

# PLSO: A generative framework for decomposing nonstationary time-series into piecewise stationary oscillatory components (Supplementary material)

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## APPENDIX

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## A. CONTINUOUS MODEL INTERPRETATION OF PLSO

We can establish the equivalent continuous model of the PLSO in Eq. 2, using stochastic different equation

$$\frac{d\tilde{\mathbf{z}}_j(t)}{dt} = \underbrace{\left( \left( -\frac{1}{l_j} \right) \oplus \begin{pmatrix} 0 & -\omega_j \\ \omega_j & 0 \end{pmatrix} \right)}_{\mathbf{F}} \tilde{\mathbf{z}}_j(t) + \varepsilon(t), \quad (\text{A.1})$$

where  $\tilde{\mathbf{z}}_j(t) : \mathbb{R} \rightarrow \mathbb{R}^2$ ,  $\oplus$  denotes the Kronecker sum and  $\varepsilon(t) \sim \mathcal{N}(0, \sigma_j^2 \mathbf{I}_{2 \times 2})$ . Discretizing the solution of Eq. A.1

at  $\Delta$ , such that  $\tilde{\mathbf{z}}_{j,k} = \tilde{\mathbf{z}}_j(k\Delta)$ , yields Eq. 2. Consequently, we obtain the following for  $\Delta > 0$

$$\begin{aligned} \exp(\mathbf{F}\Delta) &= \exp(-\Delta/l_j)\mathbf{R}(\omega_j), \\ \sigma_j^2 \int_0^\Delta \exp(\mathbf{F}(\Delta - \tau)) \exp(\mathbf{F}(\Delta - \tau))^\top d\tau \\ &= \sigma_j^2 (1 - \exp(-2\Delta/l_j)) \mathbf{I}_{2 \times 2}. \end{aligned}$$

This interpretation extends to the nonstationary PLSO. The corresponding continuous model for  $\tilde{\mathbf{z}}_{j,mN+n}$  in Eq. 3 is the same as Eq. A.1, with different variance  $\mathbb{E}[\varepsilon_j(t)\varepsilon_j^\top(t)] = \sum_{m=1}^M \sigma_{j,m}^2 \cdot \mathbf{1} \left( \left( \frac{m-1}{M} \right) T \leq t < \left( \frac{m}{M} \right) T \right) \mathbf{I}_{2 \times 2}$ .

## B. PSD FOR COMPLEX AR(1) PROCESS

We compute the steady-state covariance denoted as  $\mathbf{P}_\infty^j$ . Since we assume  $\mathbf{P}_1^j = \sigma_j^2 \mathbf{I}_{2 \times 2}$ , it is easy to show that  $\mathbf{P}_k^j$  is a diagonal matrix from  $\mathbf{R}(\omega_j)\mathbf{R}^\top(\omega_j) = \mathbf{I}_{2 \times 2}$ . Denoting  $\mathbf{P}_\infty^j = \alpha \mathbf{I}_{2 \times 2}$ , we use the discrete Lyapunov equation

$$\begin{aligned} \mathbf{P}_\infty^j &= \exp(-2\Delta/l_j)\mathbf{R}(\omega_j)\mathbf{P}_\infty^j\mathbf{R}^\top(\omega_j) \\ &\quad + \sigma_j^2 (1 - \exp(-2\Delta/l_j)) \mathbf{I}_{2 \times 2} \\ \Rightarrow \alpha &= \exp(-2\Delta/l_j)\alpha + \sigma_j^2 (1 - \exp(-2\Delta/l_j)) \\ \Rightarrow \mathbf{P}_\infty^j &= \sigma_j^2 \mathbf{I}_{2 \times 2}, \end{aligned}$$

which implies that under the assumption  $\mathbf{P}_1^j = \sigma_j^2 \mathbf{I}_{2 \times 2}$ , we are guaranteed  $\mathbf{P}_k^j = \sigma_j^2 \mathbf{I}_{2 \times 2}, \forall k$ . To compute the PSD of the stationary process  $\mathbf{z}_j$ , we now need to compute the auto-covariance. Since only  $\mathbf{z}_{j,k}^{\Re}$  contributes to  $\mathbf{y}_k$ , we compute the autocovariance of  $\mathbb{E}[\mathbf{z}_{j,k}^{\Re} \mathbf{z}_{j,k+n}^{\Re}]$  as

$$\begin{aligned} \mathbb{E}[\mathbf{z}_{j,k}^{\Re} \mathbf{z}_{j,k+n}^{\Re}] &= \mathbb{E}[\mathbf{z}_{j,k}^{\Re} \cdot \Re(\rho_j^n \exp(i\omega_j n) \mathbf{z}_{j,k})] \\ &= \rho_j^n \mathbb{E}[\mathbf{z}_{j,k}^{\Re} \mathbf{z}_{j,k}^{\Re} \cos \omega_j n] \\ &= \rho_j^n \cos \omega_j n \cdot \mathbb{E}[\{\mathbf{z}_{j,k}^{\Re}\}^2] \\ &= \rho_j^n \sigma_j^2 \cos \omega_j n, \end{aligned}$$

where  $\Re(\cdot)$  denotes the operator that extracts the real part of the complex argument and we used the fact that  $\mathbb{E}[\mathbf{z}_{j,k}^{\Re} \mathbf{z}_{j,k}^{\Im}] = 0$ . The spectra for the  $j^{\text{th}}$  component,  $S_j(\omega)$  can be written as

$$\begin{aligned} S_j(\omega) &= \sum_{n=-\infty}^{\infty} \mathbb{E} [\mathbf{z}_{j,k}^{\Re} \mathbf{z}_{j,k+n}^{\Re}] \exp(-i\omega n) \\ &= \sum_{n=-\infty}^{\infty} \rho_j^n \sigma_j^2 \cos \omega_j n \exp(-i\omega n) \\ &= \sigma_j^2 \sum_{n=-\infty}^{\infty} \rho_j^n \{\exp(i\omega_j n) + \exp(-i\omega_j n)\} \exp(-i\omega n) \\ &= \sigma_j^2 \sum_{n=-\infty}^{\infty} \rho_j^n \exp(-i(\omega \pm \omega_j)n). \end{aligned}$$

Unpacking the infinite sum for one of the terms,

$$\begin{aligned} &\sum_{n=-\infty}^{\infty} \rho_j^n \exp(-i(\omega - \omega_j)n) \\ &= 1 + \sum_{n=1}^{\infty} \rho_j^n \exp(-i(\omega - \omega_j)n) + \rho_j^n \exp(i(\omega - \omega_j)n) \\ &= 1 + \frac{\rho_j \exp(-i(\omega - \omega_j))}{1 - \rho_j \exp(-i(\omega - \omega_j))} + \frac{\rho_j \exp(i(\omega - \omega_j))}{1 - \rho_j \exp(i(\omega - \omega_j))} \\ &= 1 + \frac{2\rho_j \cos(\omega - \omega_j) - 2\rho_j^2}{(1 - \rho_j \exp(-i(\omega - \omega_j))) (1 - \rho_j \exp(i(\omega - \omega_j)))} \\ &= \frac{1 - \rho_j^2}{1 + \rho_j^2 - 2\rho_j \cos(\omega - \omega_j)}. \end{aligned}$$

