



Dynamic simulation of moderately thick annular system coupled with shape memory alloy and multi-phase nanocomposite face sheets

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Abstract

The current research work analyzes dynamics of a sandwich disk which is gently thick. The mentioned sandwich structure has honeycomb core, a couple of middle layers having fibers of shape memory alloy (SMA), and a couple of external layers of multi-scaled hybrid nanocomposite (MHC) considering in-plane force. The core in the shape of honeycomb is manufactured of aluminum due to its high stiffness and less density compared with other materials. Applying energy methods called the principle of Hamilton, we obtained governing motion equations of the mentioned structure and solved them using First-order shear-deformation-theory (FSDT), as well as generalized-differential-quadature-method (GDQM), respectively. To layers' joint, the compatibility equations have been taken into account. A parametric mathematical manipulation has been conducted to analyze the impacts of fibers of SMA, boundary conditions (BCs), internal loads, honeycomb network angle, ratio of external to internal radiuses, ratio of thickness to length of the honeycomb, weight fraction of CNTs, angle of fibers, ratio of honeycomb to face-sheet thickness on the frequency of the multi-phase sandwich disk. The outcomes derived reveal that for any amount of internal pressure and each BC, the relation of the honeycomb's thickness ratios to MHC layer (h_H/h_i) and sandwich structure's frequency is similar to quadratic function. Further results show that the effects of the fibers' angle on the frequency can be ignored for larger h_H/h_i amounts.

Keywords Honeycomb core · Porosity · Multi-layer disk · Various boundary conditions · MHC · SMA fibers

1 Introduction

Since applicable material and structures' thermomechanical response has been improved by applying the sandwich module structures, in the past decant, scientists discovered

a new and perfect method to boost the dynamic and static responses of the low-density beam, plate, disk, and shell [1–7] due to the mentioned fact, structures in honeycomb shape [8, 9] are presented to use in the pertained industries [10, 11]. The sandwich panel's frequency parameters made from Honeycomb Core have been analyzed by Mukhopadhyay et al. [12] using Energy Methods. In that paper, it is reported that core of honeycomb would enhance the free vibration characteristics and ultimately the sandwich panels' stiffness. Ref. [13] has investigated the impacts of a wide range of defects occurring to build honeycomb composite beams and determined the mechanical behavior of those structures in different frequency modes employing Fast Fourier Transform analyzer and finite element (FE) model. Considerable outcomes of this research revealed that the structure's frequency parameters would be reduced by increasing percentage of structure's defect. Frequency characteristics of honeycomb sandwich structures with various cores have been analyzed by Mozafari et al. [14] Employing experimental tests, they obtained the polyurethane foams' mechanical behaviors and evaluated the impact of the first

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resonance, the foams' impact and the shape of the state on the core's vibrations. The frequency characteristics of the structure's graded corrugated lattice core along with using an exact solution method to solve the governing motion equations has been evaluated by Ref. [15]. They studied the beam length's impact, thickness of facial leaf and graded parameters on the mentioned structure's frequency. The amplitude of the solar panel's vibrations which is manufactured by smart layers and honeycomb core has been controlled by Amini et al. [16]. Using the theories related to thin plate and employing Hamilton's principle, they obtained the BCs and equations of motion. Eventually, they discovered that elastoelectric impacts are playing prominent roles in the solar panel's frequency parameters. Ref. [17] examined the panels' post-buckling response with graphene particles reinforcements and honeycomb cores. The scientists' achievements show that the thickness of core, weight fraction of GPL, and geometric characters pertained to panel have vital roles on the sandwich panel's post-buckling. Sobhy [18] has analyzed the curved beam's bending characters reinforced with layers of graphene nanoplatelets along with honeycomb core. Employing DQM, they obtained the associated BCs and solved complicated motion equations. Sandwich panels' frequency parameters with honeycomb core employing FE model and experimental research has been conducted by Wang et al. [19]. They eventually revealed that, the face-sheet's thickness ratio, and density of filling foam have essential role in the sandwich panels' frequency with honeycomb core. Frequency investigation of a sandwich beam considering honeycomb hybrid core applying experimental techniques and FE model has been studied by Ref. [20]. Honeycomb sandwich shell's nonlinear frequency parameters has been reported by Zhang et al. [21]. They have solved the structure's governing motion equations with simply BCs through using method of homotopy perturbation. Recently, applying reinforcements of CNTs have attracted the attention of many scientists and researchers. For instance, an FG circular plate's major bending behavior boosted through using CNTs and covered by an elastic foundation has been studied by Keleshteri et al. [22]. They assert that the thick shell deformation models and von Kármán are used to present more precise results in their mathematical method. Moreover, for solving governing equations extracted through using energy methods, they applied the Mindlin–Plesha as well as GDQ approach. Their prominent result is that the CNT's volume fraction along with thickness would play an important role when it comes to the analysis of the circular disk's nonlinear frequency. Nonlinear free and forced an FG disk's frequency response employing the von Kármán model along with thin SDT has been conducted by Ansari and Torabi [23]. They highlighted basically on the enhanced GDQ approach for the solution method of

the FG disk's governing equation and presented a structure's large amplitude frequency. Torabi and Ansari [24] have presented that it is vital to develop motion equations of the FG-CNT-reinforced circular plate's large amplitude vibration based on general asymmetric's equations in the existence of elementary thermally stress for reaching precise outcomes. Finally, a huge number of researches presented the stability/sinstability of applicable structure in many articles [25–33].

Based on the highly scrutinized literature investigation, no one would assert there is a study on the sandwich disk's frequency analysis considering a core (honeycomb), including fiber of shape memory alloy (SMA), and a couple of external layers (MHC) under in-plane load. FSDT is employed to stress–strain formulation. Modified Halpin–Tsai model and the Rule of the mixture are involved in providing the efficient the MHC disk's material constant. Using energy methods, the structure's governing equations are extracted. Eventually, the impacts of some physically and geometrically factors of SMA fiber and MHC and in-plane loads on the dynamics of the mentioned structure are reported, in detail. The final results reveal that we would conclude that using the network of honeycomb as the structure's core boosts the structure's dynamics considerably.

2 Mathematical modeling

2.1 The homogenization process of MHC

The homogenization procedure is contained of a couple of basic stages due to the Halpin–Tsai theory [34], combined with a micromechanical model. The primary level is involved with calculating the efficient behaviors of the composite enhanced with CF as below [35]

$$E_{11} = V_F E_{11}^F + V_{\text{NCM}} E^{\text{NCM}}, \quad (1)$$

$$\frac{1}{E_{22}} = \frac{V_F}{E_{22}^F} + \frac{V_{\text{NCM}}}{E^{\text{NCM}}} - V_F V_{\text{NCM}} \times \frac{(v^F)^2 \frac{E^{\text{NCM}}}{E_{22}^F} + (v^{\text{NCM}})^2 \frac{E_{22}^F}{E^{\text{NCM}}} - 2v^F v^{\text{NCM}}}{V_F E_{22}^F + V_{\text{NCM}} E^{\text{NCM}}}, \quad (2)$$

$$\frac{1}{G_{12}} = \frac{V_F}{G_{12}^F} + \frac{V_{\text{NCM}}}{G^{\text{NCM}}}, \quad (3)$$

$$\rho = V_{\text{NCM}} \rho^{\text{NCM}} + V_F \rho^F, \quad (4)$$

$$v_{12} = V_{\text{NCM}} v^{\text{NCM}} + V_F v^F. \quad (5)$$

The nanocomposite matrixes [36–38] and fiber's volume fraction may be as [35]:

$$V_F + V_{NCM} = 1. \tag{6}$$

The following stage is organized to determine the efficient behaviors of the nano-scaled composite matrix enhanced with CNTs using the developed micromechanics of Halpin-Tsai [39]:

$$E^{NCM} = E^M \left(\frac{5}{8} \left(\frac{1 + 2\beta_{dd} V_{CNT}}{1 - \beta_{dd} V_{CNT}} \right) + \frac{3}{8} \left(\frac{1 + 2(l^{CNT}/d^{CNT})\beta_{dl} V_{CNT}}{1 - \beta_{dl} V_{CNT}} \right) \right). \tag{7}$$

Here, β_{dl} and β_{dd} may be calculated as the following relations

$$\beta_{dl} = \frac{(E_{11}^{CNT}/E^M)}{(l^{CNT}/2t^{CNT}) + (E_{11}^{CNT}/E^M)} - \frac{(d^{CNT}/4t^{CNT})}{(l^{CNT}/2t^{CNT}) + (E_{11}^{CNT}/E^M)},$$

$$\beta_{dd} = \frac{(E_{11}^{CNT}/E^M)}{(d^{CNT}/2t^{CNT}) + (E_{11}^{CNT}/E^M)} - \frac{(d^{CNT}/4t^{CNT})}{(d^{CNT}/2t^{CNT}) + (E_{11}^{CNT}/E^M)}. \tag{8}$$

The CNTs' volume fraction would be calculated as follow:

$$V_{CNT}^* = \frac{W_{CNT}}{W_{CNT} + \left(\frac{\rho_{CNT}}{\rho^M}\right)(1 - W_{CNT})}. \tag{9}$$

However, a broad range of MHC distribution through orientation of thickness would be written by [40]:

$$V_{CNT} = 4V_{CNT}^* \frac{|\xi_j|}{h} \quad \text{FG-X}$$

$$V_{CNT} = 2V_{CNT}^* \left(1 - 2\frac{|\xi_j|}{h}\right) \quad \text{FG-O} \tag{10}$$

$$V_{CNT} = V_{CNT}^* \quad \text{FG-UD}$$

where $\xi_j = \left(\frac{1}{2} + \frac{1}{2N_i} - \frac{j}{N_i}\right)$, $i = 1, 2, \dots, N_i$.

Moreover, the sum of V_{CNT} and V_M is equal to one as below (Fig. 1):

$$V_{CNT} + V_M = 1. \tag{11}$$

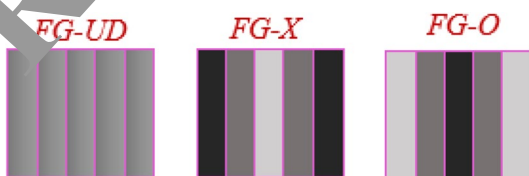


Fig. 1 Distribution of CNT through the thickness of the MHL composite

Eventually, the MHC face sheets' mechanical properties may be written as:

$$\rho^{NCM} = V_{CNT}\rho^{CNT} + V_M\rho^M \tag{12}$$

$$v^{NCM} = v^M, \tag{13}$$

$$G^{NCM} = \frac{E^{NCM}}{2(1 + v^{NCM})} \tag{14}$$

In Fig. 2, different types of porosity distribution called, X, U, along with O are given [41–44]. The shear modulus, Young's modulus, as well as density, are as follow:

$$\tilde{G}_{12}(z) = \frac{G}{2(1 - v(z))} \tag{15a}$$

$$\tilde{E}_{22} = E_{22}(1 - e_0s(z)), \tag{15b}$$

$$\tilde{E}_{11} = E_{11}(1 - e_0s(z)), \tag{15c}$$

$$\tilde{\rho}(z) = [-e_m s(z) + 1]\rho(z) + V_{ncm}\rho_{ncm}, \tag{15d}$$

where [45]:

$$s = \begin{cases} s_o & \text{PD-UD} \\ s_o \cos\left(\frac{\pi}{4} + \frac{\pi z}{2h}\right) & \text{PD-X} \\ s_o \cos\left(\frac{\pi z}{h}\right) & \text{PD-O} \end{cases} \tag{16}$$

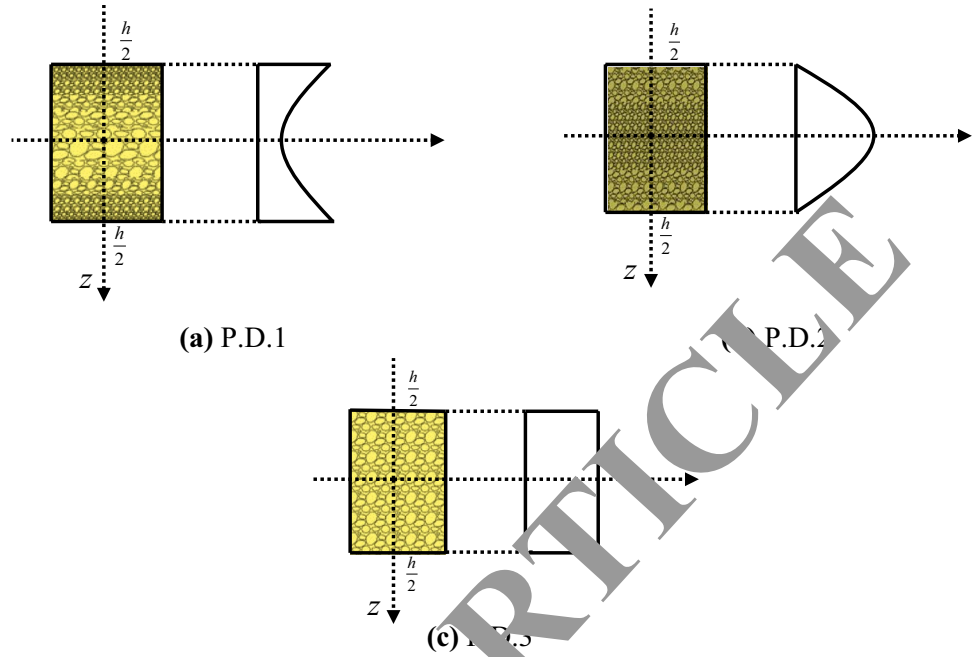
Due to the Gaussian Random Field scheme (Figs. 3 and 4), we would have [45]:

$$e_m = \frac{1.121 \left[1 - (1 - e_0s(z))^{\frac{1}{2.3}}\right]}{s(z)}, \tag{17}$$

the porous disk's Poisson's ratio pertaining to the Field of closed-cell Gaussian Random would be computed as [45]:

$$\tilde{v}_{12} = 0.221 \left(1 - \frac{\tilde{\rho}(z)}{\rho}\right) + v_{12} \left[1 + 0.342 \left(1 - \frac{\tilde{\rho}(z)}{\rho}\right)^2 - 1.21 \left(1 - \frac{\tilde{\rho}(z)}{\rho}\right)\right], \tag{18a}$$

