

# Coincidence of a high-fluence blazar outburst with a PeV-energy neutrino event

## 3 1 Supplementary Methods

4 **VLBI Imaging:** The TANAMI VLBI images of PKS B1424–418 (Fig. 1) are based on observa-  
5 tions on 2011 Nov 13, 2012 Sep 16, and 2013 Mar 14 with the Long Baseline Array (LBA) and  
6 associated telescopes in Australia, New Zealand, South Africa, Chile, and Antarctica. Correlation  
7 of the data was performed with the DiFX correlator<sup>46</sup> at Curtin University in Perth, Western Aus-  
8 tralia. The data were calibrated and imaged following Ojha et al. (2010)<sup>17</sup>. Supplementary Table 1  
9 gives an overview of the array configuration and image parameters. The noise level for each image  
10 was determined by fitting a Gaussian function to an off-source part of the image. For this purpose  
11 and for producing the final images shown here we used the software package ISIS 1.6.2<sup>47</sup>.

12 The TANAMI array varies in configuration for each observation, leading to inhomogeneous  
13 sampling of the visibility data. This results in a different noise level and restoring beam in the  
14 final images. In order to compare the images, they were created with the maximum noise level of  
15 all three images (1.1 mJy/beam) and the enclosing beam for convolution, i.e. the restoring beam  
16 which encompasses the synthesis beam of all observations (here  $2.26 \times 0.79 \text{ mas}^2$  at a position an-  
17 gle of  $9.5^\circ$ ). In addition, Fig. 1 shows the individual images restored with the individual beams and  
18 contours starting at  $3\sigma_{\text{RMS}}$ . The brightness temperature  $T_{\text{B}}$  and size of the central emission region  
19 diameter  $d_{\text{core}}$  from each observation were determined by fitting a two-dimensional Gaussian func-

20 tion to the visibility data.  $T_B$  was then calculated following standard procedures<sup>48</sup>. By adopting  
21 a cosmology<sup>49</sup> with  $H_0 = 73 \text{ km s}^{-1} \text{ Mpc}^{-1}$ ,  $\Omega_\Lambda = 0.73$  and  $\Omega_m = 0.27$  for the conversion from  
22 angular to linear scales, 1 mas corresponds to about 8.3 pc at a redshift of  $z = 1.522$ .

23 ***Fermi/LAT*  $\gamma$ -ray data analysis:** For the analysis of *Fermi/LAT*  $\gamma$ -ray data, we used the *Fermi*  
24 Science Tools (v9r32p5) with the reprocessed Pass 7 data and the P7REP\_SOURCE\_V15 instru-  
25 ment response functions<sup>50</sup> and a region of interest (ROI) of  $10^\circ$  in the energy range 100 MeV  
26 to 300 GeV. We applied a zenith angle limit of  $100^\circ$  and the GTMKTIME cut `DATA_QUAL==1`  
27 `&& LAT_CONFIG==1` `&& ABS(ROCK_ANGLE)<52`. We used the Galactic diffuse emission  
28 model `GLL_IEM_V05.FITS` and the isotropic diffuse model `ISO_SOURCE_V05.TXT`. The model  
29 consists of all 2FGL sources inside the ROI. For the spectral and light-curve analysis, the normal-  
30 ization constants of all model sources were left free. The source spectral index was left free for the  
31 spectral analysis and frozen for the light curve analysis during the modeling of the spectrum. The  
32  $\gamma$ -ray light curve of PKS B1424–418 (Fig. 1) has been calculated using 14 day bins. We applied  
33 a Bayesian-blocks analysis<sup>51</sup> to the light curve. We used the change points to determine the onset  
34 of the outburst as 2012 Jul 16, which is marked as a shaded blue region in Fig. 1.