Using the relation  $\rho_j = \exp(-\Delta/l_j)$  and unpacking the infinite sum for the other term, we have

$$\begin{aligned} S_j(\omega) &= \frac{\sigma_j^2(1 - \exp(-2\Delta/l_j))}{1 + \exp(-2\Delta/l_j) - 2\exp(-\Delta/l_j) \cos(\omega - \omega_j)} \\ &\quad + \frac{\sigma_j^2(1 - \exp(-2\Delta/l_j))}{1 + \exp(-2\Delta/l_j) - 2\exp(-\Delta/l_j) \cos(\omega + \omega_j)}. \end{aligned}$$

Since Fourier transform is a linear operator, we can conclude that  $\gamma(\omega) = \sigma_v^2 + \sum_{j=1}^J S_j(\omega)$ .

### C. PROOF FOR PROPOSITION 1 (SECTION 3.2.1)

**Proposition 1.** For a given  $m$ , as  $\Delta \rightarrow 0$ , the samples on either side of the interval boundary, which are  $\tilde{\mathbf{z}}_{j,(m+1)N}$  and  $\tilde{\mathbf{z}}_{j,(m+1)N+1}$ , converge to each other in mean square,

$$\lim_{\Delta \rightarrow 0} \mathbb{E}[\Delta \tilde{\mathbf{z}}_{j,(m+1)N} \Delta \tilde{\mathbf{z}}_{j,(m+1)N}^{\text{T}}] = 0,$$

where we use  $\Delta \tilde{\mathbf{z}}_{j,(m+1)N} = \tilde{\mathbf{z}}_{j,(m+1)N+1} - \tilde{\mathbf{z}}_{j,(m+1)N}$ .

*Proof.* To analyze Eq. 3 in the limit of  $\Delta \rightarrow 0$ , we use the equivalent continuous model. It suffices to show that  $\lim_{\Delta \rightarrow 0} \exp(\mathbf{F}\Delta) = \mathbf{I}_{2 \times 2}$  and  $\lim_{\Delta \rightarrow 0} \mathbb{E}[\varepsilon_{j,(m+1)N+1} \varepsilon_{j,(m+1)N+1}^{\text{T}}] = \mathbf{0}$ . We have,

$$\begin{aligned} \lim_{\Delta \rightarrow 0} \exp(\mathbf{F}\Delta) &= \mathbf{I}_{2 \times 2} + \lim_{\Delta \rightarrow 0} \sum_{k=1}^{\infty} \frac{\Delta^k}{k!} \mathbf{F}^k = \mathbf{I}_{2 \times 2} \\ \lim_{\Delta \rightarrow 0} \mathbb{E}[\varepsilon_{j,(m+1)N+1} \varepsilon_{j,(m+1)N+1}^{\text{T}}] / \sigma_{j,m+1}^2 &= \lim_{\Delta \rightarrow 0} \int_0^{\Delta} \exp(\mathbf{F}(\Delta - \tau)) \exp(\mathbf{F}(\Delta - \tau))^{\text{T}} d\tau = \mathbf{0}. \end{aligned}$$

Since this implies  $\lim_{\Delta \rightarrow 0} \mathbb{E}[\Delta \tilde{\mathbf{z}}_{j,(m+1)N} \Delta \tilde{\mathbf{z}}_{j,(m+1)N}^{\text{T}}] = 0$ , we have convergence in mean square.  $\square$

### D. PROOF FOR PROPOSITION 2 (SECTION 3.2.2)

**Proposition 2.** Assume  $l_j \ll N\Delta$ , such that  $\mathbf{P}_{m,N}^j = \mathbf{P}_{m,\infty}^j$ . In Eq. 3, the difference between  $\mathbf{P}_{m,\infty}^j = \sigma_{j,m}^2 \mathbf{I}_{2 \times 2}$  and  $\mathbf{P}_{m+1,\infty}^j = \sigma_{j,m+1}^2 \mathbf{I}_{2 \times 2}$  decays exponentially fast as a function of  $n$ , for  $1 \leq n \leq N$ ,

$$\mathbf{P}_{m+1,n}^j = \mathbf{P}_{m+1,\infty}^j + \exp(-2n\Delta/l_j) (\mathbf{P}_{m,\infty}^j - \mathbf{P}_{m+1,\infty}^j).$$

*Proof.* We first obtain the steady-state covariance  $\mathbf{P}_{m,\infty}^j$ , similar to **Appendix B**. Since we assume  $\mathbf{P}_{1,1}^j = \sigma_{j,1}^2 \mathbf{I}_{2 \times 2}$ , we can show that  $\forall m, n$ ,  $\mathbf{P}_{m,n}^j$  is a diagonal matrix, noting that  $\mathbf{R}(\omega_j) \mathbf{R}^{\text{T}}(\omega_j) = \mathbf{I}_{2 \times 2}$ . Denoting  $\mathbf{P}_{m,\infty}^j = \alpha \mathbf{I}_{2 \times 2}$ , we now use the discrete Lyapunov equation

$$\begin{aligned} \mathbf{P}_{m,\infty}^j &= \exp(-2\Delta/l_j) \mathbf{R}(\omega_j) \mathbf{P}_{m,\infty}^j \mathbf{R}^{\text{T}}(\omega_j) \\ &\quad + \sigma_{j,m}^2 (1 - \exp(-2\Delta/l_j)) \mathbf{I}_{2 \times 2} \\ &\Rightarrow \alpha = \exp(-2\Delta/l_j) \alpha + \sigma_{j,m}^2 (1 - \exp(-2\Delta/l_j)) \\ &\Rightarrow \mathbf{P}_{m,\infty}^j = \sigma_{j,m}^2 \mathbf{I}_{2 \times 2}. \end{aligned}$$

We now prove the proposition by induction. For fixed  $j$  and  $m$ , and for  $n = 1$ ,

$$\begin{aligned} \mathbf{P}_{m+1,1}^j &= \exp(-2\Delta/l_j) \mathbf{R}(\omega_j) \mathbf{P}_{m,N}^j \mathbf{R}^{\text{T}}(\omega_j) \\ &\quad + \sigma_{j,m+1}^2 (1 - \exp(-2\Delta/l_j)) \mathbf{I}_{2 \times 2} \\ &= \{\sigma_{j,m+1}^2 + \exp(-2\Delta/l_j) (\sigma_{j,m}^2 - \sigma_{j,m+1}^2)\} \mathbf{I}_{2 \times 2}. \end{aligned}$$

