  $\lambda > 1$ ,  $\theta_{\lambda}$  can be 'factored' into  $\lambda$  fractions of the form:

$$\binom{c_j}{x_j} + b_j \to x_{j+1} \tag{2.2}$$

Thus for 
$$\lambda_2$$
, (2.2) has two expressions;  $\binom{c_1}{x_1} + b_1 \rightarrow x_2$  and  $\binom{c_2}{x_2} + b_2 \rightarrow x_1$ 

Then by substitution and taking limits:  $x_1 = b_2 + \frac{c_2}{b_1 + \frac{c_1}{x_1}}$ ;  $x_2 = b_1 + \frac{c_1}{b_2 + \frac{c_2}{x_2}}$ 

Clearing fractions and collecting terms:

$$Q_{1} = b_{1}x_{1}^{2} - (b_{1}b_{2} - c_{1} + c_{2})x_{1} - b_{2}c_{1} = 0$$
$$Q_{2} = b_{2}x_{2}^{2} - (b_{1}b_{2} + c_{1} - c_{2})x_{2} - b_{1}c_{2} = 0$$

These are the same equations as given by  $\Phi_2$ . Note that although an infinite periodic continued fraction  $\theta_{\lambda}$  is usually thought of as converging to one solution, in this context  $\theta_{\lambda}$  converges to  $\lambda$  solutions, and all are of equal interest. Henceforth, let  $\Theta_{\lambda}$  represent the set of these solutions and the associated quadratic equations. Now we have seen so far that for  $0 < \lambda \le 2$ ,  $\Theta_{\lambda}$  and  $\Phi_{\lambda}$  converge to the same values. For  $\lambda > 2$ , a slight complication appears. To see why, take the  $\varphi_3$  formula

$$Q_1 = (b_1b_2 + c_2)x_1^2 - (b_1b_2b_3 + b_1c_3 - b_2c_1 + b_3c_2)x_1 - (b_2b_3c_1 + c_1c_3) = 0$$

Having derived this equation from  $S_1$  in  $\Phi_3$   $[b_1, b_2, b_3][c_1, c_2, c_3]$ , we can expand it into a continued fraction  $\theta_3$ . First, arrange the terms of  $Q_1$  so that each is positive and then divide by  $x_1$  for

$$(b_1b_2 + c_2)x_1 + b_2c_1 = \frac{b_2b_3c_1 + c_1c_3}{x_1} + b_1b_2b_3 + b_1c_3 + b_3c_2$$

Some rearranging and tweaking of these terms allows for division by  $b_1 + \frac{c_1}{x_1}$ :

$$\left(\frac{b_1b_2x_1 + \frac{b_2c_1x_1}{x_1} + c_2x_1 = b_1b_2b_3 + \frac{b_2b_3c_1}{x_1} + b_1c_3 + \frac{c_1c_3}{x_1} + b_3c_2}{b_1 + \frac{c_1}{x_1}}\right) = \left(b_2x_1 + \frac{c_2x_1}{b_1 + \frac{c_1}{x_1}} = b_2b_3 + \frac{b_3c_2}{b_1 + \frac{c_1}{x_1}} + c_3\right)$$

Dividing the last equation by  $b_2 + \frac{c_2}{b_1 + \frac{c_1}{x_1}}$  yields  $x_1 = b_3 + \frac{c_3}{b_2 + \frac{c_2}{b_1 + \frac{c_1}{x_1}}}$ 

This last equation in turn can be expanded indefinitely. The aforementioned complication is one of a notational nature, for it appears that  $\Phi_3$  [ $b_1$ ,  $b_2$ ,  $b_3$ ][ $c_1$ ,  $c_2$ ,  $c_3$ ] is associated with the fraction set  $\Theta_3$  [ $b_3$ ,  $b_2$ ,  $b_1$ ][ $c_3$ ,  $c_2$ ,  $c_1$ ]. But since  $\theta_{\lambda}$  is solved 'upwards', the notation could be adjusted to say that  $\Phi_3$  [ $b_1$ ,  $b_2$ ,  $b_3$ ][ $c_1$ ,  $c_2$ ,  $c_3$ ]  $\leftrightarrow \Theta_3$  [ $b_1$ ,  $b_2$ ,  $b_3$ ][ $c_1$ ,  $c_2$ ,  $c_3$ ]

Hence, finding  $Q_1$  for any  $\lambda$  merely requires clearing  $x_1 = b_{\lambda} + \frac{c_{\lambda}}{b_{\lambda-1}} + \frac{c_{\lambda-1}}{b_{\lambda-2}} + \frac{c_{\lambda-2}}{\frac{c_{\lambda-2}}{b_{\lambda-2}}} + \frac{c_{\lambda-2}}{b_{\lambda-2}} +$ 

## WORKSHEET

## For page 15

|                        | $S_1$      | $S_2$       | $S_3$      | $S_4$      |
|------------------------|------------|-------------|------------|------------|
|                        | 1247.375   | 2808.34375  | 724.71875  | 2703.8125  |
| $F_{\text{-}4\lambda}$ | -731.71875 | -1149.84375 | -1045.3125 | -1045.3125 |
|                        | 299.59375  | 1658.5      | 404.125    | 613.1875   |
|                        | -432.125   | -641.1875   | -237.0625  | -251.0625  |
|                        | 167.0625   | 376.125     | 97.0625    | 362.125    |
| $F_{-3\lambda}$        | -98        | -154        | -140       | -140       |
|                        | 40.125     | 222.125     | 54.125     | 82.125     |
|                        | -57.875    | -85.875     | -31.75     | -33.625    |
|                        | 22.375     | 50.375      | 13         | 48.5       |
| $F_{-2\lambda}$        | -13.125    | -20.625     | -18.75     | -18.75     |
|                        | 5.375      | 29.75       | 7.25       | 11         |
|                        | -7.75      | -11.5       | -4.25      | -4.5       |
|                        | 3          | 6.75        | 1.75       | 6.5        |
| $F_{-\lambda}$         | -1.75      | -2.75       | -2.5       | -2.5       |
|                        | .75        | 4           | 1          | 1.5        |
|                        | -1         | -1.5        | 5          | 5          |
|                        | .5         | 1           | .5         | 1          |
| $\Phi_4[1,1,2,3][1,2]$ | 0          | 0           | 0          | 0          |
|                        | 1          | 1           | 1          | 1          |
|                        | 1          | 1           | 2          | 3          |
|                        | 3          | 3           | 8          | 4          |
| $F_{\lambda}$          | 7          | 11          | 10         | 10         |
|                        | 27         | 14          | 26         | 24         |
|                        | 34         | 36          | 62         | 92         |
|                        | 88         | 86          | 238        | 116        |
| $F_{2\lambda}$         | 210        | 330         | 300        | 300        |
|                        | 806        | 416         | 776        | 716        |
|                        | 1016       | 1076        | 1852       | 2748       |
|                        | 2628       | 2568        | 7108       | 3464       |
| $F_{3\lambda}$         | 6272       | 9856        | 8960       | 8960       |
|                        | 24072      | 12424       | 23176      | 21384      |
|                        | 30344      | 32136       | 55312      | 82072      |
|                        | 78488      | 76696       | 212288     | 103456     |
| $F_{4\lambda}$         | 187320     | 294360      | 267600     | 267600     |
|                        | 718936     | 371056      | 692176     | 638656     |
|                        |            |             |            |            |