Fig. 2 Patterns of porosity distribution through the thickness of MHC [46]



$$\tilde{v}_{21} = \tilde{v}_{12} \frac{\tilde{E}_{11}}{\tilde{E}_{22}} \tag{18b}$$

Moreover, when the porosity distributions' various types with the disk's total masses are the same, the amount of the S_0 may be obtained as [45]:

$$s_0 = \frac{1}{e_0} \left[1 - \frac{\left(\frac{1}{h} \int_{-h/2}^{h/2} \frac{\tilde{\rho}(z)}{\rho} dz + 0.121 \right)^{2.3}}{1.121} \right] \tag{19}$$

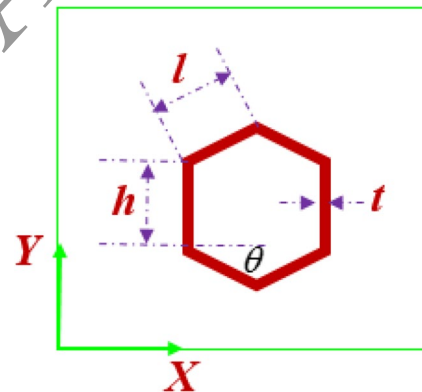


Fig. 3 The hexagonal cell geometry

3 Honeycomb core modeling

The hexagonal cell configuration is demonstrated in Fig. 3. Due to the Gibson theory, we would have [47]:

$$E^* = \frac{\sigma_1}{\epsilon_1} = E_s \left(\frac{t}{l} \right)^3 \frac{\cos(\theta_h)}{(h/l + \sin(\theta_h)) \sin^2(\theta_h)} \frac{1}{1 + (t/l)^2 \cot^2(\theta_h)} \tag{20-a}$$

$$E_{22}^* = \frac{\sigma_2}{\epsilon_2} = E_s \left(\frac{t}{l} \right)^3 \frac{(h/l + \sin(\theta_h))}{\cos^3(\theta_h)} \frac{1}{1 + (t/l)^2 \cot^2(\theta_h)} \tag{20-b}$$

$$v_{12}^* = -\frac{\epsilon_2}{\epsilon_1} = \frac{\cos^2(\theta_h)}{(h/l + \sin(\theta_h)) \sin(\theta_h)} \frac{1 - (t/l)^2}{1 + \cot^2(\theta_h)(t/l)^2} \tag{20-c}$$

$$v_{21}^* = -\frac{\epsilon_1}{\epsilon_2} = \frac{(h/l + \sin(\theta_h)) \sin(\theta_h)}{\cos^2(\theta_h)} \frac{-(t/l)^2 + 1}{1 + ((h/l) \sec^2(\theta_h) + \tan^2(\theta_h))(t/l)^2} \tag{20-d}$$

$$G_{12}^* = E_S \left(\frac{t}{l}\right)^3 \frac{(h/l + \sin(\theta_h))}{(h/l)^2 \cos(\theta_h)} \frac{1}{R} \tag{20-e}$$

The above equation η explains the bottom, top, along with core layers.

$$R = \left(2\frac{h}{l} + 1 + \left(\frac{t}{l}\right)^2 \frac{h/l + \sin(\theta_h)}{(h/l)^2} [(\sin(\theta_h) + h/l) \tan^2(\theta_h) + \sin(\theta_h)]\right) \tag{20-f}$$

$$\frac{\rho^*}{\rho_S} = \frac{\left(\frac{t}{l}\right)\left(\frac{h}{l} + 2\right)}{2 \cos(\theta_h)\left(\frac{h}{l} + \sin(\theta_h)\right)} \tag{20-g}$$

4 SMA characteristics in the composite's matrix

The SMA properties are reported by Ref. [48] the present research, supposed that the fibers of SMA have been distributed uniformly. The composite disk's elastic parameters with fibers of SMA may be derived by [48]:

$$\tilde{E}_{11}^{sma} = V_m E_m^m + V_S E_S \tag{21a}$$

$$\tilde{E}_{22}^{sma} = \frac{E_S E_{22}^m}{V_m E_{22}^m + V_m E_m} \tag{21b}$$

$$\tilde{G}_{12}^{sma} = \frac{G_S G_{12}^m}{G_S G_{12}^m + G_S V_m} \tag{21c}$$

$$\tilde{G}_{23}^{sma} = G_{23}^m V_m + G_S V_S \tag{21d}$$

$$\tilde{G}_{13}^{sma} = \tilde{G}_{23}^{sma} \tag{21e}$$

$$\tilde{\nu}_{12}^{sma} = V_m \nu_{12}^m + V_S \nu_S \tag{21f}$$

$$\tilde{\rho}^{sma} = V_m \rho_m + V_S \rho_S \tag{21g}$$

5 The disk's displacement fields

According to the shear stress theories [49–56], and due to the FSDT, the fields of displacement would be expressed by the following equations relations [57]:

$$\begin{aligned} \alpha^n &= \alpha_0^n + z \xi_R^n \\ \beta^n &= \beta_0^n + z \xi_\theta^n \\ \chi^n &= \chi_0^n \end{aligned} \tag{22}$$

6 Stress–Strain honeycomb core's relations

Due to the FSDT, the stress–strain relation may be given as [57–63]:

$$\begin{bmatrix} \sigma_{RR} \\ \sigma_{\theta\theta} \\ \sigma_{R\theta} \\ \sigma_{Rz} \\ \sigma_{\theta z} \end{bmatrix}^c = \begin{bmatrix} Q_{11} & Q_{12} & 0 & 0 & 0 \\ Q_{12} & Q_{22} & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & Q_{55} & 0 \\ 0 & 0 & 0 & 0 & Q_{44} \end{bmatrix}^c \begin{bmatrix} \epsilon_{RR} \\ \epsilon_{\theta\theta} \\ \gamma_{R\theta} \\ \gamma_{Rz} \\ \gamma_{\theta z} \end{bmatrix}^c$$

$$Q_{11}^c = \frac{E_{11}^*}{1 - \nu_{21}^* \nu_{12}^*}, \quad Q_{22}^c = \frac{E_{22}^*}{1 - \nu_{12}^* \nu_{21}^*}, \quad Q_{12}^c = \frac{\nu_{21}^* E_{22}^*}{1 - \nu_{12}^* \nu_{21}^*},$$

$$Q_{44}^c = G_{22}^*, \quad Q_{55}^c = G_{13}^*, \quad Q_{66}^c = G_{23}^*, \quad G_{13}^* = G_{23}^* = G_{12}^* \tag{23}$$

So, the strain elements may be defined as [64]:

$$\begin{bmatrix} \epsilon_{RR} \\ \epsilon_{\theta\theta} \\ \gamma_{R\theta} \\ \gamma_{Rz} \\ \gamma_{\theta z} \end{bmatrix}^\psi = \begin{bmatrix} \epsilon_{RR}^0 \\ \epsilon_{\theta\theta}^0 \\ \gamma_{R\theta}^0 \\ \gamma_{Rz}^0 \\ \gamma_{\theta z}^0 \end{bmatrix}^\psi + z \begin{bmatrix} \kappa_{RR} \\ \kappa_{\theta\theta} \\ \kappa_{R\theta} \\ \kappa_{Rz} \\ \kappa_{\theta z} \end{bmatrix}^\psi \tag{24}$$

Equation (24) would be rewritten as

$$\begin{bmatrix} \epsilon_{RR}^0 \\ \epsilon_{\theta\theta}^0 \\ \gamma_{R\theta}^0 \\ \gamma_{Rz}^0 \\ \gamma_{\theta z}^0 \end{bmatrix}^c = \begin{bmatrix} \frac{\partial \alpha_0}{\partial R} \\ \frac{\alpha_0}{R} + \frac{\partial \beta_0}{R \partial \theta} \\ \frac{\partial \alpha_0}{R \partial \theta} + \frac{\partial \beta_0}{\partial R} - \frac{\beta_0}{R} \\ \xi_R + \frac{\partial \chi}{\partial R} \\ \xi_\theta + \frac{\partial \chi}{R \partial \theta} \end{bmatrix}^c, \quad \begin{bmatrix} \kappa_{RR} \\ \kappa_{\theta\theta} \\ \kappa_{R\theta} \\ \kappa_{Rz} \\ \kappa_{\theta z} \end{bmatrix}^c = \begin{bmatrix} \frac{\partial \xi_R}{\partial R} \\ \frac{\xi_R}{R} + \frac{\partial \xi_\theta}{R \partial \theta} \\ \frac{\partial \xi_R}{R \partial \theta} + \frac{\partial \xi_\theta}{\partial R} - \frac{\xi_\theta}{R} \\ 0 \\ 0 \end{bmatrix}^c \tag{25}$$

MHC angle-ply laminated disk's Stress–strain equations [65–69] and SMA layer would be given as below [57]:

$$\begin{bmatrix} \sigma_{RR} \\ \sigma_{\theta\theta} \\ \sigma_{R\theta} \\ \sigma_{Rz} \\ \sigma_{\theta z} \end{bmatrix}^\psi = \begin{bmatrix} \hat{Q}_{11} & \hat{Q}_{12} & 0 & 0 & \hat{Q}_{16} \\ \hat{Q}_{12} & \hat{Q}_{22} & 0 & 0 & \hat{Q}_{26} \\ 0 & 0 & \hat{Q}_{44} & \hat{Q}_{45} & 0 \\ 0 & 0 & \hat{Q}_{45} & \hat{Q}_{55} & 0 \\ \hat{Q}_{16} & \hat{Q}_{26} & 0 & 0 & \hat{Q}_{66} \end{bmatrix}^\psi \begin{bmatrix} \epsilon_{RR} \\ \epsilon_{\theta\theta} \\ \gamma_{R\theta} \\ \gamma_{Rz} \\ \gamma_{\theta z} \end{bmatrix}^\psi \tag{26}$$

In the aforementioned relation, ψ explains the bottom and top layers of the structure.

Here

$$\hat{Q}_{11}^{\psi} = \sin^4 \theta_f \tilde{Q}_{22}^{\psi} + \cos^4 \theta_f \tilde{Q}_{11}^{\psi} + 2 \sin^2 \theta_f \cos^2 \theta_f (\tilde{Q}_{12}^{\psi} + 2\tilde{Q}_{66}^{\psi}) \tag{27a}$$

$$\hat{Q}_{12}^{\psi} = +(\sin^4 \theta_f + \cos^4 \theta_f) \tilde{Q}_{12}^{\psi} + \sin^2 \theta_f \cos^2 \theta_f (\tilde{Q}_{11}^{\psi} + \tilde{Q}_{22}^{\psi} - 4\tilde{Q}_{66}^{\psi}) \tag{27b}$$

$$\hat{Q}_{16}^{\psi} = \cos \theta_f \sin^3 \theta_f (\tilde{Q}_{66}^{\psi} + 2\tilde{Q}_{12}^{\psi} - 2\tilde{Q}_{22}^{\psi}) + \cos^3 \theta_f \sin \theta_f (2\tilde{Q}_{11}^{\psi} - 2\tilde{Q}_{12}^{\psi} - \tilde{Q}_{66}^{\psi}) \tag{27c}$$

$$\hat{Q}_{22}^{\psi} = 2 \sin^2 \theta_f \cos^2 \theta_f \tilde{Q}_{12}^{\psi} + \sin^4 \theta_f \tilde{Q}_{11}^{\psi} + \cos^4 \theta_f \tilde{Q}_{22}^{\psi} + 2 \sin^2 \theta_f \cos^2 \theta_f (\tilde{Q}_{12}^{\psi} + 2\tilde{Q}_{66}^{\psi}) \tag{27d}$$

$$\hat{Q}_{26}^{\psi} = \cos^3 \theta_f \sin \theta_f (2\tilde{Q}_{12}^{\psi} - 2\tilde{Q}_{22}^{\psi} + \tilde{Q}_{66}^{\psi}) + \cos \theta_f \sin^3 \theta_f (2\tilde{Q}_{11}^{\psi} - 2\tilde{Q}_{12}^{\psi} - \tilde{Q}_{66}^{\psi}) \tag{27e}$$

$$\hat{Q}_{44}^{\psi} = \sin^2 \theta_f \tilde{Q}_{55}^{\psi} + \cos^2 \theta_f \tilde{Q}_{44}^{\psi} \tag{27f}$$

$$\hat{Q}_{45}^{\psi} = \cos \theta_f \sin \theta_f (-\tilde{Q}_{44}^{\psi} + \tilde{Q}_{55}^{\psi}) \tag{27g}$$

$$\hat{Q}_{55}^{\psi} = \cos^2 \theta_f \tilde{Q}_{55}^{\psi} + \sin^2 \theta_f \tilde{Q}_{44}^{\psi} \tag{27h}$$

$$\hat{Q}_{66}^{\psi} = \tilde{Q}_{66}^{\psi} (\cos^2 \theta_f - \sin^2 \theta_f)^2 + 4 \sin^2 \theta_f \cos^2 \theta_f (\tilde{Q}_{11}^{\psi} \tilde{Q}_{22}^{\psi} - 2\tilde{Q}_{12}^{\psi}) \tag{27i}$$