Assuming the same holds for  $n = n' - 1$ , we have for

$$n = n',$$

$$\begin{aligned} \mathbf{P}_{m+1, n'}^j &= \exp(-2\Delta/l_j) \mathbf{R}(\omega_j) \mathbf{P}_{m, n'-1}^j \mathbf{R}^T(\omega_j) \\ &\quad + \sigma_{j, m+1}^2 (1 - \exp(-2\Delta/l_j)) \mathbf{I}_{2 \times 2} \\ &= \exp(-2\Delta/l_j) \sigma_{j, m+1}^2 \mathbf{I}_{2 \times 2} \\ &\quad + \exp(-2n'\Delta/l_j) (\sigma_{j, m}^2 - \sigma_{j, m+1}^2) \mathbf{I}_{2 \times 2} \\ &\quad + \sigma_{j, m+1}^2 (1 - \exp(-2\Delta/l_j)) \mathbf{I}_{2 \times 2} \\ &= \{\sigma_{j, m+1}^2 + \exp(-2n'\Delta/l_j) (\sigma_{j, m}^2 - \sigma_{j, m+1}^2)\} \mathbf{I}_{2 \times 2}. \end{aligned}$$

By the principle of induction, Eq. holds for  $1 \leq n \leq N$ .  $\square$

## E. INITIALIZATION & ESTIMATION FOR PLSO (SECTION 4)

### E.1 INITIALIZATION

As noted in the main text, we use AIC to determine the optimal number of  $J$ . For a given number of components  $J$ , we first construct the spectrogram of the data using STFT and identify the frequency bands with prominent power, i.e., frequency bands whose average power exceeds pre-determined threshold. The center frequencies of these bands serve as the initial center frequencies  $\{\omega_j^{\text{init}}\}_j$ , which are either fixed throughout the algorithm or further refined through the estimation algorithm in the main text. If  $J$  exceeds the number of identified frequency bands from the spectrogram, 1) we first place  $\{\omega_j^{\text{init}}\}_j$  in the prominent frequency bands and 2) we then place the remaining components uniformly spread out in  $[0, \omega_c]$ , where  $\omega_c$  is a cutoff frequency to be further determined in the next section. As for  $\{l_j^{\text{init}}\}_j$ , we set it to be a certain fraction of the corresponding  $\{\omega_j^{\text{init}}\}_j$ . We then fit  $\{\sigma_{j, m}^2\}_{j, m}$  and  $\theta$  with  $\lambda = 0$ , through the procedure explained in Stage 1. We finally use these estimates as initial values for other values of  $\lambda$ .

### E.2 ESTIMATION FOR $\sigma_\nu^2$

There are two possible ways to estimate the observation noise variance  $\sigma_\nu^2$ . The first approach is to perform maximum likelihood estimation of  $f(\{\sigma_{j, m}^2\}_{j, m}; \theta)$  with respect to  $\sigma_\nu^2$ . The second approach, which we found to work *better* in practice and use throughout the manuscript, is to directly estimate it from the Fourier transform of the data. Given a cutoff frequency  $\omega_c$ , informed by domain knowledge, we take the average power of the Fourier transform of  $\mathbf{y}$  in  $[\omega_c, f_s/2]$ . For instance, it is widely known that the spectral content below 40 Hz in anesthesia EEG dataset is physiologically relevant and we use  $\omega_c \simeq 40$  Hz.

## F. OPTIMIZATION FOR $\{\sigma_{j, m}^2\}_{j, m}$ VIA PROXIMAL GRADIENT UPDATE (SECTION 4.1.1)

We discuss the algorithm to obtain a local optimal solution of  $\{\sigma_{j, m}^2\}_{j, m}$  to the MAP optimization problem in Eq. 8. We define  $\psi_{j, m} = \log \sigma_{j, m}^2$  and  $\Psi = [\psi_{1, 1}, \dots, \psi_{1, M}, \dots, \psi_{J, M}] \in \mathbb{R}^{JM}$  for notational decluttering, to be used in this section.

We rewrite Eq. 8 as

$$\begin{aligned} &\min_{\Psi} \underbrace{-\log p(\Psi | \mathbf{y}, \theta)}_{h(\Psi; \theta)} \\ &= \min_{\Psi} \underbrace{\frac{1}{2} \sum_{m=1}^M \sum_{n=1}^N \left\{ \log \gamma^{(m)}(\omega_n) + \frac{I^{(m)}(\omega_n)}{\gamma^{(m)}(\omega_n)} \right\}}_{-f(\Psi; \theta)} \\ &\quad + \underbrace{\frac{\lambda}{2} \sum_{j=1}^J \sum_{m=1}^M (\psi_{j, m} - \psi_{j, m-1})^2}_{-g(\Psi; \theta)} \\ &= \min_{\Psi} -f(\Psi; \theta) - g(\Psi; \theta). \end{aligned}$$

The algorithm is described in Algorithm 1. It follows the steps outlined in the inexact accelerated proximal gradient algorithm [Li and Lin, 2015]. For faster convergence, we use larger step sizes with the Barzilai-Borwein (BB) step size initialization rule [Barzilai and Borwein, 1988]. For rest of this section, we drop dependence on  $\theta$ . The main novelty of our algorithm is the proximal gradient update

$$\begin{aligned} &\mathbf{u}^{(l+1)} \\ &= \text{prox}_{-\alpha_{\mathbf{w}}^{(l)} g}(\mathbf{w}^{(l)} + \alpha_{\mathbf{w}}^{(l)} \nabla f(\mathbf{w}^{(l)})) \\ &= \arg \min_{\Psi} \frac{1}{2\alpha_{\mathbf{w}}^{(l)}} \|\Psi - (\mathbf{w}^{(l)} + \alpha_{\mathbf{w}}^{(l)} \nabla f(\mathbf{w}^{(l)}))\|^2 - g(\Psi) \\ &= \arg \min_{\Psi} \sum_{j=1}^J \sum_{m=1}^M \frac{\left( \left( \mathbf{w}_{j, m}^{(l)} + \alpha_{\mathbf{w}}^{(l)} \cdot \frac{\partial f(\mathbf{w}^{(l)})}{\partial \mathbf{w}_{j, m}} \right) - \psi_{j, m} \right)^2}{2\alpha_{\mathbf{w}}^{(l)}} \\ &\quad + \frac{\lambda}{2} (\psi_{j, m} - \psi_{j, m-1})^2, \end{aligned}$$

where the same holds for  $\mathbf{x}^{(l+1)} = \text{prox}_{-\alpha_{\Psi}^{(l)} g}(\Psi^{(l)} + \alpha_{\Psi}^{(l)} \nabla f(\Psi^{(l)}))$ . The auxiliary variables  $\mathbf{w}, \mathbf{u}, \mathbf{x} \in \mathbb{R}^{JM}$  ensure convergence of  $\Psi$ . We use  $\mathbf{w}_{j, m}^{(l)}$  to denote  $((m-1)J + j)^{\text{th}}$  entry of  $\mathbf{w}^{(l)}$ . As mentioned in the main text, this can be solved in a computationally efficient manner by using Kalman filter/smoothing.