35 All  $\gamma$ -ray sources detected by *Fermi/LAT* with flares that exceeded the pre-defined threshold  
36 of  $10^{-6} \text{ cm}^{-2} \text{ s}^{-1}$  ( $> 100 \text{ MeV}$ ) are included in the *LAT Monitored Source List*. Light curves of  
37 all these sources are publicly available at [http://fermi.gsfc.nasa.gov/ssc/data/](http://fermi.gsfc.nasa.gov/ssc/data/access/lat/mssl_lc/)  
38 [access/lat/mssl\\_lc/](http://fermi.gsfc.nasa.gov/ssc/data/access/lat/mssl_lc/). By definition, this sample represents the brightest *Fermi/LAT* sources.  
39 We used this list to identify the sources with the highest  $\gamma$ -ray fluences and analyzed their  $\gamma$ -ray

40 light curves in detail as described above. We integrated the light curves from 2012 Jul 16 through  
41 2013 Apr 30 to determine the highest-fluence sources during the months around the arrival time  
42 of the 2 PeV IceCube neutrino event. PKS B1424–418 is the brightest blazar in terms of  $\gamma$ -ray  
43 fluence:  $F = (30.5 \pm 0.3) \text{ cm}^{-2}$  (the uncertainty is statistical only). The next on the list are the  
44 blazars CTA 102 ( $F = (20.0 \pm 0.3) \text{ cm}^{-2}$ ), and B2 1633+38 ( $F = (18.2 \pm 0.3) \text{ cm}^{-2}$ ), which are  
45 both in the northern hemisphere.

46 **Broadband SEDs:** The broadband SED of PKS B1424–418 (Fig. 2) has been built from quasi-  
47 simultaneous data from *Fermi*/LAT, *Swift*/XRT, *Swift*/UVOT, SMARTS<sup>52</sup>, ALMA, ATCA and the  
48 LBA. Archival data from 2MASS<sup>53</sup> and WISE<sup>54</sup> catalogs are included as well.

49 *Swift*/XRT data were reduced with standard methods, using the most recent software pack-  
50 ages (HEASOFT 6.15.1) and calibration databases. Spectra were grouped to a minimum signal-  
51 to-noise ratio of 5. Spectral fitting was performed with ISIS 1.6.2<sup>47</sup>. We fitted the (0.5–10) keV  
52 energy band with an absorbed power law model, which yielded good results for all time ranges.  
53 The source showed no evidence of intrinsic X-ray absorption in excess of the Galactic value<sup>55</sup>. X-  
54 ray data were deabsorbed using the best available abundances<sup>56</sup> and cross-sections<sup>57</sup>. *Swift*/UVOT  
55 data were extracted following standard methods. Optical, infrared, and ultraviolet data were dered-  
56 dened using the same absorbing columns<sup>58</sup>. At hard X-ray energies, we used data from *Swift*/BAT  
57 and *INTEGRAL*/IBIS from the energy bands of 20 keV to 100 keV and 20 keV to 200 keV, respec-  
58 tively. The source was not included in the BAT 70-month catalog<sup>59</sup>, which was based on data  
59 between 2004 Dec and 2010 Dec. Using stacked BAT maps from the 104-month survey (2004

60 Dec through 2013 Aug), we detected the source and calculated the flux by applying a power-law  
61 fit to the spectrum. For *INTEGRAL*/IBIS we extracted a  $2.5\sigma$  upper limit from the variance of a  
62 mosaic combining all pointed observations in the time-range from 2011 Dec to 2013 Sep, where  
63 the source was located within  $14^\circ$  from the pointing direction. The 104-month average BAT flux  
64 value is in good agreement with the 2LAC SED, whose data time period overlaps (2008 Aug to  
65 2010 Aug). It is marginally consistent with the SEDs corresponding to the post-2010 short-flare  
66 and the HESE period and well below the SED for the high-fluence period between 2012 Jul and  
67 2013 Apr.