The terms involved in Eq. (27a–27i) would be obtained as [70–74]:

$$\tilde{Q}_{11}^{\psi} = \frac{\tilde{E}_{11}^{\zeta}}{1 - \tilde{\nu}_{12}^{\zeta} \tilde{\nu}_{21}^{\zeta}}, \tilde{Q}_{12}^{\psi} = \frac{\tilde{\nu}_{12}^{\zeta} \tilde{E}_{22}^{\zeta}}{-\tilde{\nu}_{12}^{\zeta} \tilde{\nu}_{21}^{\zeta}}, \tilde{Q}_{22}^{\psi} = \frac{\tilde{E}_{22}^{\zeta}}{1 - \tilde{\nu}_{12}^{\zeta} \tilde{\nu}_{21}^{\zeta}}, \tilde{Q}_{44}^{\psi} = \tilde{G}_{12}^{\zeta}, \tilde{Q}_{55}^{\psi} = \tilde{G}_{23}^{\zeta}, \tilde{Q}_{66}^{\psi} = \tilde{G}_{13}^{\zeta}.$$

The aforementioned relations define ζ as the layer of MHLc and SMA. So, the strain elements could be given as [75, 76]:

$$\begin{Bmatrix} \epsilon_{RR} \\ \epsilon_{\theta\theta} \\ \gamma_{R\theta} \\ \gamma_{Rz} \\ \gamma_{\theta z} \end{Bmatrix}^{\psi} = \begin{Bmatrix} \epsilon_{RR}^0 \\ \epsilon_{\theta\theta}^0 \\ \gamma_{R\theta}^0 \\ \gamma_{Rz}^0 \\ \gamma_{\theta z}^0 \end{Bmatrix}^{\psi} + z \begin{Bmatrix} \kappa_{RR} \\ \kappa_{\theta\theta} \\ \kappa_{R\theta} \\ \kappa_{Rz} \\ \kappa_{\theta z} \end{Bmatrix}^{\psi} \tag{28}$$

Equation (28) can be rewritten as

$$\begin{Bmatrix} \epsilon_{RR}^0 \\ \epsilon_{\theta\theta}^0 \\ \gamma_{R\theta}^0 \\ \gamma_{Rz}^0 \\ \gamma_{\theta z}^0 \end{Bmatrix}^{\psi} = \begin{Bmatrix} \frac{\partial \alpha_0}{\partial R} \\ \frac{\alpha_0}{R} + \frac{\partial \beta_0}{R \partial \theta} \\ \frac{\partial \alpha_0}{R \partial \theta} + \frac{\partial \beta_0}{\partial R} - \frac{\beta_0}{R} \\ \xi_R + \frac{\partial \chi}{\partial R} \\ \xi_{\theta} + \frac{\partial \chi}{R \partial \theta} \end{Bmatrix}^{\psi}, \begin{Bmatrix} \kappa_{RR} \\ \kappa_{\theta\theta} \\ \kappa_{R\theta} \\ \kappa_{Rz} \\ \kappa_{\theta z} \end{Bmatrix}^{\psi} = \begin{Bmatrix} \frac{\partial \xi_R}{\partial R} \\ \frac{\xi_R}{R} + \frac{\partial \xi_{\theta}}{R \partial \theta} \\ \frac{\partial \xi_R}{R \partial \theta} + \frac{\partial \xi_{\theta}}{\partial R} - \frac{\xi_{\theta}}{R} \\ 0 \\ 0 \end{Bmatrix}^{\psi} \tag{29}$$

7 Compatibility equations

The conditions of compatibility supposing excellent bonding between the layers of composite and the core that would be expressed as below:

$$\alpha^c(z_c = -h_c/2) = \alpha^{SMAB}(z_{SMAB} = h_{SMAB}/2), \tag{30a}$$

$$\beta^c(z_c = -h_c/2) = \beta^{SMAB}(z_{SMAB} = h_{SMAB}/2), \tag{30b}$$

$$\chi^c(z_c = -h_c/2) = \chi^{SMAB}(z_{SMAB} = h_{SMAB}/2), \tag{30c}$$

$$\alpha^c(z_c = h_c/2) = \alpha^{SMAI}(z_{SMAI} = -h_{SMAI}/2), \tag{30d}$$

$$\beta^c(z_c = h_c/2) = \beta^{SMAI}(z_{SMAI} = -h_{SMAI}/2), \tag{30e}$$

$$\chi^c(z_c = h_c/2) = \chi^{SMAI}(z_{SMAI} = -h_{SMAI}/2), \tag{30f}$$

$$\alpha^{MHCi}(z_{MHCi} = -h_{cMHCi}/2) = \alpha^{SMAI}(z_{SMAI} = h_{SMAI}/2), \tag{30g}$$

$$\beta^{MHCt}(z_{MHCt} = -h_{MHCt}/2) = \beta^{SMAt}(z_{SMAt} = h_{SMAt}/2), \quad (30h)$$

$$\chi^{MHCt}(z_{MHCt} = -h_{MHCt}/2) = \chi^{SMAt}(z_{SMAt} = h_{SMAt}/2), \quad (30i)$$

$$\alpha^{MHCb}(z_{MHCb} = h_{MHCb}/2) = \alpha^{SMAb}(z_{SMAb} = -h_{SMAb}/2), \quad (30j)$$

$$\chi^{MHCb}(z_{MHCb} = h_{MHCb}/2) = \beta^{SMAb}(z_{SMAb} = -h_{SMAb}/2), \quad (30k)$$

$$\chi^{MHCb}(z_{MHCb} = h_{MHCb}/2) = \chi^{SMAb}(z_{SMAb} = -h_{SMAb}/2). \quad (30l)$$

$$\delta U^{*\eta} = \frac{1}{2} \iiint_V \sigma_{ij}^{\eta} \delta \epsilon_{ij}^{\eta} dV$$

$$= \int_A \left[\begin{aligned} & \left(P_{RR} \frac{\partial \delta \alpha_0}{\partial R} + Q_{RR} \frac{\partial \delta \alpha_1}{\partial R} \right) \\ & + \left(P_{\theta\theta} \frac{\partial \delta \beta_0}{R \partial \theta} + Q_{\theta\theta} \frac{\partial \delta \beta_1}{R \partial \theta} + P_{\theta\theta} \frac{\delta \alpha_0}{R} + Q_{\theta\theta} \frac{\delta \alpha_1}{R} \right) \\ & + \left(P_{R\theta} \frac{\partial \delta \beta_0}{\partial R} + Q_{R\theta} \frac{\partial \delta \beta_1}{\partial R} + P_{R\theta} \frac{\partial \delta \alpha_0}{R \partial \theta} + Q_{R\theta} \frac{\partial \delta \alpha_1}{R \partial \theta} \right) \\ & + \left(-P_{R\theta} \frac{\delta \beta_0}{R} - Q_{R\theta} \frac{\delta \beta_1}{R} \right) \\ & + \left((P_{Rz}) \left(\delta \alpha_1 + \frac{\partial \delta \chi_0}{\partial R} \right) + (P_{\theta z}) \left(\delta \alpha_1 + \frac{\partial \delta \chi_0}{R \partial \theta} \right) \right) \end{aligned} \right] dA \quad (35)$$

8 Developed principle of Hamilton

Due to the energy method, there are equations between motion equations and BCs defined as [77, 78]:

$$\int_{t_1}^{t_2} (\delta T^* - \delta U^* + \delta W^*)^{\eta} dt = 0 \quad (31)$$

The related rotating system's kinetic energy may be calculated as [79–84]:

$$T^{*\eta} = \int_V \frac{1}{2} \rho^{\eta} \left[\left(\frac{\partial \alpha}{\partial t} \right)^2 + \left(\frac{\partial \beta}{\partial t} \right)^2 + \left(\frac{\partial \chi}{\partial t} \right)^2 \right]^{\eta} dV \quad (32)$$

$$\delta T^{*\eta} = \int_V \rho^{\eta} \left(\frac{\partial U}{\partial t} \frac{\partial \delta U}{\partial t} + \frac{\partial V}{\partial t} \frac{\partial \delta V}{\partial t} + \frac{\partial W}{\partial t} \frac{\partial \delta W}{\partial t} \right)^{\eta} dV :$$

$$\delta T^{*\eta} = \int_{R_1}^{R_2} \int_0^{\theta} \left[\begin{aligned} & \left\{ -I_0 \frac{\partial^2 \alpha_0}{\partial t^2} - I_1 \frac{\partial^2 \xi_R}{\partial t^2} \right\} \delta \alpha + \left\{ -I_1 \frac{\partial^2 \alpha_0}{\partial t^2} - I_2 \frac{\partial^2 \xi_R}{\partial t^2} \right\} \delta \xi_R \\ & + \left\{ c_1 I_3 \frac{\partial^2 \alpha_0}{\partial t^2} + c_1 I_4 \frac{\partial^2 \xi_R}{\partial t^2} \right\} \delta \xi_R + \left\{ -I_0 \frac{\partial^2 \beta_0}{\partial t^2} - I_1 \frac{\partial^2 \xi_{\theta}}{\partial t^2} \right\} \delta \beta \\ & + \left\{ -I_1 \frac{\partial^2 \beta_0}{\partial t^2} - I_2 \frac{\partial^2 \xi_{\theta}}{\partial t^2} \right\} \delta \xi_{\theta} + \left\{ -I_0 \frac{\partial^2 \chi_0}{\partial t^2} \right\} \delta \chi \end{aligned} \right] R dR d\theta \quad (33)$$

where:

$$\{I_i\} = \int_V \rho^{\eta} (z^i) dV, \quad i = 0 : 4 \quad (34)$$

However, the presented composite structure's strain energy [85–88] could be determined as:

which [89]:

$$\{P_{RR}, Q_{RR}\}^{\eta} = \int_z \{ \sigma_{RR}, z \sigma_{RR} \}^{\eta} dz;$$

$$\{P_{\theta\theta}, Q_{\theta\theta}\}^{\eta} = \int_z \{ \sigma_{\theta\theta}, z \sigma_{\theta\theta} \}^{\eta} dz;$$

$$\{P_{Rz}, Q_{Rz}, T_{Rz}\}^{\eta} = \int_z \{ \sigma_{Rz}, z \sigma_{Rz}, z^2 \sigma_{Rz} \}^{\eta} dz;$$

$$\{P_{R\theta}, Q_{R\theta}, T_{R\theta}\}^{\eta} = \int_z \{ \sigma_{R\theta}, z \sigma_{R\theta}, z^2 \sigma_{R\theta} \}^{\eta} dz;$$

$$\{P_z, Q_{\theta z}, T_{\theta z}\}^{\eta} = \int_z \{ \sigma_{\theta z}, z \sigma_{\theta z}, z^2 \sigma_{\theta z} \}^{\eta} dz.$$

Additionally, the variation of the work done [82, 90–93]

by mechanical force (internal load) can be obtained as

$$W^{*\eta} = - \int_A P^A \left(\frac{\partial^2 \chi_0^{\eta}}{\partial R^2} \right) dA \quad (37)$$

Ultimately, corresponding BCs and governing equations would be extracted by inserting Eqs. (37), (35), and (33) in principle of Hamilton (Eq. (31)) which may be written as:

$\delta\alpha_0^n$:

$$\frac{\partial}{\partial R} P_{RR}^n - \frac{P_{\theta\theta}^n - P_{RR}^n}{R} + \frac{\partial}{R\partial\theta} P_{R\theta}^n = I_0^n \frac{\partial^2 \alpha_0^n}{\partial t^2} + I_1^n \frac{\partial^2 \alpha_1^n}{\partial t^2} \quad (38a)$$

$\delta\beta_0^n$:

$$\frac{\partial}{R\partial\theta} P_{\theta\theta}^n + \frac{2P_{R\theta}^n}{R} + \frac{\partial}{\partial R} P_{R\theta}^n = I_0^n \frac{\partial^2 \beta_0^n}{\partial t^2} + I_1^n \frac{\partial^2 \beta_1^n}{\partial t^2} \quad (38b)$$

$\delta\chi_0^n$:

$$2 \frac{\partial}{\partial R} (P_{Rz})^n + 2 \frac{\partial}{R\partial\theta} (P_{\theta z})^n - P^A \frac{\partial^2 \chi_0}{\partial R^2} = \left(I_0 \frac{\partial^2 \chi_0}{\partial t^2} \right)^n \quad (38c)$$

$\delta\xi_R^n$:

$$\frac{\partial Q_{RR}}{\partial R} - \frac{Q_{\theta\theta}}{R} + \frac{1}{R} \frac{\partial Q_{R\theta}}{\partial\theta} - (P_{Rz}) = I_1 \frac{\partial^2 \alpha_0}{\partial t^2} + I_2 \frac{\partial^2 \xi_R}{\partial t^2} \quad (38d)$$