## G. LIPSCHITZ CONSTANT FOR $\nabla f(\{\sigma_{j,m}^2\}_{j,m}; \theta)$ (SECTION 4.1.1)

In this section, we prove that under some assumptions, we can show that the log-likelihood  $f(\{\sigma_{j,m}^2\}_{j,m}; \theta)$  has Lipschitz continuous gradient with the Lipschitz constant  $C$ . As in the previous section, we use  $\psi_{j,m} = \log \sigma_{j,m}^2$  and  $\Psi = [\psi_{1,1}, \dots, \psi_{1,M}, \dots, \psi_{J,M}] \in \mathbb{R}^{JM}$ .

Let us start by restating the definition of Lipschitz continuous gradient.

**Definition** A continuously differentiable function  $f : \mathcal{S} \rightarrow \mathbb{R}$  is Lipschitz continuous gradient if

$$\|\nabla f(\mathbf{x}) - \nabla f(\mathbf{y})\|_2 \leq C \|\mathbf{x} - \mathbf{y}\|_2 \quad \text{for every } \mathbf{x}, \mathbf{y} \in \mathcal{S},$$

where  $\mathcal{S}$  is a compact subset of  $\mathbb{R}^{JM}$  and  $C > 0$  is the Lipschitz constant.

Our goal is to find the constant  $C$  for the Whittle likelihood  $f(\Psi; \theta)$

$$\begin{aligned} f(\Psi; \theta) &= -\frac{1}{2} \sum_{m=1}^M \sum_{n=1}^N \left\{ \log \gamma_{m,n} + \frac{I_{m,n}}{\gamma_{m,n}} \right\} \\ &= -\frac{1}{2} \sum_{m=1}^M \sum_{n=1}^N \log \left( \underbrace{\sigma_\nu^2 + \sum_{j=1}^J \exp(\psi_{j,m}) \alpha_{j,n}}_{f_1(\Psi; \theta)} \right) \\ &\quad - \frac{1}{2} \sum_{m=1}^M \sum_{n=1}^N \frac{I_{m,n}}{\underbrace{\left( \sigma_\nu^2 + \sum_{j=1}^J \exp(\psi_{j,m}) \alpha_{j,n} \right)}_{f_2(\Psi; \theta)}}, \end{aligned}$$

where we use  $\gamma_{m,n} = \gamma^{(m)}(\omega_n)$  and  $I_{m,n} = I^{(m)}(\omega_n)$  for notational simplicity and

$$\alpha_{j,n} = \frac{(1 - \exp(-2\Delta/l_j))}{1 + \exp(-2\Delta/l_j) - 2 \exp(-\Delta/l_j) \cos(\omega_n - \omega_j)}.$$

We make the following assumptions

1.  $I_{m,n}$  is bounded, i.e.,  $I_{m,n} \leq C_I$ . With the real-world signal, we can reasonably assume that  $I_{m,n}$  or energy of the signal is bounded.
2.  $\forall j, m, \psi_{j,m}$  is bounded, i.e.,  $|\psi_{j,m}| \leq \log C_\psi$  for some  $C_\psi > 1$ . This implies  $1/C_\psi \leq \exp(\psi_{j,m}) \leq C_\psi$ .

In addition, we have the following facts

1.  $I_{m,n}$ ,  $\alpha_{j,n}$ , and  $\gamma_{m,n}$  are nonnegative.
2.  $I_{m,n}$  and  $\gamma_{m,n}$  are bounded. This follows from the aforementioned assumptions.
3. For given  $\{l_j\}_j$ , we have bounded  $\alpha_{j,n}$ . To see this,

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### Algorithm 1: Inference for $\Psi$ via inexact APG

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**Result:**  $\hat{\Psi}$

**Initialize**  $\Psi^{(0)} = \Psi^{(1)} = \mathbf{u}^{(1)}$ ,  
 $\beta^{(0)} = 0, \beta^{(1)} = 1, \delta > 0, \rho < 1$

**for**  $l \leftarrow 1$  **to**  $L$  **do**

$$\mathbf{w}^{(l)} = \Psi^{(l)} + \frac{\beta^{(l-1)}}{\beta^{(l)}} (\mathbf{u}^{(l)} - \Psi^{(l-1)}) + \frac{\beta^{(l-1)} - 1}{\beta^{(l)}} (\Psi^{(l)} - \Psi^{(l-1)})$$

**(BB step size initialization rule)**

$$\mathbf{s}^{(l)} = \mathbf{u}^{(l)} - \mathbf{w}^{(l-1)},$$

$$\mathbf{r}^{(l)} = -\nabla f(\mathbf{u}^{(l)}) + \nabla f(\mathbf{w}^{(l-1)})$$

$$\alpha_{\mathbf{w}}^{(l)} = ((\mathbf{s}^{(l)})^T \mathbf{s}^{(l)}) / ((\mathbf{s}^{(l)})^T \mathbf{r}^{(l)})$$

$$\mathbf{s}^{(l)} = \mathbf{x}^{(l)} - \Psi^{(l-1)},$$

$$\mathbf{r}^{(l)} = -\nabla f(\mathbf{x}^{(l)}) + \nabla f(\Psi^{(l-1)})$$

$$\alpha_{\Psi}^{(l)} = ((\mathbf{s}^{(l)})^T \mathbf{s}^{(l)}) / ((\mathbf{s}^{(l)})^T \mathbf{r}^{(l)})$$

**(Proximal update step)**

**repeat**

$$\mathbf{u}^{(l+1)} = \text{prox}_{-\alpha_{\mathbf{w}}^{(l)} g}(\mathbf{w}^{(l)} + \alpha_{\mathbf{w}}^{(l)} \nabla f(\mathbf{w}^{(l)}))$$

$$\alpha_{\mathbf{w}}^{(l)} = \rho \cdot \alpha_{\mathbf{w}}^{(l)}$$

**until**  $h(\mathbf{u}^{(l+1)}) \leq h(\mathbf{w}^{(l)}) - \delta \|\mathbf{u}^{(l+1)} - \mathbf{w}^{(l)}\|^2$ ;

**repeat**

$$\mathbf{x}^{(l+1)} = \text{prox}_{-\alpha_{\Psi}^{(l)} g}(\Psi^{(l)} + \alpha_{\Psi}^{(l)} \nabla f(\Psi^{(l)}))$$

$$\alpha_{\Psi}^{(l)} = \rho \cdot \alpha_{\Psi}^{(l)}$$

**until**  $h(\mathbf{x}^{(l+1)}) \leq h(\Psi^{(l)}) - \delta \|\mathbf{x}^{(l+1)} - \Psi^{(l)}\|^2$ ;

$$\beta^{(l+1)} = \frac{1 + \sqrt{4(\beta^{(l)})^2 + 1}}{2}$$

$$\Psi^{(l+1)} = \begin{cases} \mathbf{u}^{(l+1)} & \text{if } h(\mathbf{u}^{(l+1)}) \leq h(\mathbf{x}^{(l+1)}) \\ \mathbf{x}^{(l+1)} & \text{otherwise} \end{cases}$$