68 The broadband spectra were fit with two logarithmic parabolas<sup>60</sup> in order to parametrize the  
69 characteristic double humped structure. The spectral-model fit is performed in detector space while  
70 the fluxes in each energy bin in the top panel of Fig. 2 are calculated independent of the model<sup>61</sup>,  
71 under the assumption of constant flux in each bin. Note that for curved spectra and wide energy  
72 bins, this approach can lead to an apparent positive offset of the flux data points over the best-fit  
73 model as seen, e.g., in the *Fermi*/LAT data points of the ‘short flare’ in Fig. 2. X-ray absorption and  
74 optical extinction were included in the broadband fit. Upper limits are generally not considered in  
75 the fits. The low  $\gamma$ -ray upper limits at the highest *Fermi*/LAT energies indicate a possible cutoff.  
76 Including the upper limits as values in the SED fits forces the parabolas to increase in peak flux in  
77 order to still describe the slope at X-ray energies and reduces the overall fit quality. The increase  
78 in peak flux leads to a higher neutrino estimate (see below), with the exception of the short flare,  
79 so that our method to ignore upper limits in the fit is more conservative for our purposes.

80 **Neutrino event prediction:** To derive maximal-possible neutrino events to be expected from the  
81 measured SEDs, we applied the techniques that we developed for our analysis of the AGN in the  
82 fields of the first two PeV IceCube events. We used 2FGL GeV  $\gamma$ -ray spectra<sup>62</sup> and extracted  
83 X-ray spectra taken during the same period by *Swift*<sup>63</sup>. For sources not observed by *Swift*, we  
84 used X-ray data or upper limits from the *ROSAT* all-sky survey<sup>64</sup>. For PKS B1424–418, we  
85 calculated in addition *Fermi*/LAT  $\gamma$ -ray spectra for the various time ranges of interest according  
86 to the methods described above. We fit a single log parabola to the high-energy data. From the  
87 integrated 1 keV to 5 GeV SEDs, we can derive a maximal-possible neutrino flux for each source.  
88 From this, we can derive the maximal-possible number of counts for the 988-day HESE livetime  
89 (see Table 3) using the appropriate effective area for northern/southern-hemisphere sources for  
90 electron neutrinos of the IceCube 988-day analysis<sup>3</sup>. Formal uncertainties of neutrino counts can  
91 be derived from the uncertainties of the integrated electromagnetic energy output but are dominated  
92 by systematic uncertainties due to the simplicity of the model and method. In particular, note that  
93 the overall shape of the SED changes during long time ranges and that the derived neutrino output  
94 from averaged LAT spectra is not necessarily equal to the sum of neutrino output from additive  
95 segments of the full time range, especially in the presence of a strong outburst. This systematic  
96 uncertainty is illustrated by the fact the maximal-possible neutrino output derived from the average  
97 3-yr SED of PKS B1424–418 is 4.5 events, while the separate consideration of the 27-month pre-  
98 outburst phase and the 9-month outburst yields a higher maximal-possible neutrino output of 5.7  
99 events (cf. Table 2).