$\delta\xi_\theta$:

$$\frac{\partial}{R\partial\theta} Q_{\theta\theta} + \frac{2}{R} Q_{R\theta} + \frac{\partial}{\partial R} Q_{R\theta} - Q_{\theta z} = I_1^n \frac{\partial^2 \beta_0}{\partial t^2} + I_2^n \frac{\partial^2 \xi_\theta}{\partial t^2} \quad (38e)$$

Furthermore, general pertained BCs would be written as:

$$\delta\alpha^n = 0 \text{ or } P_{RR}^n \hat{n}_R + \frac{P_{R\theta}^n}{R} \hat{n}_\theta = 0 \quad (39a)$$

$$\delta\beta^n = 0 \text{ or } P_{R\theta}^n \hat{n}_R + \frac{P_{\theta\theta}^n}{R} \hat{n}_\theta = 0 \quad (39b)$$

$$\delta\chi^n = 0 \text{ or } [2(P_{Rz})]^n \hat{n}_R + \left[\frac{2}{R} (P_{\theta z}) \right]^n \hat{n}_\theta = 0 \quad (39c)$$

$$\delta\chi_R^n = 0 \text{ or } [Q_{RR}]^n \hat{n}_R + \left[\frac{Q_{R\theta}}{R} \right]^n \hat{n}_\theta = 0 \quad (39d)$$

$$\delta\chi_\theta^n = 0 \text{ or } [Q_{R\theta}]^n \hat{n}_R + \left[\frac{Q_{\theta\theta}}{R} \right]^n \hat{n}_\theta = 0 \quad (39e)$$

By inserting eq. (36) into Eqs. (38a–e), the governing equations of the MHLCD are obtained as

$$\begin{aligned} \delta\alpha_0^n : & \frac{\partial}{\partial R} \left(\left\{ \mathbb{R}_{11} \frac{\partial \alpha_0^n}{\partial R} + \mathbb{Z}_{11} \frac{\partial \xi_R^n}{\partial R} \right\} + \left\{ \frac{\mathbb{R}_{12}}{R} \alpha_0^n + \frac{\mathbb{Z}_{12}}{R} \xi_R^n + \frac{\mathbb{R}_{12}}{R} \frac{\partial \beta_0^n}{\partial\theta} + \frac{\mathbb{Z}_{12}}{R} \frac{\partial \xi_\theta^n}{\partial\theta} \right\} \right) \\ & + \frac{1}{R} \frac{\partial}{\partial\theta} \left(\frac{\mathbb{R}_{66}}{R} \frac{\partial \alpha_0^n}{\partial\theta} + \frac{\mathbb{Z}_{66}}{R} \frac{\partial \xi_R^n}{\partial\theta} + \mathbb{R}_{66} \frac{\partial \beta_0^n}{\partial R} + \mathbb{Z}_{66} \frac{\partial \xi_\theta^n}{\partial R} - \frac{\mathbb{R}_{66}}{R} \beta_0^n - \frac{\mathbb{Z}_{66}}{R} \xi_\theta^n \right) \\ & - \frac{1}{R} \left(\left\{ \mathbb{R}_{12} \frac{\partial \alpha_0^n}{\partial R} + \mathbb{Z}_{12} \frac{\partial \xi_R^n}{\partial R} \right\} + \left\{ \frac{\mathbb{R}_{22}}{R} \alpha_0^n + \frac{\mathbb{Z}_{22}}{R} \xi_R^n + \frac{\mathbb{R}_{22}}{R} \frac{\partial \beta_0^n}{\partial\theta} + \frac{\mathbb{Z}_{22}}{R} \frac{\partial \xi_\theta^n}{\partial\theta} \right\} \right) \\ & = I_0^n \frac{\partial^2 \alpha_0^n}{\partial t^2} + I_1^n \frac{\partial^2 \alpha_1^n}{\partial t^2} \end{aligned} \quad (40a)$$

$$\begin{aligned} \delta\beta_0^n : & \frac{1}{R} \frac{\partial}{\partial\theta} \left(\left\{ \mathbb{R}_{12} \frac{\partial \alpha_0^n}{\partial R} + \mathbb{Z}_{12} \frac{\partial \xi_R^n}{\partial R} \right\} + \left\{ \frac{\mathbb{R}_{22}}{R} \alpha_0^n + \frac{\mathbb{Z}_{22}}{R} \xi_R^n + \frac{\mathbb{R}_{22}}{R} \frac{\partial \beta_0^n}{\partial\theta} + \frac{\mathbb{Z}_{22}}{R} \frac{\partial \xi_\theta^n}{\partial\theta} \right\} \right) \\ & + \frac{\partial}{\partial\theta} \left(\frac{\mathbb{R}_{66}}{R} \frac{\partial \alpha_0^n}{\partial\theta} + \frac{\mathbb{Z}_{66}}{R} \frac{\partial \xi_R^n}{\partial\theta} + \mathbb{R}_{66} \frac{\partial \beta_0^n}{\partial R} + \mathbb{Z}_{66} \frac{\partial \xi_\theta^n}{\partial R} - \frac{\mathbb{R}_{66}}{R} \beta_0^n - \frac{\mathbb{Z}_{66}}{R} \xi_\theta^n \right) \\ & + \frac{1}{R} \left(\frac{\mathbb{R}_{66}}{R} \frac{\partial \alpha_0^n}{\partial\theta} + \frac{\mathbb{Z}_{66}}{R} \frac{\partial \xi_R^n}{\partial\theta} + \mathbb{R}_{66} \frac{\partial \beta_0^n}{\partial R} + \mathbb{Z}_{66} \frac{\partial \xi_\theta^n}{\partial R} - \frac{\mathbb{R}_{66}}{R} \beta_0^n - \frac{\mathbb{Z}_{66}}{R} \xi_\theta^n \right) \\ & = I_0^n \frac{\partial^2 \beta_0^n}{\partial t^2} + I_1^n \frac{\partial^2 \beta_1^n}{\partial t^2} \end{aligned} \quad (40b)$$

$$\begin{aligned} \delta w : & \frac{\partial}{\partial R} \left((\mathbb{R}_{55}) \left(\xi_R^n + \frac{\partial \chi_0^n}{\partial R} \right) \right) + \frac{1}{R} \frac{\partial}{\partial\theta} \left((\mathbb{R}_{44}) \left(\xi_\theta^n + \frac{\partial \chi_0^n}{R\partial\theta} \right) \right) - \frac{P}{R^2} \frac{\partial^2 \chi_0^n}{\partial\theta^2} \\ & = \left(I_0 \frac{\partial^2 \chi_0}{\partial t^2} \right)^n \end{aligned} \quad (40c)$$

$$\begin{aligned} \delta \xi_R^\eta : & \frac{\partial}{\partial R} \left(\left\{ Z_{11} \frac{\partial \alpha_0^\eta}{\partial R} + C_{11} \frac{\partial \xi_R^\eta}{\partial R} \right\} + \left\{ \frac{Z_{12}}{R} \alpha_0^\eta + \frac{C_{12}}{R} \xi_R^\eta + \frac{Z_{12}}{R} \frac{\partial \beta_0^\eta}{\partial \theta} + \frac{C_{12}}{R} \frac{\partial \xi_\theta^\eta}{\partial \theta} \right\} \right) \\ & - \frac{1}{R} \left(\left\{ Z_{12} \frac{\partial \alpha_0^\eta}{\partial R} + C_{12} \frac{\partial \xi_R^\eta}{\partial R} \right\} + \left\{ \frac{Z_{22}}{R} \alpha_0^\eta + \frac{C_{22}}{R} \xi_R^\eta + \frac{Z_{22}}{R} \frac{\partial \beta_0^\eta}{\partial \theta} + \frac{C_{22}}{R} \frac{\partial \xi_\theta^\eta}{\partial \theta} \right\} \right) \\ & + \frac{1}{R} \frac{\partial}{\partial \theta} \left(\frac{Z_{66}}{R} \frac{\partial \alpha_0^\eta}{\partial \theta} + \frac{C_{66}}{R} \frac{\partial \xi_R^\eta}{\partial \theta} + Z_{66} \frac{\partial \beta_0^\eta}{\partial R} + C_{66} \frac{\partial \xi_\theta^\eta}{\partial R} - \frac{Z_{66}}{R} \beta_0^\eta - \frac{C_{66}}{R} \xi_\theta^\eta \right) \\ & - (\mathbb{R}_{55}) \left(\xi_R^\eta + \frac{\partial \chi_0^\eta}{\partial R} \right) = I_1 \frac{\partial^2 \alpha_0}{\partial t^2} + I_2 \frac{\partial^2 \xi_R}{\partial t^2} \end{aligned} \tag{40d}$$

$$\begin{aligned} \delta \xi_\theta : & \frac{1}{R} \frac{\partial}{\partial \theta} \left(\left\{ Z_{12} \frac{\partial \alpha_0^\eta}{\partial R} + C_{12} \frac{\partial \xi_R^\eta}{\partial R} \right\} + \left\{ \frac{Z_{22}}{R} \alpha_0^\eta + \frac{C_{22}}{R} \xi_R^\eta + \frac{Z_{22}}{R} \frac{\partial \beta_0^\eta}{\partial \theta} + \frac{C_{22}}{R} \frac{\partial \xi_\theta^\eta}{\partial \theta} \right\} \right) \\ & + \frac{\partial}{\partial R} \left(\frac{Z_{66}}{R} \frac{\partial \alpha_0^\eta}{\partial \theta} + \frac{C_{66}}{R} \frac{\partial \xi_R^\eta}{\partial \theta} + Z_{66} \frac{\partial \beta_0^\eta}{\partial R} + C_{66} \frac{\partial \xi_\theta^\eta}{\partial R} - \frac{Z_{66}}{R} \beta_0^\eta - \frac{C_{66}}{R} \xi_\theta^\eta \right) \\ & + \frac{1}{R} \left(\frac{Z_{66}}{R} \frac{\partial \alpha_0^\eta}{\partial \theta} + \frac{C_{66}}{R} \frac{\partial \xi_R^\eta}{\partial \theta} + Z_{66} \frac{\partial \beta_0^\eta}{\partial R} + C_{66} \frac{\partial \xi_\theta^\eta}{\partial R} - \frac{Z_{66}}{R} \beta_0^\eta - \frac{C_{66}}{R} \xi_\theta^\eta \right) \\ & + \frac{1}{R} \left(\frac{Z_{66}}{R} \frac{\partial \alpha_0^\eta}{\partial \theta} + \frac{C_{66}}{R} \frac{\partial \xi_R^\eta}{\partial \theta} + Z_{66} \frac{\partial \beta_0^\eta}{\partial R} + C_{66} \frac{\partial \xi_\theta^\eta}{\partial R} - \frac{Z_{66}}{R} \beta_0^\eta - \frac{C_{66}}{R} \xi_\theta^\eta \right) \\ & - \left((\mathbb{R}_{44}) \left(\xi_\theta^\eta + \frac{\partial \chi_0^\eta}{R \partial \theta} \right) \right) = I_1^\eta \frac{\partial^2 \beta_0}{\partial t^2} + I_2^\eta \frac{\partial^2 \xi_\theta}{\partial t^2} \end{aligned} \tag{40e}$$

where: $\{C_{ij}, Z_{ij}, \mathbb{R}_{ij}\} = \int_{-\frac{h}{2}}^{\frac{h}{2}} \bar{Q}_{ij} \{z^2, z^1, 1\} dz$.

It is highlighted that, due to the equations of compatibility (Eq. (30)), the amounts of variables which are unknown would be declined from by 8. So, the total amounts of unknowns in the core and face sheets is reduced from 25 to 13.