**end**

$$\hat{\Psi} = \Psi^{(L)}$$


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note that the maximum of  $\alpha_{j,n}$  is achieved at  $\omega_n = \omega_j$ ,

$$\begin{aligned}\max \alpha_{j,n} &= \frac{(1 - \exp(-2\Delta/l_j))}{1 + \exp(-2\Delta/l_j) - 2 \exp(-\Delta/l_j)} \\ &= \frac{(1 + \exp(-\Delta/l_j))}{(1 - \exp(-\Delta/l_j))}.\end{aligned}$$

Therefore, denoting  $l_{\max} = \max_j \{l_j\}$ ,

$$\max \alpha_{j,n} \leq \frac{(1 + \exp(-\Delta/l_{\max}))}{(1 - \exp(-\Delta/l_{\max}))} = C_\alpha.$$

Finally, we define  $\mathcal{S} = [-\log C_\psi, \log C_\psi] \subset \mathbb{R}^{JM}$ .

We want to compute the Lipschitz constant for  $\nabla f_1(\Psi)$  and  $\nabla f_2(\Psi)$  for  $\Psi, \bar{\Psi} \in \mathcal{S}$ , i.e.,

$$\begin{cases} \|\nabla f_1(\Psi) - \nabla f_1(\bar{\Psi})\|_2 \leq C_1 \|\Psi - \bar{\Psi}\|_2 \\ \|\nabla f_2(\Psi) - \nabla f_2(\bar{\Psi})\|_2 \leq C_2 \|\Psi - \bar{\Psi}\|_2, \end{cases}$$

where we dropped dependence on  $\theta$  for notational simplicity. Consequently, the triangle inequality yields

$$\begin{aligned}\|\nabla f(\Psi) - \nabla f(\bar{\Psi})\|_2 &\leq \|\nabla f_1(\Psi) - \nabla f_1(\bar{\Psi})\|_2 + \|\nabla f_2(\Psi) - \nabla f_2(\bar{\Psi})\|_2 \\ &\leq (C_1 + C_2) \|\Psi - \bar{\Psi}\|_2.\end{aligned}$$

### LIPSCHITZ CONSTANT $C_1$ FOR $\nabla f_1(\Psi)$

Let us examine  $f_1(\Psi)$  first. The derivative with respect to  $\psi_{j,m}$  is given as

$$\begin{aligned}\frac{\partial f_1(\Psi)}{\partial \psi_{j,m}} &= \sum_{n=1}^N \alpha_{j,n} \cdot \underbrace{\frac{\exp(\psi_{j,m})}{\sigma_\nu^2 + \sum_{j'=1}^J \exp(\psi_{j',m}) \alpha_{j',n}}}_{\tilde{f}(\psi_{j,m})} \\ &= \sum_{n=1}^N \alpha_{j,n} \tilde{f}(\psi_{j,m}).\end{aligned}$$

We now have

$$\left| \frac{\partial f_1(\Psi)}{\partial \psi_{j,m}} - \frac{\partial f_1(\bar{\Psi})}{\partial \psi_{j,m}} \right| = \sum_{n=1}^N |\alpha_{j,n}| \cdot \left| \tilde{f}(\psi_{j,m}) - \tilde{f}(\bar{\psi}_{j,m}) \right|.$$

Without loss of generality, we assume  $\psi_{j,m} \geq \bar{\psi}_{j,m}$ . We now apply the mean value theorem (MVT) to  $\tilde{f}(\psi_{j,m})$

$$\tilde{f}(\psi_{j,m}) - \tilde{f}(\bar{\psi}_{j,m}) = \tilde{f}'(\psi'_{j,m})(\psi_{j,m} - \bar{\psi}_{j,m}),$$

where  $\psi'_{j,m} \in [\bar{\psi}_{j,m}, \psi_{j,m}]$ . We can compute and bound  $\tilde{f}'(\psi'_{j,m}) = d\tilde{f}(\psi'_{j,m})/d\psi'_{j,m}$  as follows

$$\begin{aligned}\tilde{f}'(\psi'_{j,m}) &= \frac{\exp(\psi'_{j,m})}{\sigma_\nu^2 + \sum_{j'=1}^J \exp(\psi'_{j',m}) \alpha_{j',n}} \\ &\quad - \frac{\alpha_{j,n} \exp^2(\psi'_{j,m})}{(\sigma_\nu^2 + \sum_{j'=1}^J \exp(\psi'_{j',m}) \alpha_{j',n})^2} \\ &= \frac{\exp(\psi'_{j,m})}{\sigma_\nu^2 + \sum_{j'=1}^J \exp(\psi'_{j',m}) \alpha_{j',n}} \\ &\quad \times \underbrace{\left( 1 - \frac{\alpha_{j,n} \exp(\psi'_{j,m})}{\sigma_\nu^2 + \sum_{j'=1}^J \exp(\psi'_{j',m}) \alpha_{j',n}} \right)}_{\leq 1} \\ &\leq \frac{\exp(\psi'_{j,m})}{\sigma_\nu^2 + \sum_{j'=1}^J \exp(\psi'_{j',m}) \alpha_{j',n}} \\ &\leq \frac{C_\psi}{\sigma_\nu^2}.\end{aligned}$$

Combining both, we have

$$\begin{aligned}\sum_{n=1}^N |\alpha_{j,n}| \cdot \left| \tilde{f}(\psi_{j,m}) - \tilde{f}(\bar{\psi}_{j,m}) \right| &= \sum_{n=1}^N |\alpha_{j,n}| \cdot \left| \tilde{f}'(\psi'_{j,m})(\psi_{j,m} - \bar{\psi}_{j,m}) \right| \\ &\leq \frac{NC_\alpha C_\psi}{\sigma_\nu^2} |\psi_{j,m} - \bar{\psi}_{j,m}|.\end{aligned}$$

We thus have,

$$\begin{aligned}\|\nabla f_1(\Psi) - \nabla f_1(\bar{\Psi})\|_2^2 &= \sum_{j=1}^J \sum_{m=1}^M \left( \frac{\partial f_1(\Psi)}{\partial \psi_{j,m}} - \frac{\partial f_1(\bar{\Psi})}{\partial \psi_{j,m}} \right)^2 \\ &\leq \left( \frac{JMN C_\alpha C_\psi}{\sigma_\nu^2} \right)^2 \|\Psi - \bar{\Psi}\|_2^2.\end{aligned}$$