100 **Calculation of chance coincidences of major blazar outbursts and PeV neutrino events:** We  
101 estimate the number of chance coincidences of major blazar outbursts and PeV neutrino events  
102 *a posteriori* using  $N_{\text{coinc}} = \dot{N}_{\nu} \times \dot{N}_{\text{outburst}} \times \tau \times T$ . Here,  $\dot{N}_{\nu}$  is the number of PeV-neutrinos  
103 per year in the southern sky,  $\Omega_{\text{south}} = 2\pi$ ,  $\dot{N}_{\text{outburst}}$  is the number of major outbursts per year  
104 within a cone covered by the typical median uncertainty of shower-like PeV neutrino events of  
105  $\Omega_{\nu} \sim 0.16$  sr (averaged median angular uncertainty of the three observed PeV events  $\sim 13^\circ$ ),  $\tau$   
106 is the coincidence time window defined by the outburst time, and  $T$  is the IceCube observation  
107 time of  $T = 3$  yr, during which IceCube has observed three PeV events. We therefore obtain  
108 for  $\dot{N}_{\nu} \sim 1 \text{ yr}^{-1}$ . High-fluence blazar outbursts occur on time scales of months so we choose,  
109 conservatively,  $\tau = 1$  yr for the time window for chance coincidences. In the *Fermi*/LAT blazar  
110 light curves for six years and for the whole sky, we find a total of eight 1 yr periods, during which  
111 an individual blazar reached or exceeded the GeV fluence of the PKS B1424–418 outburst. These  
112 high-fluence phases occurred in the sources PKS B1424–418, PKS B 1222+216, 3C 273, 3C 454.3  
113 (2 periods), and PKS B1510–089 (3 periods). Hence, the rate of such outbursts occurring within  
114  $\Omega_{\nu}$  is expected to be  $\dot{N}_{\text{outburst}} \sim 0.1 \text{ yr}^{-1} \text{ sr}^{-1} \times \Omega_{\nu} = 0.016 \text{ yr}^{-1}$ . Using these numbers, we end  
115 up with  $N_{\text{coinc}} \sim 0.05$  as the mean number of chance coincidences. The Poisson probability to  
116 observe one or more such coincidences is about 5%. However, because of the lack of a pre-defined  
117 statistical test, we cannot formally use this value to test the hypothesis of a chance coincidence.

## 118 2 Supplementary discussions

119 **Coincidences of sub-PeV neutrino events with high-fluence blazars:** All high-fluence blazars  
120 from Table 3 are in coincidence with more than one IceCube sub-PeV neutrino event, which is  
121 not surprising given the high probability of chance coincidences within the large  $R_{50}$  radii of  
122 the cascade-like IceCube events. In addition, up to 25 of the 34 sub-PeV events are of possible  
123 atmospheric-neutrino or cosmic-ray muon origin<sup>3</sup>. The two most interesting cases involve HESE-  
124 7, which coincides with long-lasting outbursts of the two high-fluence blazars PKS B0537–441  
125 and PKS B2326–502. Both sources are located within the  $\Omega_{\text{HESE-7}}^{R_{50}}$  field. The integrated fluence of  
126 PKS B0537–441 and PKS B2326–502 predicts a maximum of  $\sim 5$  neutrino events, i.e., applying  
127 the scaling factor of 0.025, the Poisson probability for the detection of more than zero PeV events  
128 is about 12%. The IceCube events HESE-16 and HESE-25 are close ( $1.9R_{50}$  and  $1.2R_{50}$ , respec-  
129 tively) to the position of the top-ranked high-fluence source, PKS B1510–089. It is remarkable  
130 that the two observed events occurred close to the highest peaks of the GeV light curve but we em-  
131 phasize again that a considerable fraction of IceCube events at these comparatively low energies  
132 can be of atmospheric origin and that the field of HESE-25 is particularly large.

133 **Coincidences of track neutrino events with  $\gamma$ -ray blazars:** We note that the IceCube track  
134 events, which have a much better angular resolution than the shower events, are all of rather low  
135 energy and/or are located at far northern declinations. Most of them are thus likely of atmospheric  
136 origin. In fact, event HESE-28 showed associated hits in the IceTop surface air shower array<sup>2,3</sup>.  
137 The IceCube team considers three additional track events particularly ambiguous, which started