9 Solution process

Recently, computer-aid simulation and semi numerical approach become well-known to simulation different problems [94–97]. In this stage of the current paper, we introduce a solution procedure called GDQM to solve the formulation of the current problem. In this method have:

$$\frac{\partial f}{\partial r} \Big|_{r=r_i, \theta=\theta_j} = \sum_{m=1}^{N_r} \sum_{n=1}^{N_\theta} A_{im}^r I_{jn}^\theta f_{mn} \tag{41a}$$

$$\frac{\partial f}{\partial \theta} \Big|_{r=r_i, \theta=\theta_j} = \sum_{m=1}^{N_r} \sum_{n=1}^{N_\theta} I_{im}^r A_{jn}^\theta f_{mn} \tag{41b}$$

$$\frac{\partial}{\partial r} \left(\frac{\partial f}{\partial \theta} \Big|_{r=r_i, \theta=\theta_j} \right) = \sum_{m=1}^{N_r} \sum_{n=1}^{N_\theta} A_{im}^r A_{jn}^\theta f_{mn} \tag{41c}$$

$$\frac{\partial^2 f}{\partial r^2} \Big|_{r=r_i, \theta=\theta_j} = \sum_{m=1}^{N_r} \sum_{n=1}^{N_\theta} B_{im}^r I_{jn}^\theta f_{mn} \tag{41d}$$

$$\frac{\partial^2 f}{\partial \theta^2} \Big|_{r=r_i, \theta=\theta_j} = \sum_{m=1}^{N_r} \sum_{n=1}^{N_\theta} I_{im}^r B_{jn}^\theta f_{mn} \tag{41e}$$

Also, I_{im}^r and I_{jn}^θ would equal to 1 while $m = i$ and $n = j$ otherwise, may be equal to 0. Furthermore, $B_{jn}^\theta, B_{im}^r, A_{jn}^\theta$ and A_{im}^r may be weighting factors of the second and first-order derivatives along with the θ and r orientation, respectively, and would be taken into consideration as

$$A_{im}^{(1)} = \begin{cases} \frac{\xi(r_i)}{(r_i - r_m)\xi(r_m)} & \text{when } i \neq m \\ -\sum_{k=1, k \neq i}^{N_r} A_{ik}^{(1)} & \text{when } i = m \end{cases} \quad i, m = 1, 2, \dots, N_r \quad (42a)$$

$$B_{im}^{(r)} = r \left(A_{ii}^{(r-1)} A_{im}^{(1)} - \frac{A_{im}^{(r-1)}}{(r_i - r_m)} \right) \quad i, m = 1, 2, \dots, N_r, \quad i \neq m \quad (44a)$$

$$A_{jn}^{(1)} = \begin{cases} \frac{\xi(\theta_j)}{(\theta_j - \theta_n)\xi(\theta_n)} & \text{when } j \neq n \\ -\sum_{k=1, k \neq j}^{N_\theta} A_{jk}^{(1)} & \text{when } j = n \end{cases} \quad j, n = 1, 2, \dots, N_\theta \quad (42b)$$

$$B_{jn}^{(\theta-1)} = \theta \left(A_{jj}^{(\theta-1)} A_{jn}^{(1)} - \frac{A_{jn}^{(\theta-1)}}{(\theta_j - \theta_n)} \right) \quad j, n = 1, 2, \dots, N_\theta, \quad j \neq n \quad (44b)$$

$$B_{ii}^{(r)} = -\sum_{k=1, k \neq i}^{N_r} B_{ik}^{(r)}, \quad i = 1, 2, \dots, N_r, \quad i = m \quad (44c)$$

$$B_{jj}^{(\theta)} = -\sum_{k=1, k \neq j}^{N_\theta} B_{jk}^{(\theta)}, \quad j = 1, 2, \dots, N_\theta, \quad j = n \quad (44d)$$

in which

$$\xi(r_i) = \prod_{k=1, k \neq i}^{N_r} (r_i - r_k) \quad (43a)$$

$$\xi(\theta_j) = \prod_{k=1, k \neq j}^{N_\theta} (\theta_j - \theta_k) \quad (43b)$$

In addition, applying points of Chebyshev polynomials greed, the seed as well as r and θ directions would be distributed as [103]

And

Table 1 Comparison of non-dimensional natural frequency ($\bar{\omega}^2 = \omega b^2 \sqrt{\rho/E}$) of the multi-scale nanocomposite face sheets for different CNT pattern

CNT distribution	N_{CNT}		0.06	
	Current research	Ref. [108]	Current research	Ref. [108]
FG-X	3.2647	3.2604	3.4805	3.4808
FG-V	3.3775	3.3817	3.5795	3.5783
FG-A	3.4275	3.4291	3.6292	3.6294
FG-UL	3.4174	3.4156	3.6197	3.6170

Table 2 The material properties of core [18] and composite layer [109]

CNTs	Matrix (Epoxy)	Fiber (Carbon)	Metal (Core)
$\nu_{12} = 0.33$	$E^m (Gpa) = 3.51$	$\nu^f = 0.2$	$\nu^s = 0.34$
$\rho^{cnt} (kg/m^3) = 1350$	$\nu^m = 0.34$	$\rho^f (kg/m^3) = 1750$	$\rho^s (kg/m^3) = 2700$
$E^{cnt} (Gpa) = 640$	$\rho^m (kg/m^3) = 1200$	$E_{11}^f (GPa) = 233.05$	$E^s (Gpa) = 70$
$d^{cnt} (m) = 1.4 \times 10^{-9}$		$E_{11}^f (GPa) = 23.1$	
$t^{cnt} (m) = 0.34 \times 10^{-9}$		$G_{11}^f (GPa) = 8.96$	
$l^{cnt} (m) = 25 \times 10^{-6}$			

Table 3 Convergence study of the GDQM for solving the current research for different boundary conditions (B.Cs)

B.Cs	$N=4$	$N=6$	$N=8$	$N=10$	$N=12$	$N=14$
S-S	0.0370	0.0419	0.0421	0.0422	0.0422	0.0422
C-S	0.1015	0.1009	0.1009	0.1009	0.1009	0.1009
C-C	0.1398	0.1390	0.1385	0.1385	0.1385	0.1385

Table 4 Convergence study of the GDQM for solving the current research for different CNT-distribution pattern and PD

PD	CNT- distribution	Nt=2	Nt=4	Nt=6	Nt=8	Nt=10	Nt=12	Nt=14	Nt=16	Nt=100
X	FG-X	0.1385	0.1363	0.1370	0.1373	0.1374	0.1374	0.1374	0.1375	0.1375
	FG-O	0.1411	0.1354	0.1349	0.1347	0.1347	0.1346	0.1346	0.1346	0.1346
	FG-UD	0.1401	0.1358	0.1358	0.1358	0.1357	0.1357	0.1357	0.1357	0.1358
O	FG-X	0.1247	0.1360	0.1363	0.1363	0.1363	0.1363	0.1363	0.1363	0.1363
	FG-O	0.1286	0.1344	0.1346	0.1346	0.1346	0.1346	0.1346	0.1346	0.1347
	FG-UD	0.1275	0.1349	0.1351	0.1351	0.1351	0.1351	0.1351	0.1351	0.1351
UD	FG-X	0.1325	0.1352	0.1358	0.1359	0.1360	0.1360	0.1361	0.1361	0.1361
	FG-O	0.1357	0.1342	0.1338	0.1338	0.1337	0.1337	0.1337	0.1337	0.1337
	FG-UD	0.1346	0.1346	0.1346	0.1346	0.1346	0.1346	0.1346	0.1346	0.1346

Table 5 Effects of three types’ method for reinforcing the structure on the system’s frequency with consideration three porosity coefficient and boundary conditions

	Simply-Simply			Clamped-Simply			Clamped-Clamped		
	FG-X	FG-O	FG-UD	FG-X	FG-O	FG-UD	FG-X	FG-O	FG-UD
Without imperfection ($e_0=0$)									
MHC/HC/MHC ¹	0.0388	0.0378	0.0381	0.0918	0.0918	0.0926	0.1296	0.1263	0.1275
CNT/HC/CNT ²	0.0489	0.0395	0.0465	0.0825	0.0697	0.0785	0.1190	0.9998	0.1155
With imperfection ($e_0=0.3$)									
MHC/HC/MHC ¹	0.0396	0.0386	0.0390	0.0961	0.0938	0.0946	0.1322	0.1291	0.1302
CNT/HC/CNT ²	0.0451	0.0321	0.0407	0.0875	0.0594	0.0702	0.1140	0.0895	0.1042
With imperfection ($e_0=0.5$)									
MHC/HC/MHC ¹	0.0406	0.0396	0.0400	0.0984	0.0963	0.0971	0.1355	0.1325	0.1337
CNT/HC/CNT ²	0.0556	0.0407	0.0507	0.0950	0.0732	0.0842	0.1329	0.1042	0.1258

¹Multi-scale hybrid nanocomposite reinforced disk /Honey-Comb/ Multi-scale hybrid nanocomposite reinforced disk
²Carbon nanotubes reinforced disk/Honey-Comb/ Carbon nanotubes reinforced disk

$$r_i = \frac{-R_i + R_0}{2} \left(-\cos \left(\frac{(i-1)}{(N_r-1)} \pi + 1 \right) \right) + R_i \quad i = 1, 2, 3, \dots, N_r \tag{45a}$$

$$\theta_j = \left(-\cos \left(\frac{(j-1)}{(N_\theta-1)} \pi + 1 \right) \right) \frac{\chi}{2} \quad j = 1, 2, 3, \dots, N_\theta \tag{45b}$$

By employing the GDQM as solution method of this problem, an algebraic eigenvalue [79, 104–107] would be created. These relations would be given by:

$$\left\{ \begin{bmatrix} M_{aa} & M_{ab} \\ M_{ba} & M_{bb} \end{bmatrix} \right\} \omega_n^2 + \left\{ \begin{bmatrix} K_{dd} & K_{db} \\ K_{bd} & K_{bb} \end{bmatrix} \right\} \begin{Bmatrix} \delta_d \\ \delta_b \end{Bmatrix} = 0 \tag{46}$$

Eventually, Though Eq. (46) the novel system are

$$[K_{dd}] \delta_d + [K_{db}] \delta_b = 0 \tag{47a}$$

$$[K_{bd}] \delta_d + [K_{bb}] \delta_b = 0 \tag{47b}$$

here, the freedom degrees’ vector will be expressed as:

$$\delta_b = -\delta_d \frac{[K_{dd}]}{[K_{db}]} \tag{48}$$

By inserting of Eq. (48) into Eq. (47b):

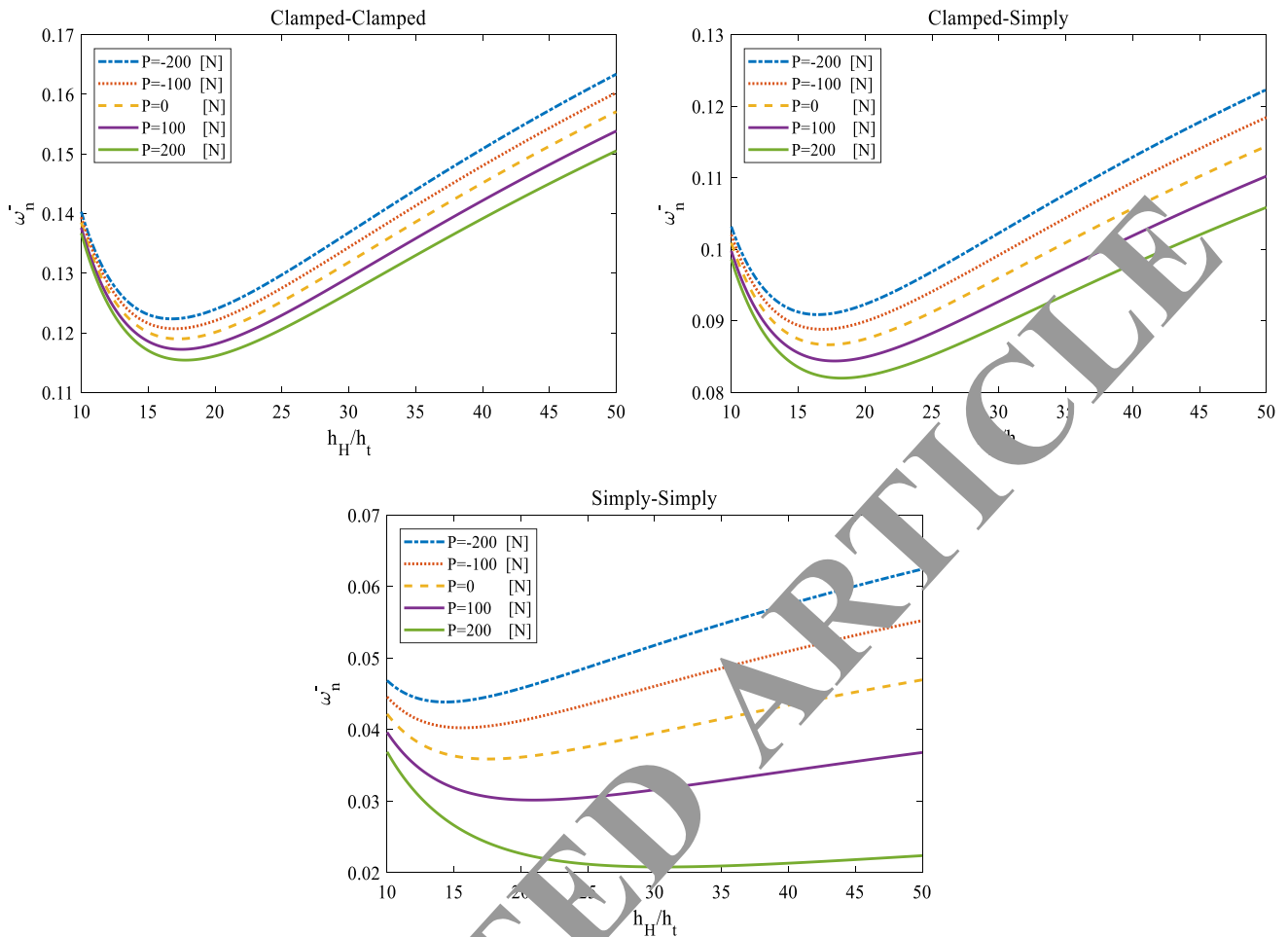


Fig. 4 Frequency of the disk versus to the honeycomb core to composite face-sheet thickness ratio for five value of internal force and three kinds of boundary conditions

$$\left(-[K_{bb}][K_{db}]^{-1}[K_{dd}] + [K_{bd}]\right)\delta_d = 0 \tag{49}$$

So

$$K^* = [K_{bd}] - [K_{bb}][K_{db}]^{-1}[K_{dd}] \tag{50}$$

and

$$[M_{bd}]\delta_d + [M_{bb}]\delta_b = 0 \tag{51}$$

Also by applying of Eq. (48) into Eq. (51):

$$\left(-[M_{bb}][K_{db}]^{-1}[K_{dd}] + [M_{bd}]\right)\delta_d = 0 \tag{52}$$

Thereby

$$M^* = -[M_{bb}][K_{db}]^{-1}[K_{dd}] + [M_{bd}] \tag{53}$$

Ultimately, through solving the following relation, the structure's displacement fields and frequency information would be derived by GDQM.