### LIPSCHITZ CONSTANT $C_2$ FOR $\nabla f_2(\Psi)$

Computing  $C_2$  proceeds in a similar manner to computing  $C_1$ . The derivative with respect to  $\psi_{j,m}$  is given as

$$\begin{aligned}\frac{\partial f_2(\Psi)}{\partial \psi_{j,m}} &= - \sum_{n=1}^N I_{m,n} \alpha_{j,n} \cdot \underbrace{\frac{\exp(\psi_{j,m})}{(\sigma_\nu^2 + \sum_{j'=1}^J \exp(\psi_{j',m}) \alpha_{j',n})^2}}_{\tilde{f}(\psi_{j,m})}.\end{aligned}$$

We now have

$$\begin{aligned}\left| \frac{\partial f_2(\Psi)}{\partial \psi_{j,m}} - \frac{\partial f_2(\bar{\Psi})}{\partial \psi_{j,m}} \right| &= \sum_{n=1}^N |I_{m,n} \alpha_{j,n}| \cdot \left| -\tilde{f}(\psi_{j,m}) + \tilde{f}(\bar{\psi}_{j,m}) \right|.\end{aligned}$$

Without loss of generality, assume  $\psi_{j,m} \geq \bar{\psi}_{j,m}$ . To apply MVT, we need to compute and bound  $\tilde{f}'(\psi'_{j,m})$

$$\begin{aligned}\tilde{f}'(\psi'_{j,m}) &= \frac{\exp(\psi'_{j,m})}{(\sigma_\nu^2 + \sum_{j'=1}^J \exp(\psi'_{j',m}) \alpha_{j',n})^2} \\ &\quad - \frac{2\alpha_{j,n} \exp(\psi'_{j,m})}{(\sigma_\nu^2 + \sum_{j'=1}^J \exp(\psi'_{j',m}) \alpha_{j',n})^3} \\ &= \frac{\exp(\psi'_{j,m})}{(\sigma_\nu^2 + \sum_{j'=1}^J \exp(\psi'_{j',m}) \alpha_{j',n})^2} \\ &\quad \times \underbrace{\left(1 - \frac{2\alpha_{j,n} \exp(\psi'_{j,m})}{\sigma_\nu^2 + \sum_{j'=1}^J \exp(\psi'_{j',m}) \alpha_{j',n}}\right)}_{\leq 1} \\ &\leq \frac{\exp(\psi'_{j,m})}{(\sigma_\nu^2 + \sum_{j'=1}^J \exp(\psi'_{j',m}) \alpha_{j',n})^2} \\ &\leq \frac{C_\psi}{\sigma_\nu^4}.\end{aligned}$$

Applying MVT,

$$\begin{aligned}&\sum_{n=1}^N |I_{m,n} \alpha_{j,n}| \cdot \left| -\tilde{f}(\psi_{j,m}) + \tilde{f}(\bar{\psi}_{j,m}) \right| \\ &= \sum_{n=1}^N |I_{m,n} \alpha_{j,n}| \cdot \left| \tilde{f}'(\psi'_{j,m})(\psi_{j,m} - \bar{\psi}_{j,m}) \right| \\ &\leq \frac{NC_I C_\alpha C_\psi}{\sigma_\nu^4} |\psi_{j,m} - \bar{\psi}_{j,m}|.\end{aligned}$$

We thus have,

$$\begin{aligned}&\|\nabla f_2(\Psi) - \nabla f_2(\bar{\Psi})\|_2^2 \\ &= \sum_{j=1}^J \sum_{m=1}^M \left( \frac{\partial f_2(\Psi)}{\partial \psi_{j,m}} - \frac{\partial f_2(\bar{\Psi})}{\partial \psi_{j,m}} \right)^2 \\ &\leq \left( \frac{JMNC_\alpha C_\psi C_I}{\sigma_\nu^4} \right)^2 \|\Psi - \bar{\Psi}\|_2^2.\end{aligned}$$

Collecting the Lipschitz constants  $C_1$  and  $C_2$ , we finally have

$$\|\nabla f(\Psi) - \nabla f(\bar{\Psi})\|_2 \leq \underbrace{\frac{JMNC_\alpha C_\psi}{\sigma_\nu^2}}_C \left(1 + \frac{C_I}{\sigma_\nu^2}\right) \|\Psi - \bar{\Psi}\|_2.$$

## H. INFERENCE WITH

$p(\{\mathbf{z}_j\}_j \mid \{\sigma_{j,m}^2\}_{j,m}, \mathbf{y}, \theta)$  (SECTION 4.2)

We present the details for performing inference with  $p(\{\mathbf{z}_j\}_j \mid \{\sigma_{j,m}^2\}_{j,m}, \mathbf{y}, \theta)$ , given the estimates  $\{\hat{\sigma}_{j,m}^2\}_{j,m}$  and  $\hat{\theta}$  from window-level inference. First, we present the

Kalman filter/smoothing algorithm to compute the mean posterior trajectory, and the credible interval. Next, we present the forward filtering backward sampling (FFBS) algorithm Carter and Kohn [1994], Lindsten and Schön [2013] to generate Monte Carlo (MC) sample trajectories.

First, we define additional notations.

$$1) \tilde{\mathbf{z}}_{j,k|k'} = \mathbb{E}[\tilde{\mathbf{z}}_{j,k} \mid \{\hat{\sigma}_{j,m}^2\}_{j,m}, \mathbf{y}_{1:k'}, \hat{\theta}] \in \mathbb{R}^2$$

Posterior mean of  $\tilde{\mathbf{z}}_{j,k}$ . We are primarily concerned with the following three types: 1)  $\tilde{\mathbf{z}}_{j,k|k-1}$ , the one-step prediction estimate, 2)  $\tilde{\mathbf{z}}_{j,k|k}$ , the Kalman filter estimate, and 3)  $\tilde{\mathbf{z}}_{j,k|MN}$ , the Kalman smoother estimate.

$$2) \tilde{\mathbf{z}}_{k|k'} = [(\tilde{\mathbf{z}}_{1,k|k'})^\top, \dots, (\tilde{\mathbf{z}}_{J,k|k'})^\top]^\top \in \mathbb{R}^{2J}$$

A collection of  $\{\tilde{\mathbf{z}}_{j,k|k'}\}_j$  in a single vector.

$$3) \mathbf{P}_{j,k|k'} = \mathbb{E}[(\tilde{\mathbf{z}}_{j,k} - \tilde{\mathbf{z}}_{j,k|k'}) (\tilde{\mathbf{z}}_{j,k} - \tilde{\mathbf{z}}_{j,k|k'})^\top \mid \{\hat{\sigma}_{j,m}^2\}_{j,m}, \mathbf{y}_{1:k'}, \hat{\theta}] \in \mathbb{R}^{2 \times 2}$$