138 near the detector boundary and were down going (HESE-3, HESE-8, HESE-18, HESE-23). The  
139 positions of the remaining three (HESE-5, HESE-13, and HESE-37) which can be considered  
140 the most-likely track events to be of true extraterrestrial origin, have been compared with posi-  
141 tions of known  $\gamma$ -ray AGN. One of these events (HESE-5) is spatially coincident with the blazar  
142 PKS 0723–008, but as pointed out previously<sup>65</sup>, the chance probability for a *Fermi*/LAT detected  
143 AGN to coincide with any given field of about  $1 \text{ deg}^2$  is high. We estimate the Poisson probability  
144 for HESE-5 from the known 1017  $\gamma$ -ray AGN in the 2LAC catalog<sup>24</sup> distributed over the full sky  
145 (excluding the  $|b| < 10^\circ$  region around the Galactic plane) to be 41 %. However, as mentioned  
146 already above, only half of all IceCube events are expected to hold the coordinates of their true  
147 astrophysical sources inside their  $R_{50}$  regions, while the other half are located at larger positional  
148 offsets. In this small sample of three track events, we thus expect only one or two coincidences  
149 within  $R_{50}$ . If we consider radii twice as large, we find one 2LAC blazar for HESE-13 (4C +41.11  
150 at  $1.8R_{50}$ ) and one for HESE-37 (RBS 0958 at  $1.6R_{50}$ ). The *a posteriori* chance probability for  
151 all three track events to coincide within  $2R_{50}$  with a 2LAC blazar is 22 %. None of the candi-  
152 date blazars is of particularly high fluence. It is likely that they represent the large population of  
153 ‘typical’ 2LAC  $\gamma$ -ray blazars, which have low individual detection probabilities. However, their  
154 integrated emission plus the contribution of fainter unresolved blazars to the EGB is high and may  
155 lead to the detected events.



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201 **4 Supplementary Tables**

**Supplementary Table 1: Details of TANAMI VLBI observations of PKS B1424–418 and image parameters**

Date	$S_{\text{total}}$	$S_{\text{peak}}$	$\sigma_{\text{RMS}}$	$b_{\text{maj}}$	$b_{\text{min}}$	P.A.	$T_{\text{B,core}}$	$d_{\text{core}}$
[yyyy-mm-dd]	[Jy]	[Jy/beam]	[mJy/beam]	[mas]	[mas]	[°]	[ $10^{11}$ K]	[pc]
2011 Nov 13	$2.33 \pm 0.23$	$1.80 \pm 0.18$	0.2	1.95	0.50	3.8	4.2	2.5
2012 Sep 16	$3.17 \pm 0.32$	$2.89 \pm 0.29$	0.4	2.20	0.71	4.0	7.4	2.2
2013 Mar 14	$6.23 \pm 0.62$	$5.61 \pm 0.56$	1.1	2.26	0.79	9.5	13	2.3

Note: Tabulated image parameters refer to the formal values after hybrid imaging using natural weighting<sup>17</sup>.

The three images shown in Fig. 1 are all convolved with the same restoring beam and use uniform contour lines and color scaling. See Fig. 1 for the individual images restored with the individual beams and contours starting at  $3\sigma_{\text{RMS}}$ . Core FWHM diameters  $d_{\text{core}}$  and  $T_{\text{B,core}}$  have been derived from fitting an elliptical Gaussian component to the visibility data.

Array configurations: PK-AT-HO-CD-TC-TD-HH-MP-WW (2011 Nov 13); PK-AT-HO-CD-TC-HH-AK-KE (2012 Sep 16); PK-AT-HO-CD-TC-TI-HH-WW-AK-KE (2013 Mar 14); PK: Parkes (64 m), AT: ATCA (5x22 m), HO: Hobart (26 m), CD: Ceduna (30 m), TC: TIGO (6 m), TI: Tidbinbilla (70 m), TD: Tidbinbilla (34 m), MP: Mopra (22 m), HH: Hartebeesthoek (26 m), WW: Warkworth (12 m), AK: ASKAP (12 m), KE: Katherine (12 m).

202 **5 Supplementary Figures**

203 **Figure 1 Supplementary Figure 1** – Individual TANAMI VLBI images of the core region  
204 of PKS B1424–418. Each image corresponds to the properties in Supplementary Table  
205 1 for the individual observations. All contours start at  $3\sigma_{\text{RMS}}$  and increase logarithmically  
206 by factors of 2 .