$$M^*\omega^2 + K^* = 0 \tag{54}$$

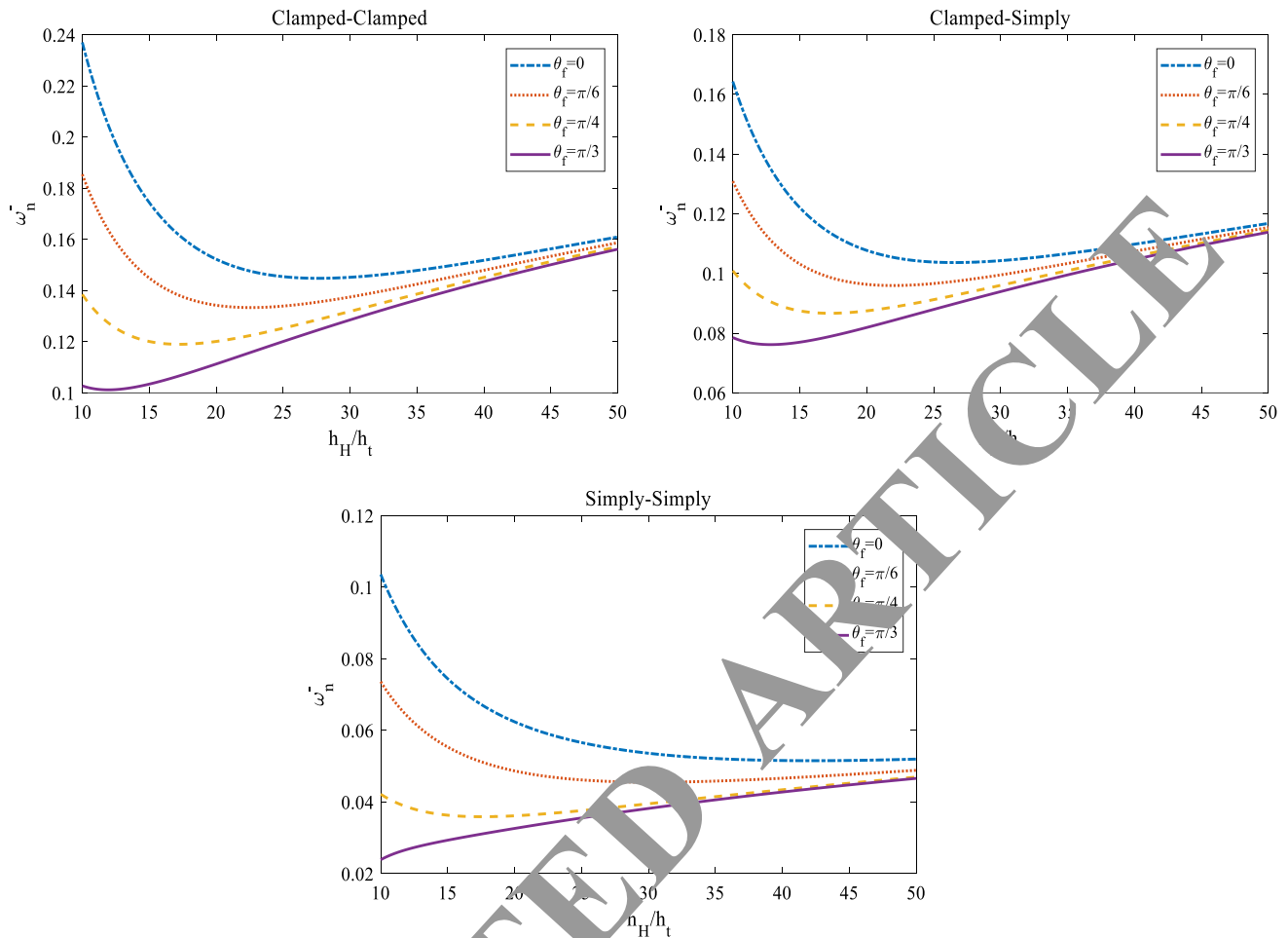


Fig. 5 Frequency of the disk versus to the honeycomb core to composite face-sheet thickness ratio for three kinds of boundary conditions and four θ_f

10 Results and discussion

For evaluating the reliability of the current method, the dimensionless natural frequency of the annular plate obtained in this research is compared with outcomes provided by Ref. [108] in Table 1. As illustrated in a broad range of W_{CNT} and CNT-distribution pattern, there is a good agreement between the outcomes. Regarding the table, this study anticipates the vibration characteristics of the annular plate very precise and similar to the outcomes given in Ref. [106].

The mechanical properties for both of the reinforcements (Carbon Fiber and Carbon nanotube) and the epoxy matrix are shown in Table 2.

Also, the properties of SMA and matrix layer illustrated in Ref. [48]. Convergence number of grid points for having independent results with respect to the three kinds of boundary conditions is investigated in Table 3. As illustrated in Table 3, as grid points' number in the GDQ approach is more than eleven, the error for the calculation of the disk's natural frequency becomes zero and this matter is a fact for all boundary conditions.

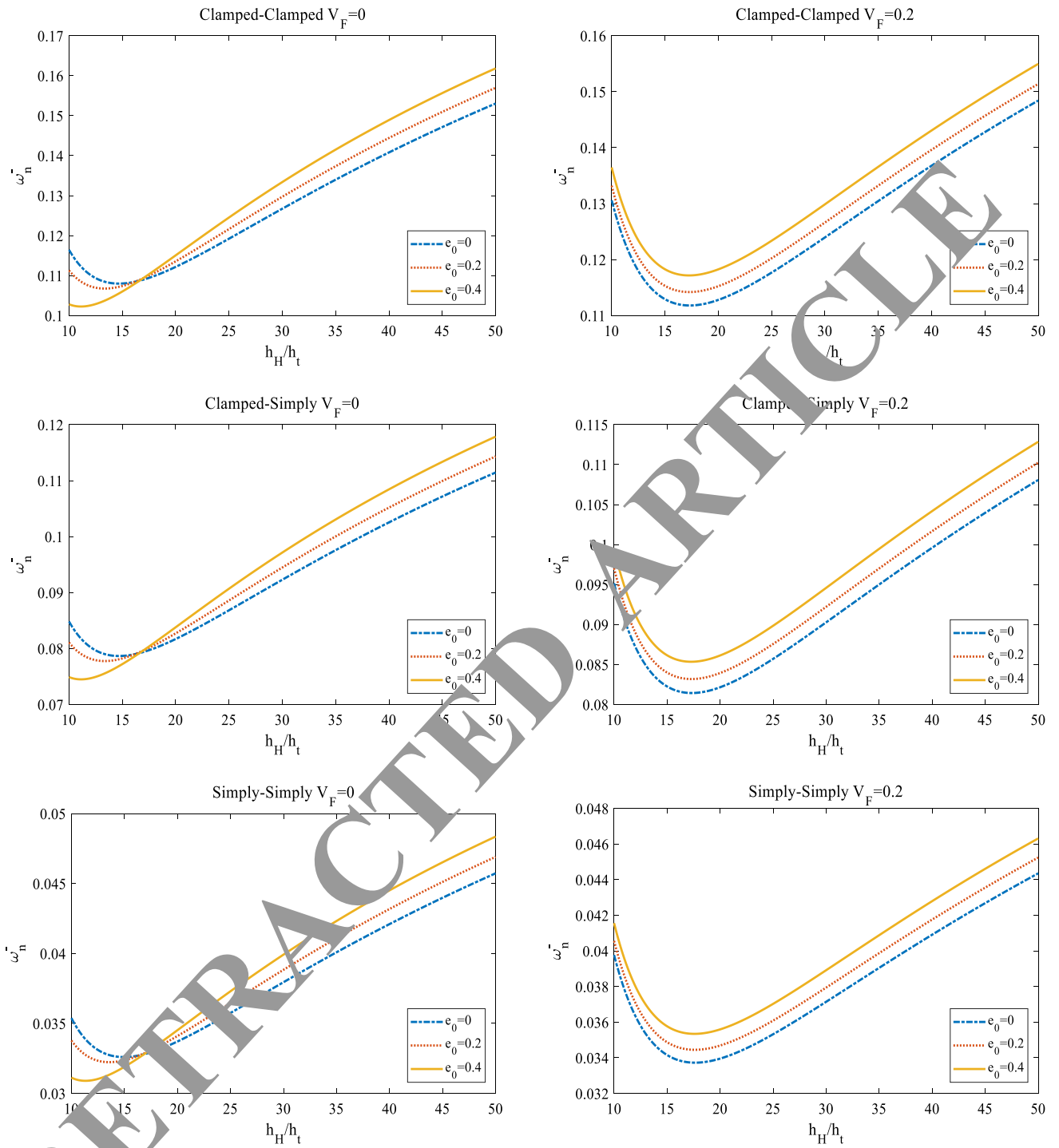


Fig. 6 Frequency of the disk versus to h_H/h_t for three kinds of boundary conditions

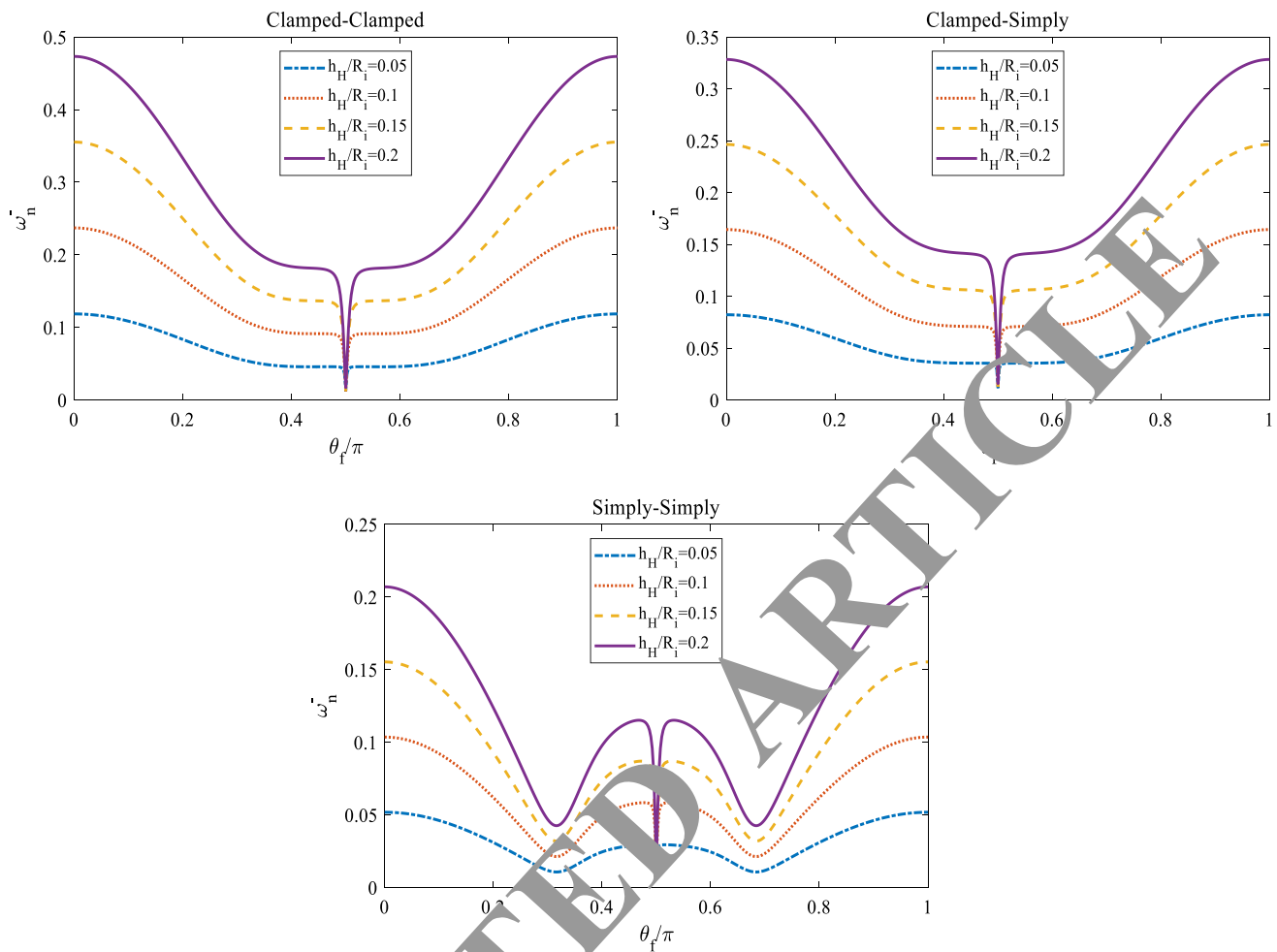


Fig. 7 Frequency of the disk versus to θ_f/π for three kinds of boundary conditions and four h_H/R_i

Convergence information for effect of the number of nanocomposite layer (N_l) and different functionally graded distribution on the frequency of the multi-phase disk with respect to the porosity patterns is depicted in Table 4.

Based on Table 4, we can state that the structure will have the best dynamic response by considering the PD-X and FG-O patterns. Also the layer's number in the compositionally face sheets must not exceed nine for all FG and porosity patterns, due to the $N_l \geq 10$ we are not able to observe any change in the structure's frequency. Analysis of the impact of three types' method for reinforcing the structure on the system's frequency with consideration three porosity coefficient and boundary conditions is argued in

Table 5. The ends of Table 5 are that for clamped-simply, and clamped-clamped not only MHC/HC/MHC reinforced disk have the greatest natural frequency compared to CNTs/HC/CNTs reinforcements but also growing the imperfection effect is a reason to decrease the systems' frequency.