Posterior covariance of  $\tilde{\mathbf{z}}_{j,k}$ . Just as in  $\tilde{\mathbf{z}}_{j,k|k'}$ , we are interested in three types, i.e.,  $\mathbf{P}_{j,k|k-1}$ ,  $\mathbf{P}_{j,k|k}$ , and  $\mathbf{P}_{j,k|MN}$ .

$$4) \mathbf{P}_{k|k'} = \text{blkdiag}(\mathbf{P}_{j,k|k'}) \in \mathbb{R}^{2J \times 2J}$$

A block diagonal matrix of  $J$  posterior covariance matrices.

$$5) \mathbf{A} = \text{blkdiag}(\exp(-\Delta/l_j) \mathbf{R}(\omega_j)) \in \mathbb{R}^{2J \times 2J}$$

A block diagonal transition matrix.

$$6) \mathbf{H} = (1, 0, \dots, 1, 0) \text{ The observation gain.}$$

## KALMAN FILTER/SMOOTHER

The Kalman filter equations are given as

$$\begin{aligned}&\tilde{\mathbf{z}}_{j,mN+n|mN+(n-1)} \\ &= \exp(-\Delta/l_j) \mathbf{R}(\omega_j) \tilde{\mathbf{z}}_{j,mN+(n-1)|mN+(n-1)} \\ &\mathbf{P}_{j,mN+n|mN+(n-1)} \\ &= \exp(-2\Delta/l_j) \mathbf{R}(\omega_j) \mathbf{P}_{j,mN+(n-1)|mN+(n-1)} \mathbf{R}^\top(\omega_j) \\ &\quad + \sigma_{j,m}^2 (1 - \exp(-2\Delta/l_j)) \\ &\mathbf{K}_{mN+n} \\ &= \mathbf{P}_{mN+n|mN+(n-1)} \mathbf{H}^\top (\mathbf{H} \mathbf{P}_{mN+n|mN+(n-1)} \mathbf{H}^\top + \sigma_\nu^2)^{-1} \\ &\tilde{\mathbf{z}}_{mN+n|mN+n} \\ &= \tilde{\mathbf{z}}_{mN+n|mN+(n-1)} \\ &\quad + \mathbf{K}_{mN+n} (\mathbf{y}_{mN+n} - \mathbf{H} \tilde{\mathbf{z}}_{mN+n|mN+(n-1)}) \\ &\mathbf{P}_{mN+n|mN+n} = (\mathbf{I}_{2J \times 2J} - \mathbf{K}_{mN+n} \mathbf{H}) \mathbf{P}_{mN+n|mN+(n-1)}.\end{aligned}$$

Subsequently, the Kalman smoother equations are given as

$$\begin{aligned}
& \mathbf{C}_{mN+n} \\
&= \mathbf{P}_{mN+n|mN+n} \mathbf{A}^T \mathbf{P}_{mN+(n+1)|mN+n}^{-1} \in \mathbb{R}^{2J \times 2J} \\
& \tilde{\mathbf{z}}_{mN+n|MN} \\
&= \tilde{\mathbf{z}}_{mN+n|mN+n} \\
&\quad + \mathbf{C}_{mN+n} (\tilde{\mathbf{z}}_{mN+(n+1)|MN} - \tilde{\mathbf{z}}_{mN+(n+1)|mN+n}) \\
& \mathbf{P}_{mN+n|MN} \\
&= \mathbf{P}_{mN+n|mN+n} \\
&\quad + \mathbf{C}_{mN+n} \mathbf{P}_{mN+(n+1)|MN} \mathbf{C}_{mN+n}^T \\
&\quad - \mathbf{C}_{mN+n} \mathbf{P}_{mN+(n+1)|mN+n} \mathbf{C}_{mN+n}^T
\end{aligned}$$

To obtain the mean reconstructed trajectory for the  $j^{\text{th}}$  oscillatory component,  $\{\hat{\mathbf{y}}_{j,k}\}_k$ , we take the real part of the  $j^{\text{th}}$  component of the smoothed mean,  $\hat{\mathbf{y}}_{j,k} = \mathbf{e}_{2j-1}^T \cdot \tilde{\mathbf{z}}_{k|MN}$ , where  $\mathbf{e}_{2j-1} \in \mathbb{R}^{2J}$  is a unit vector with the only non-zero value, equal to 1, at the entry  $2j-1$ .

The 95% credible interval for  $\hat{\mathbf{y}}_{j,k}$ , denoted as  $\text{CI}_{j,k}^{\text{lower}}/\text{CI}_{j,k}^{\text{upper}}$  for the upper/lower end, respectively, is given as

$$\begin{aligned}
\text{CI}_{j,k}^{\text{upper}} &= \mathbf{e}_{2j-1}^T \cdot \tilde{\mathbf{z}}_{k|MN} + 1.96 \cdot \sqrt{\mathbf{e}_{2j-1}^T \mathbf{P}_{k|MN} \mathbf{e}_{2j-1}} \\
\text{CI}_{j,k}^{\text{lower}} &= \mathbf{e}_{2j-1}^T \cdot \tilde{\mathbf{z}}_{k|MN} - 1.96 \cdot \sqrt{\mathbf{e}_{2j-1}^T \mathbf{P}_{k|MN} \mathbf{e}_{2j-1}}.
\end{aligned}$$

## FFBS ALGORITHM FOR $p(\{\mathbf{z}_j\}_j | \{\sigma_{j,m}^2\}_{j,m}, \mathbf{y}, \theta)$

To generate the MC trajectory samples from the posterior distribution  $p(\{\mathbf{z}_j\}_j | \{\sigma_{j,m}^2\}_{j,m}, \mathbf{y}, \theta)$ , we use the FFBS algorithm. The steps are summarized in Algorithm 2, which uses the Kalman estimates derived in the previous section. We denote  $s = 1, \dots, S$  as the MC sample index.

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### Algorithm 2: FFBS algorithm

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**Result:**  $\{\tilde{\mathbf{z}}_k^{(s)}\}_{k,s}^{MN,S}$

**for**  $s \leftarrow 1$  **to**  $S$  **do**

Sample  $\tilde{\mathbf{z}}_{MN}^{(s)}$  from  $\mathcal{N}(\tilde{\mathbf{z}}_{MN|MN}, \mathbf{P}_{MN|MN})$ .

**for**  $k \leftarrow MN-1$  **to**  $1$  **do**

Sample  $\tilde{\mathbf{z}}_k^{(s)}$  from  $\mathcal{N}(\tilde{\mu}_k, \tilde{\mathbf{P}}_k)$ , where

$\tilde{\mu}_k = \tilde{\mathbf{z}}_{k|k} + \mathbf{P}_{k|k} \mathbf{A}^T \mathbf{P}_{k+1|k}^{-1} (\tilde{\mathbf{z}}_{k+1}^{(s)} - \mathbf{A} \tilde{\mathbf{z}}_{k|k})$

$\tilde{\mathbf{P}}_k = \mathbf{P}_{k|k} - \mathbf{P}_{k|k} \mathbf{A}^T \mathbf{P}_{k+1|k}^{-1} \mathbf{A} \mathbf{P}_{k|k}$ .

