

Microscopic Polarization in Bilayer Graphene

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I. Numerical Data Analysis of Landau Levels

The low energy band structure of graphene consists of four bands; two bands meet at the charge neutrality point, and two higher energy bands related to the A_2 - B_1 dimer bonds are offset by the interlayer energy γ_1 . We adopt a two-band model, which includes only the lowest energy bands, to quantitatively analyze the LLs spectra^{2,3}. The analysis of LLs is limited to the low energy regime (less than 150 meV), which meets the requirement of the two-band model; $E < \gamma_1 = 0.377$ eV. We have found that this simple two-band model gives results within 5 % of those from the full four-band model⁴. Thus, the introduced error for using the simpler model is less than our experimental uncertainties. The simple analytic expression for the LL energies is derived from the two-band model as given by Eq. 1 in the main text, which we use to fit the LLs spectra with the potential energy asymmetry, ΔU and Fermi velocity, v_F as free parameters.

Figure S1a shows the individual dI/dV spectrum obtained for the P1 electron puddle at $V_g = 33$ V. Up to nine Landau levels are observed over the displayed spectra range in Fig. S1a. A fit of the LL energies with Eq. 1 yields an energy asymmetry, $\Delta U = (-34.8 \pm 2.4)$ meV⁵, and Fermi velocity of $v_F = (1.00 \pm 0.01) \times 10^6$ m s⁻¹ as shown in Fig. 3d in the main text.

The band gaps qualitatively follow the separation between $LL_{(0,+)(1,+)}$ and $LL_{(2,+)(2,-)}$. We plot in Fig. S1b the difference in LL energies between $LL_{(2,+)(2,-)}$ - $LL_{(0,+)(1,+)}$, which can be obtained directly from the gate maps (Fig. 3c). It is clear that the determined energy asymmetries from the fitting of the LLs with Eq. 1 (Fig. S1c, d) match the overall trend seen in the difference between LLs. Additionally, we have confirmed that using the Fermi velocity as a variable fit parameter only weakly affects the gap determination as demonstrated in Fig. S1c, d. Using the velocity as a fitting parameter, however, we observe a few percent variation in Fermi velocity as displayed in Fig. S1e, which can be ascribed to the velocity renormalization effects due to electron-electron interactions⁶.

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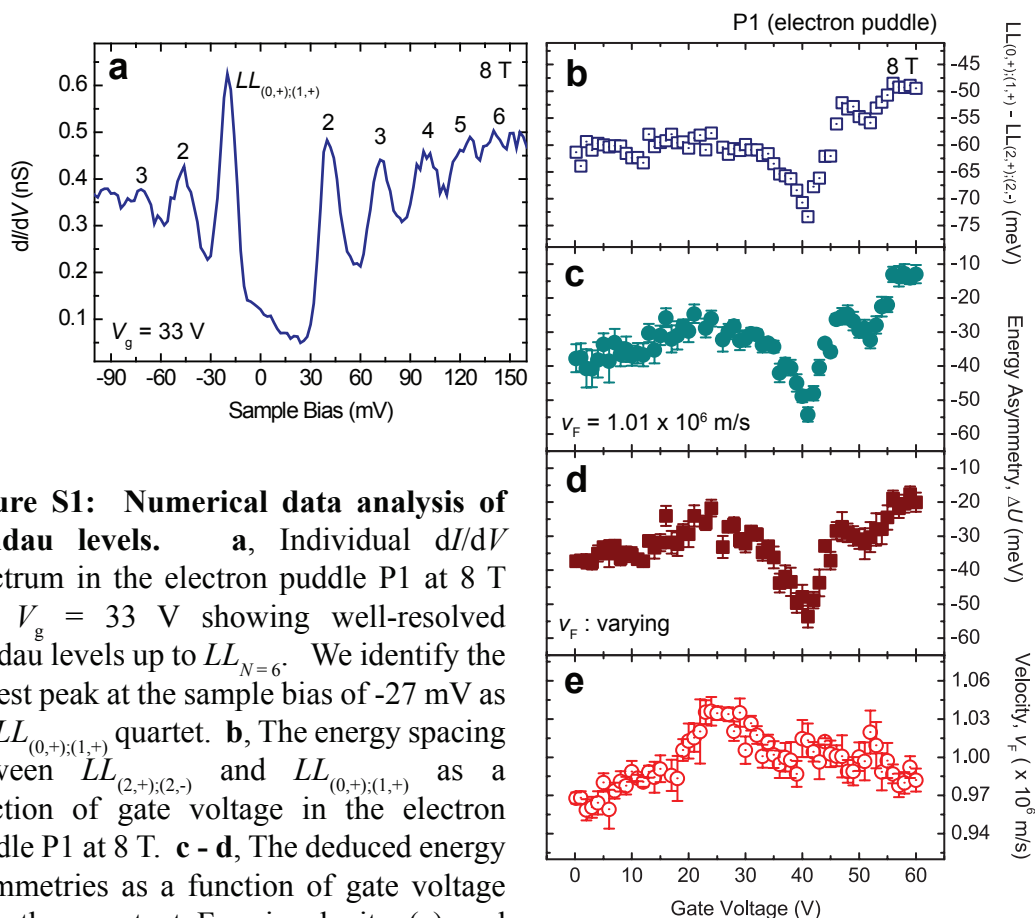


Figure S1: Numerical data analysis of Landau levels. **a**, Individual dI/dV spectrum in the electron puddle P1 at 8 T and $V_g = 33$ V showing well-resolved Landau levels up to $LL_{N=6}$. We identify the largest peak at the sample bias of -27 mV as the $LL_{(0,+);(1,+)}$ quartet. **b**, The energy spacing between $LL_{(2,+);(2,-)}$ and $LL_{(0,+);(1,+)}$ as a function of gate voltage in the electron puddle P1 at 8 T. **c - d**, The deduced energy asymmetries as a function of gate voltage with the constant Fermi velocity (**c**) and with the Fermi velocity as a fitting parameter (**d**). The band gaps follow the evolution of LL spacing displayed in **b**. **e**, The extracted Fermi velocity variation as a function of gate voltage from the LL fitting. The velocity renormalization, manifested by a few percent variations in magnitude does not affect the main features of bilayer energy asymmetry (Fig. S1c, d). Error bars one standard deviation, in (c) - (e).