In accord with Table 5, it can be concluded that applying the honeycomb network as the core of the structure will enhance the structure's dynamic response, impressively. The graphs and data in Fig. 4 depicts the internal force, three types of boundary conditions, and honeycomb core to composite face-sheet thickness ratio (h_H/h_f) effects on the vibrational response of the compositionally disk. It is as a fact

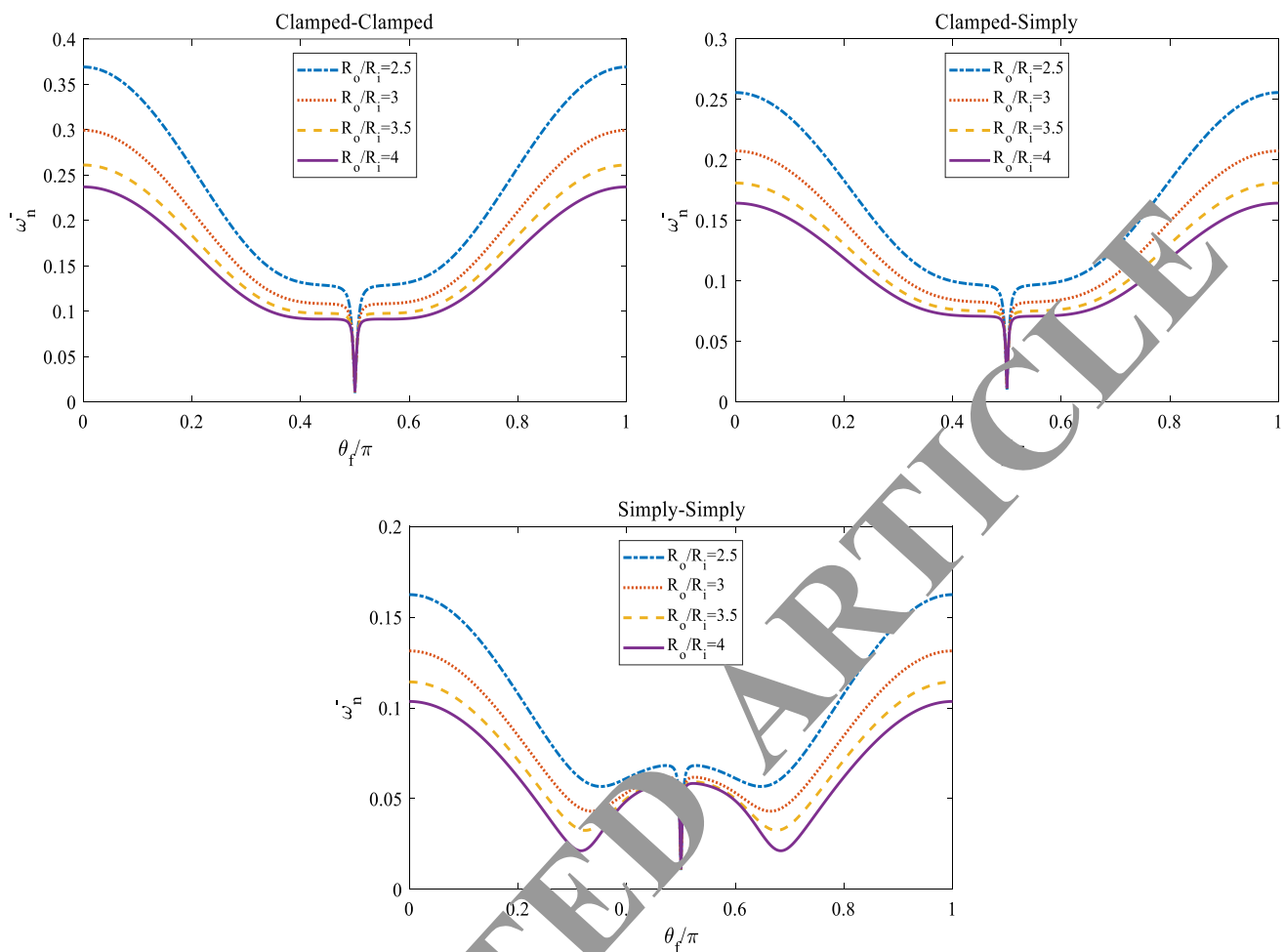


Fig. 8 Frequency of the disk versus to $\frac{\theta_f}{\pi}$ for three kinds of boundary conditions and four R_o/R_i

for various boundary condition in which with each increase in the applied internal force can claim that the frequency of the structure tends to improve, particularly in the bigger value of h_H/h_t and simply-simply boundary conditions. For each boundary condition and any value of internal pressure, the relation between h_H/h_t and the sandwich structure's frequency is the same as quadratic function. For more explanation by ever-increasing in the h_H/h_t at first, the honeycomb core structure's frequency drops, exponentially and after a specific value for the h_H/h_t the structure's dynamic response improves, linearly.

The Fig. 8 illustrates the fibers angel (θ_f), three types of boundary conditions, and honeycomb core to composite face-sheet thickness ratio (h_H/h_t) effects on the vibrational behaviors

of the compositionally disk. It is as a fact for various boundary condition in which with each decrease in the θ_f can claim that the frequency of the structure tends to improve, particularly for the smaller values of h_H/h_t and simply-simply boundary conditions. Also, the impact of the fibers angle on the frequency can be overlooked for large h_H/h_t . For each boundary conditions and any value of fibers angel, the relation among h_H/h_t and the sandwich structure's frequency is the same as quadratic function. To clarify it more, by ever-increasing in the h_H/h_t at first, the honeycomb core structure's frequency drops, exponentially and after the specific value of the structure's dynamic response improves, linearly. It is important to mention that the specific value for the h_H/h_t grows by increasing the θ_f .

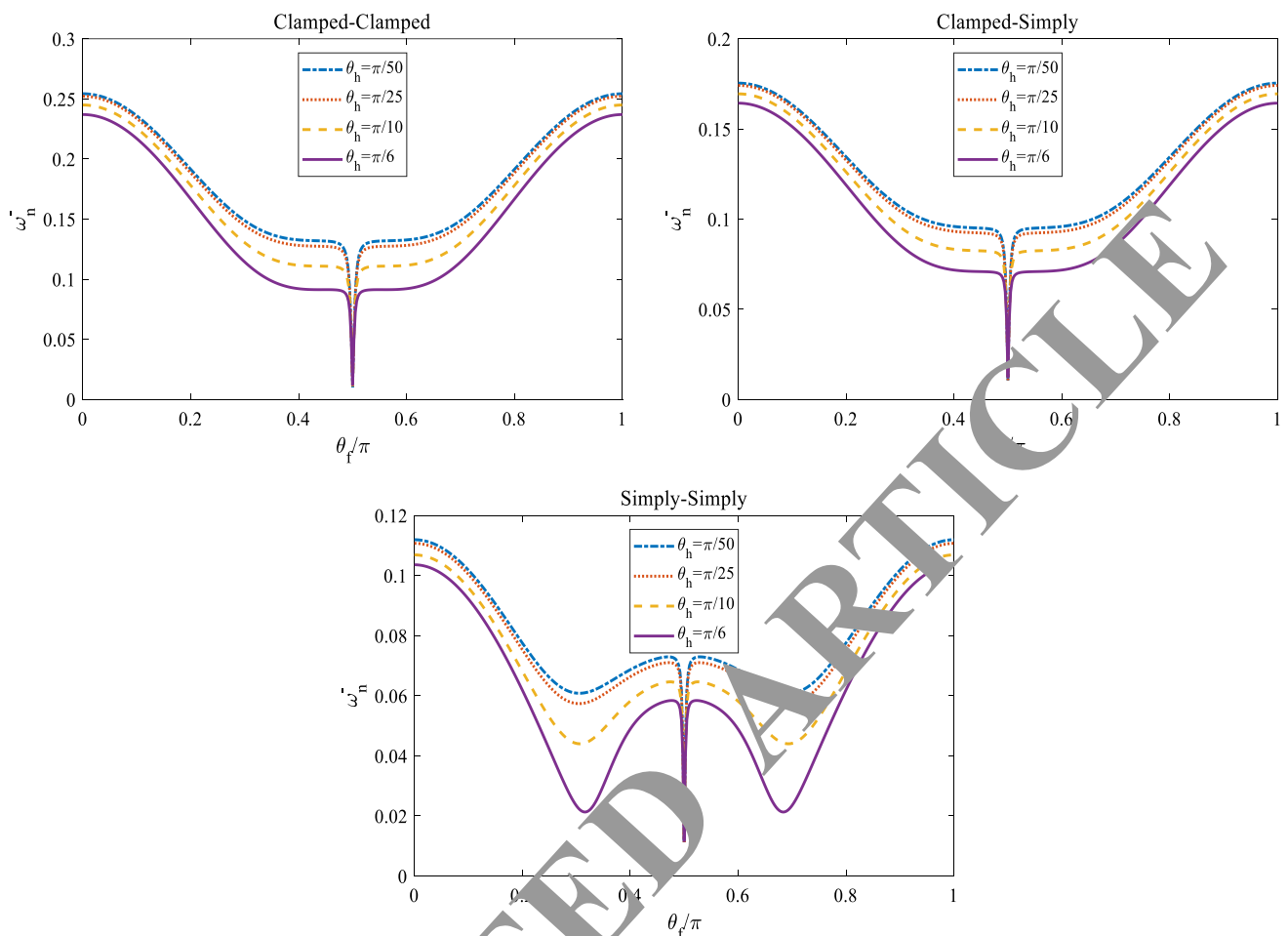


Fig. 9 Frequency of the disk versus to $\frac{\theta_f}{\pi}$ for three kinds of boundary conditions and four θ_h

The depicted graph in Fig. 9 displays the fibers angel (V_f), three types of boundary conditions, porosity factor, and honeycomb core to composite thickness ratio of face sheet (h_H/h_i) effects on the vibrational behaviors of the compositionally disk. As common result in all below figure by ever-increasing in the h_H/h_i at first, the honeycomb core structure's frequency drops exponentially and after the specific value the structure's dynamic response improves, linearly.

Figure 9 shows that the fibers angel (θ_f/π), three types of boundary conditions, and honeycomb core thickness to inner radius ratio (h_H/R_i) effects on the vibrational behaviors of the compositionally disk. As Fig. 7 presents the relation between fibers angle and frequency is similar to a

bell-shaped function with the positive concave. When the θ_f/π becomes close to 0.5, there is no change in the disk's frequency, especially for the smaller h_H/R_i and C–C and C–S boundary conditions. In addition, for each value of θ_f/π , dynamic stability of the structure will improve by increasing h_H/R_i . As another general outcome, the critical or minimum frequency for the sandwich disk displays when the $(\theta_f/\pi)^{C-C}$ and $(\theta_f/\pi)^{C-S}$ are 0.5 while for S–S edges, θ_f/π is equal to 0.5, 0.32, and 0.64.

The given information in Fig. 8 shows the fibers angel (θ_f/π), three kinds of boundary conditions, and outer to inner radius ratio (R_o/R_i) affect the vibrational behavior of the compositionally disk. Figure 8 presents that each value

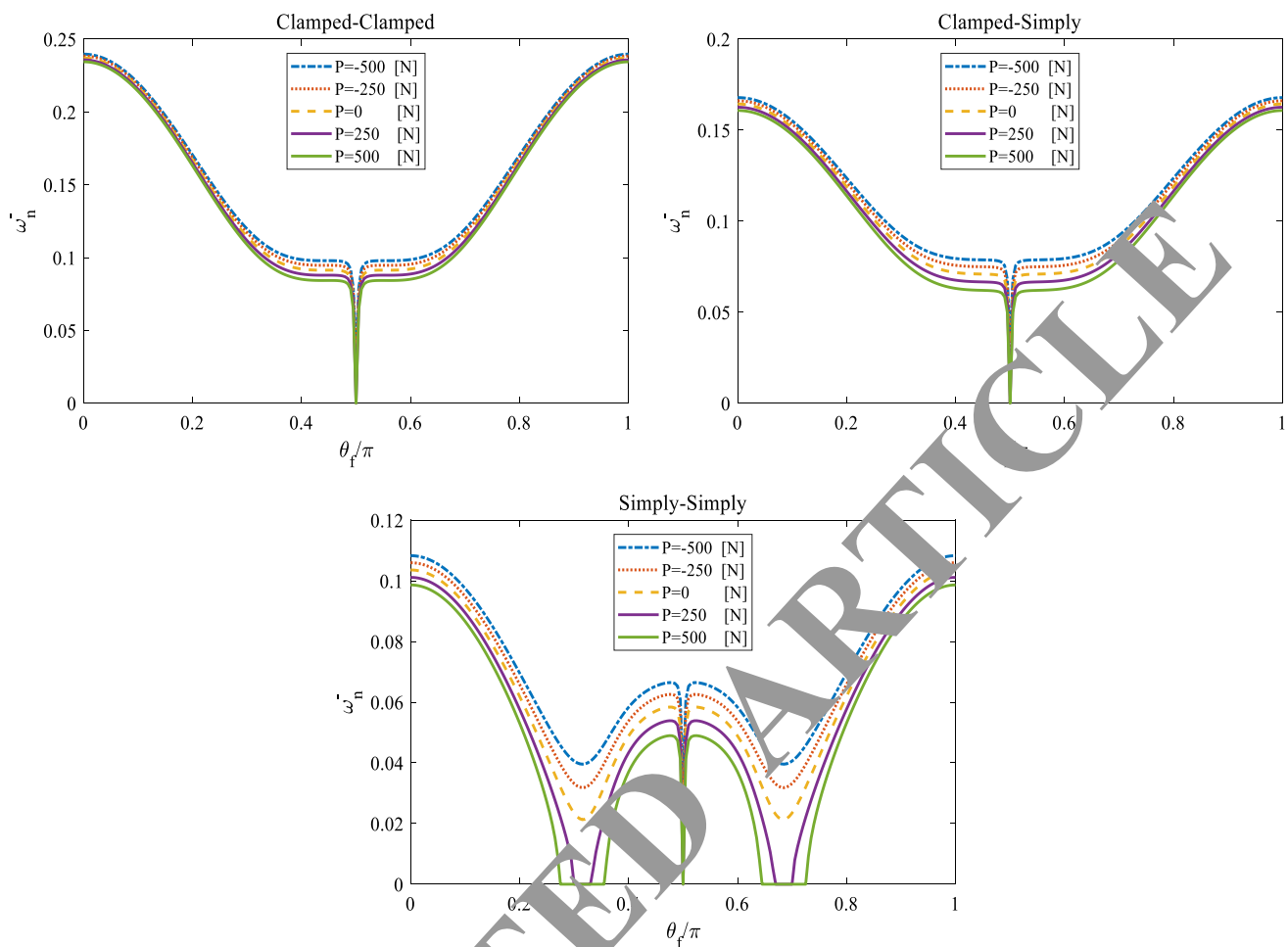


Fig. 10 Frequency of the disk versus to θ_f/π for three kinds of boundary conditions and four p

of θ_f/π , the structure's dynamic stability will improve through an increase in R_f parameter.