**end**

**end**

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## I. COMPUTATIONAL EFFICIENCY OF PLSO VS. GP-PS

We show the runtime of PLSO and piecewise stationary GP (GP-PS) for inference of the mean trajectory of the hippocampus data ( $f_s = 1, 250$  Hz,  $J = 5$ , 2-second window) for varying data lengths (50, 100, 200 seconds corresponding to  $K = 6.25 \times 10^4, 1.25 \times 10^5, 2.5 \times 10^5$  sample points, respectively).

As noted in Section 5, the computational complexity of PLSO is  $O(J^2 K)$ , whereas the computational complexity of GP-PS is  $O(N^2 K)$ . Since  $N$ , the number of samples per window, is fixed (2,500 samples), we expect both PLSO and GP-PS to be linear in terms of the number of samples  $K$ . Table 1 indeed confirms that this is the case. However, we observe that PLSO is much more efficient than GP-PS.

Table 1: Runtime (s) for PLSO and GP-PS for varying length

	PLSO	GP-PS
$T=50$	1.7	346.8
$T=100$	3.1	700.6
$T=200$	6.5	1334.0

## J. SIMULATION EXPERIMENT (SECTION 5.1)

We simulate from the following model for  $1 \leq k \leq K$

$$\mathbf{y}_k = 10 \left( \frac{K-k}{K} \right) \mathbf{z}_{1,k}^{\Re} + 10 \cos^4(2\pi\omega_0 k) \mathbf{z}_{2,k}^{\Re} + \nu_k,$$

where  $\mathbf{z}_{1,k}$  and  $\mathbf{z}_{2,k}$  are from the PLSO stationary generative model, i.e.,  $\sigma_{j,m}^2 = \sigma_j^2$ . The parameters are  $\omega_0/\omega_1/\omega_2 = 0.04/1/10$  Hz,  $f_s = 200$  Hz,  $T = 100$  seconds,  $l_1 = l_2 = 1$ , and  $\nu_k \sim \mathcal{N}(0, 25)$ . This stationary process comprises two amplitude-modulated oscillations, namely one modulated by a slow-frequency ( $\omega_0 = 0.04$  Hz) sinusoid and the other a linearly-increasing signal Ba et al. [2014]. We assume a 2-second PS interval. For PLSO, we use  $J = 2$  components and 5 block coordinate descent iterations for optimizing  $\theta$  and  $\{\sigma_{j,m}^2\}_{j,m}$ .

## K. DETAILS OF THE TVAR MODEL

As explained in Section 5, the TVAR model is defined as

$$\mathbf{y}_k = \sum_{p=1}^P a_{p,k} \mathbf{y}_{k-p} + \varepsilon_k,$$

which can alternatively be written as

$$\begin{pmatrix} \mathbf{y}_k \\ \vdots \\ \mathbf{y}_{k-P+1} \end{pmatrix} = \underbrace{\begin{pmatrix} a_{1,k} & a_{2,k} & \cdots & a_{P-1,k} & a_{P,k} \\ 1 & 0 & \cdots & 0 & 0 \\ 0 & 1 & \cdots & 0 & 0 \\ \vdots & & & & \vdots \\ 0 & 0 & \cdots & 1 & 0 \end{pmatrix}}_{A_k} \begin{pmatrix} \mathbf{y}_{k-1} \\ \vdots \\ \mathbf{y}_{k-P} \end{pmatrix} + \varepsilon_k.$$

It is the transition matrix  $A_k$  that determines the oscillatory component profile at time  $k$ , such as the number of components and the center frequencies. Specifically,  $\{A_k\}_k$  are first fit to the data  $\mathbf{y}$  and then eigen-decomposition is performed on each of the estimated  $\{A_k\}_k$ . More in-depth technical details can be found in West et al. [1999].

We use publicly available code for the TVAR implementation<sup>1</sup>. We use TVAR order of  $p = 70$ , as the models with lower orders than this value did not capture the theta-band signal - The lowest frequency band in these cases were the gamma band ( $> 30$  Hz). Even with the higher orders of  $p$ , and various hyperparameter combinations, we observed that the slow and theta frequency band was still explained by a single oscillatory component. For the discount factor, we used  $\beta = 0.999$  to ensure that the TVAR coefficients and, consequently, the decomposed oscillatory components do not fluctuate much.

## L. ANESTHESIA EEG DATASET (SECTION 5.3)

We show spectral analysis results for the EEG data of a subject anesthetized with propofol (This is a different subject from the main text.) The data last  $T = 2,500$  seconds, sampled at  $f_s = 250$  Hz. We assume a 4-second PS interval, use  $J = 9$  components and 5 block coordinate descent iterations for optimizing  $\theta$  and  $\{\sigma_{j,m}^2\}_{j,m}$ .

Fig. 1 shows the STFT and the PLSO-estimated spectrograms. As noted in the main text, PLSO with stationarity assumption is too restrictive and fails to capture the time-varying spectral pattern. Both PLSO with  $\lambda = 0$  and  $\lambda = \lambda_{CV}$  are more effective in capturing such patterns, with the latter able to remove the artifacts and better recover the smoother dynamics.

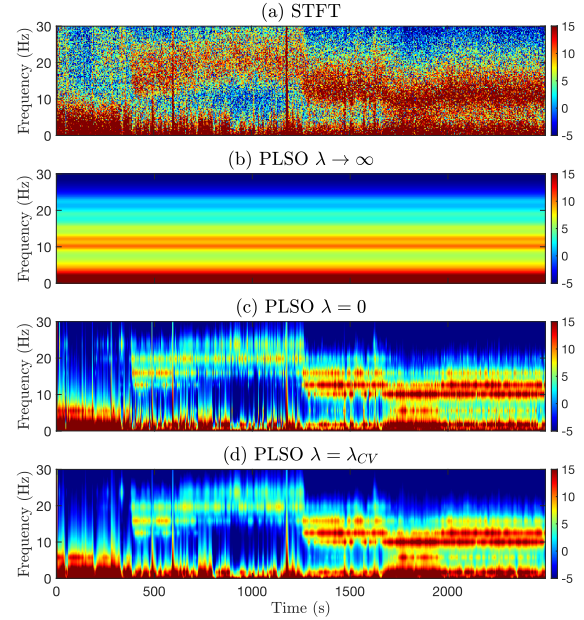


Figure 1: Spectrogram (in dB) under propofol anesthesia. (a) STFT of the data (b) PLSO with  $\lambda \rightarrow \infty$  (c) PLSO with  $\lambda = 0$  (d) PLSO with  $\lambda = \lambda_{CV}$ .

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<sup>1</sup><https://www2.stat.duke.edu/mw/mwsoftware/TVAR/index.html>