The given information in Fig. 9 shows the fibers angle (θ_f/π), three types of boundary conditions, and honeycomb network angle (θ_h) affect the vibrational behaviors of the compositionally disk. As Fig. 9 depicts by considering the close angles for honeycomb network, θ_h have significant role in the disk's frequency as an enhancement.

The Fig. 10 indicates that the fibers angle (θ_f/π), three kinds of boundary conditions, and external applied load (p) effects on the vibrational behaviors of the compositionally disk. By considering external applied load, the impact

of compressive or tensile load on the structure's dynamic response is impressive as the rigidity of the edges decreases and the dimensionless fibers angle becomes close to 0.5. As a remarkable result, when the structure is encountered with clamped edges, there is a critical fibers angle in which is equal to 0.5, but if we consider pure simply edges and compressive load, there can see three critical fibers angles in which are equal to 0.5, 0.33, and 0.66. In addition, by considering tensile applied load there is a range for critical fibers angle and this range expands by increasing the value of applied load.

Figure 11 depicts the effects of three types of boundary conditions, various thickness of the SMA reinforced face

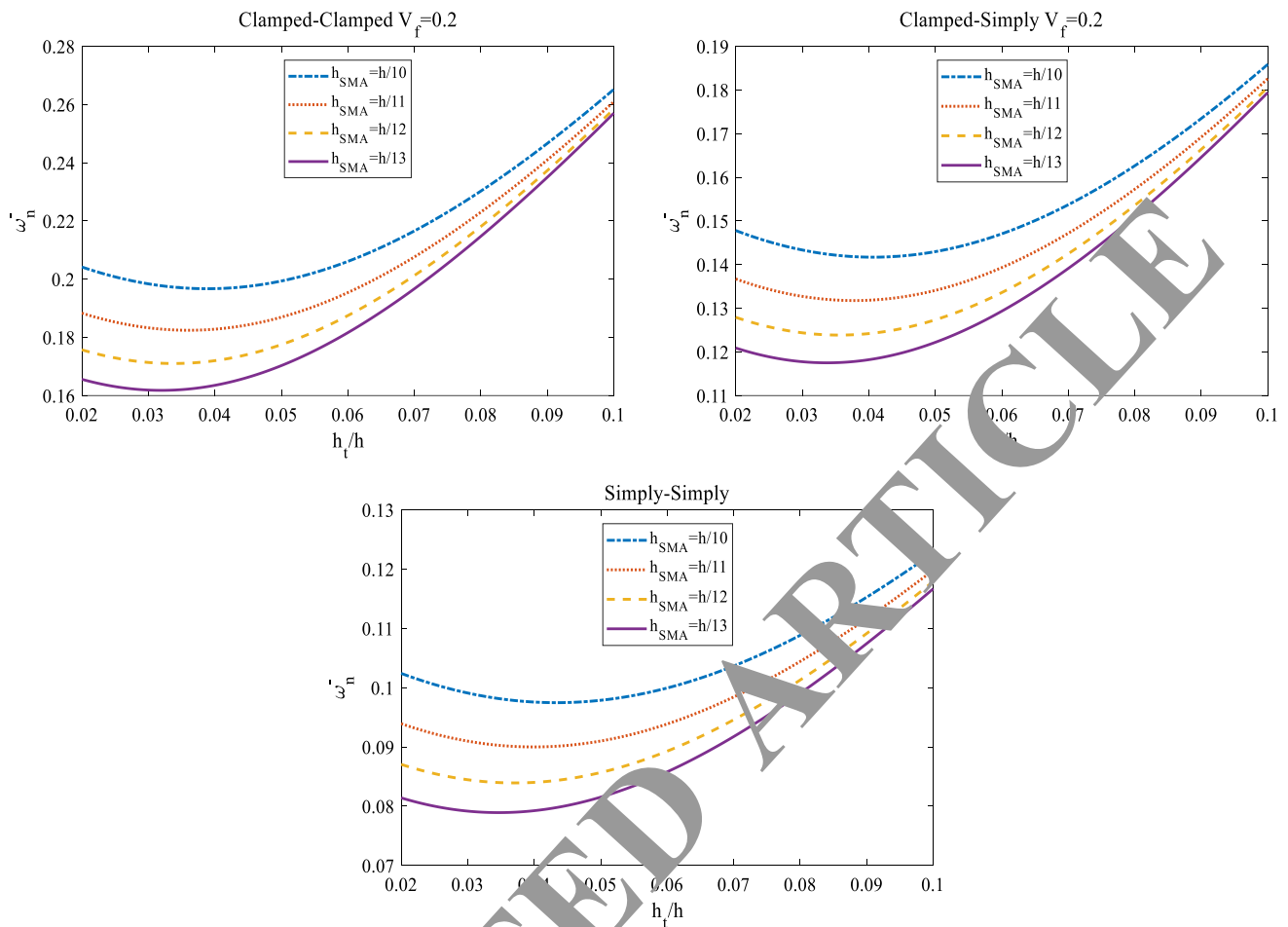


Fig. 11 Frequency of the sandwich disk versus to h_t/h for three kinds of boundary conditions and four h_{SMA}

sheet (h_{SMA}), and MHC face sheet to total thickness ratio (h_t/h) on the vibrational response of the compositionally disk. For various boundary condition, with each increase in the h_{SMA} parameter the frequency of the structure tends to improve, particularly in the lowest value of h_t/h and simply-simply boundary condition. Also, for boundary conditions and h_{SMA} values, the relation between h_t/h and sandwich structure's frequency is the same as quadratic function. To clarify it, by ever-increasing in the h_t/h at first the honeycomb core structure's frequency drops, exponentially and

after an specific value for the h_t/h the dynamic response of the structure improves.

In Fig. 12 shows the effects of three boundary condition types, various SMA (V_S) value fraction, and MHC face sheet to total thickness ratio (h_t/h) on the vibrational response of the compositionally disk. For various boundary condition, with each increase in the V_S parameter the frequency of the structure tends to improve. As an important report, the impact of V_S element on the dynamics of the structure is more considerable at the initial h_t/h value.

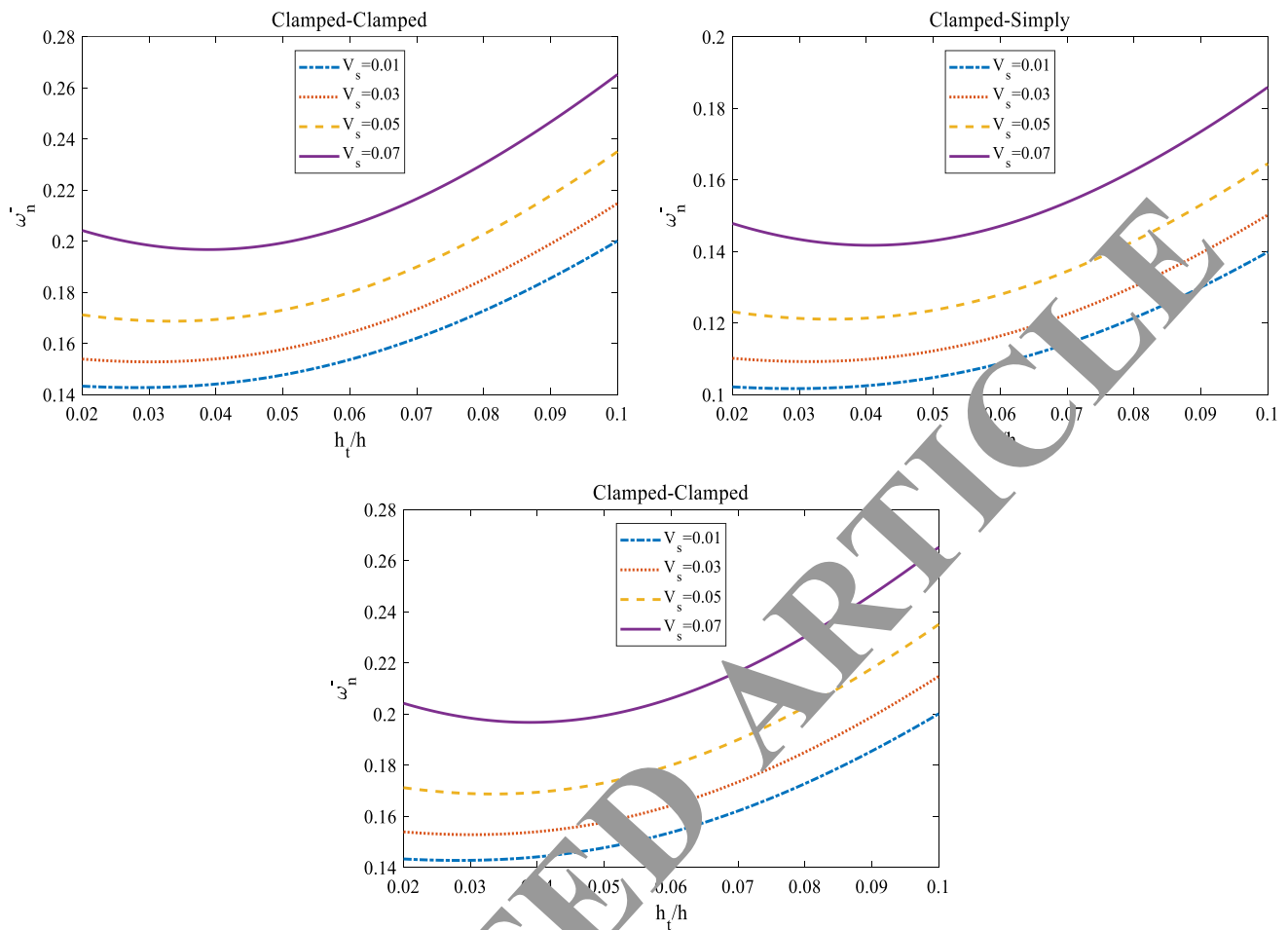


Fig. 12 Frequency of the sandwich disk versus h_t/h parameter for three kinds of boundary conditions and four V_s

11 Conclusion

A porous sandwich disk's vibrational specifications with a core (honeycomb), two middle layers including SMA fiber, and two outer layers of MR under in-plane pressure is studied. By employing FEM, the stress and strain relations are acquired. Halpin's modified model–Tsai and the mixture rule are combined to provide the efficient material constant of the introduced porous sandwich structure. Eventually, the key results of this study are as follows:

- for each boundary conditions and internal pressure's values, the relation among sandwich structure' frequency and h_H/h_t is as same as quadratic function.
- As an important report, the impact of V_s element on the dynamics of the structure is more considerable at the initial h_t/h value.
- For boundary conditions and h_{SMA} values, the relation between sandwich structure's frequency and h_t/h is as same as quadratic function.
- for various boundary conditions, with each increase in the h_{SMA} parameter the frequency of the structure tends to improve, particularly in the lowest value of h_t/h and simply-simply boundary conditions.
- the effect of the fibers angle on the frequency can be overlooked for lager value of h_H/h_t
- For various boundary condition, with each increase in the h_{SMA} parameter the frequency of the structure tends to improve, particularly in the lowest value of h_t/h and simply-simply boundary conditions
- By an increase in the h_H/h_t at first, the honeycomb core structure's frequency drops, exponentially and after the specific value of the structure's dynamic response improves, linearly.
- when the θ_f/π parameter becomes close to $\pi/2$, there is no change in the disk's frequency, especially for the smaller value of h_H/R_i and C–C boundary conditions.
- by considering external applied load, the impact of compressive or tensile load on the structure's dynamic response is impressive as the edges' rigidity decreases and the dimensionless fibers angle becomes close to $\pi/2$.

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