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Tidally disrupted dusty clumps as the origin of broad emission lines in active galactic nuclei

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The broad-line region of active galactic nuclei: tidally disrupted clumps from the torus

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Supplementary Table 1

A list of parameters involved in the present model

Parameter	Physical meanings	Valid ranges	Uncertainties ($\langle\langle\Delta X/X\rangle\rangle$)
ζ_0	$R_T^0 \omega_A^0 / v_A^0$: determines trajectories of clumps	$0 \sim 50$	1.6
γ_A	cloud A distribution index: $N_A = N_A^0 v_A^{\gamma_A}$	$0 - 4$	0.5
ξ_A	$v_A^0 = \xi_A (\zeta_0^2 + 1)^{-1/2} v_K^0$, v_K^0 is the Keplerian velocity at R_T^0 .	< 1	0.2
R_0/R_g	the circulaized radius of type A clouds	$\sim 6.2 \times 10^3$	0.3
R_{in}	radius of type B clouds merging with accretion disks	$< R_{BLR}$	0.6
γ_B	cloud B distribution index: $N_B = N_B^0 v_B^{\gamma_B}$	$0 - 4$	0.5
f_B	fraction of type B clouds	$0 \sim 1$	0.3
ξ_C	$v_C^0 = \xi_C v_K^0$, v_C^0 is the ejection velocity of type C clouds at R_T^0 .	> 1	0.3
f_C	fraction of type C clouds, properties shown by Equation (7)	$0 \sim 1$	0.6
i	inclination angle of observers, $i = 0^\circ$ (edge-on) and $i = 90^\circ$ (face-on)	$\gtrsim 40^\circ$	0.2
Auxiliary Parameters (the fixed values are listed by the numbers in bold fonts)			
Θ_{torus}	half angle of torus given by the $\Theta_{torus} - L_{[OIII]}$ relation	Ref. ¹	
α_A	cloud A velocity index: $v_A = v_A^0 (R/R_T^0)^{-\alpha_A}$	0.45 ; ≤ 0.5	
α_C	cloud C velocity index: $v_C = v_C^0 (R/R_T^0)^{-\alpha_C}$	0.30 ; ≤ 0.5	
β_A	cloud A angular velocity index: $\omega_A = \omega_A^0 (R/R_T^0)^{-\beta_A}$	1.8 ; ≥ 1.5	
γ_C	cloud C distribution index: $N_C = N_C^0 v_C^{\gamma_C}$	0 ; ~ 0.5	
γ	vertical distribution of clumps in torus: $\propto (\cos \theta_C)^\gamma$	1 ; ~ 1	
Γ	$1 + \alpha_A - \beta_A$		
R_T^0/R_g	tidal disruption radius of clumps.	10^4 ; (R_0, R_{torus})	
R_{BLR}	the emissivity-averaged radius of the BLR determined by RM.		
φ_0	position angle of tidal event.	$\varphi_0 = \pi/2$	
ω_A^0	angular velocity of type A clouds at R_T^0 .	absorbed by ζ_0	

Note: the averaged $\langle\langle\Delta X/X\rangle\rangle$ is obtained by the statistic of individual $\Delta X/X$ of the sample. Here X is any one of the fitting parameters in the model.

Supplementary Table 2

Resultant parameters of best-fittings of four quasars

Parameter	PG1354+213	PG2251+113	PG1351+640	PG1700+518
ζ_0	6.1_{-6}^{+10}	$0.71_{-0.5}^{+1.5}$	$0.02_{-0.02}^{+5}$	38_{-31}^{+13}
ξ_A	$0.36_{-0.1}^{+0.3}$	$0.59_{-0.1}^{+0.2}$	$0.42_{-0.1}^{+0.1}$	$0.63_{-0.1}^{+0.1}$
γ_A	$0.79_{-0.79}^{+1.4}$	$3.6_{-1.8}^{+0.9}$	$0.51_{-0.2}^{+0.9}$	$1.5_{-0.5}^{+0.9}$
$R_0(10^3 R_g)$	$3.2_{-0.9}^{+1.2}$	$2.6_{-0.7}^{+1.0}$	$1.4_{-0.8}^{+1.5}$	$7.5_{-2.5}^{+1.5}$
$R_{in}(10^3 R_g)$	$0.31_{-0.1}^{+0.2}$	$0.38_{-0.1}^{+0.1}$	$0.39_{-0.1}^{+0.1}$	$3.6_{-0.9}^{+1.3}$
γ_B	$1.2_{-0.1}^{+0.1}$	$2.1_{-1.0}^{+0.9}$	$2.8_{-1.8}^{+0.9}$	$3.0_{-1.3}^{+0.9}$
f_B	$0.95_{-0.6}^{+0.1}$	$0.78_{-0.2}^{+0.1}$	$0.73_{-0.1}^{+0.1}$	$0.24_{-0.1}^{+0.0}$
ξ_C	$2.7_{-1.1}^{+0.5}$	$1.1_{-0.5}^{+0.2}$	$2.2_{-0.4}^{+0.1}$	$2.9_{-0.1}^{+0.1}$
f_C	$0.02_{-0.02}^{+0.1}$	$0.18_{-0.1}^{+0.1}$	$0.04_{-0.01}^{+0.01}$	$0.44_{-0.1}^{+0.1}$
$i(^{\circ})$	75_{-13}^{+8}	66_{-9}^{+9}	74_{-9}^{+5}	61_{-1}^{+2}
Line profiles				
$\mathcal{A}_{H\beta}$	4.85	0.15	12.8	2.90
$\mathcal{S}_{H\beta}$	0.30	0.22	1.41	1.07
$\mathcal{Z}_{H\beta}$	1.34	1.12	1.77	2.15

Parameters of all the PG quasars will be provided on request.

Supplementary Table 3

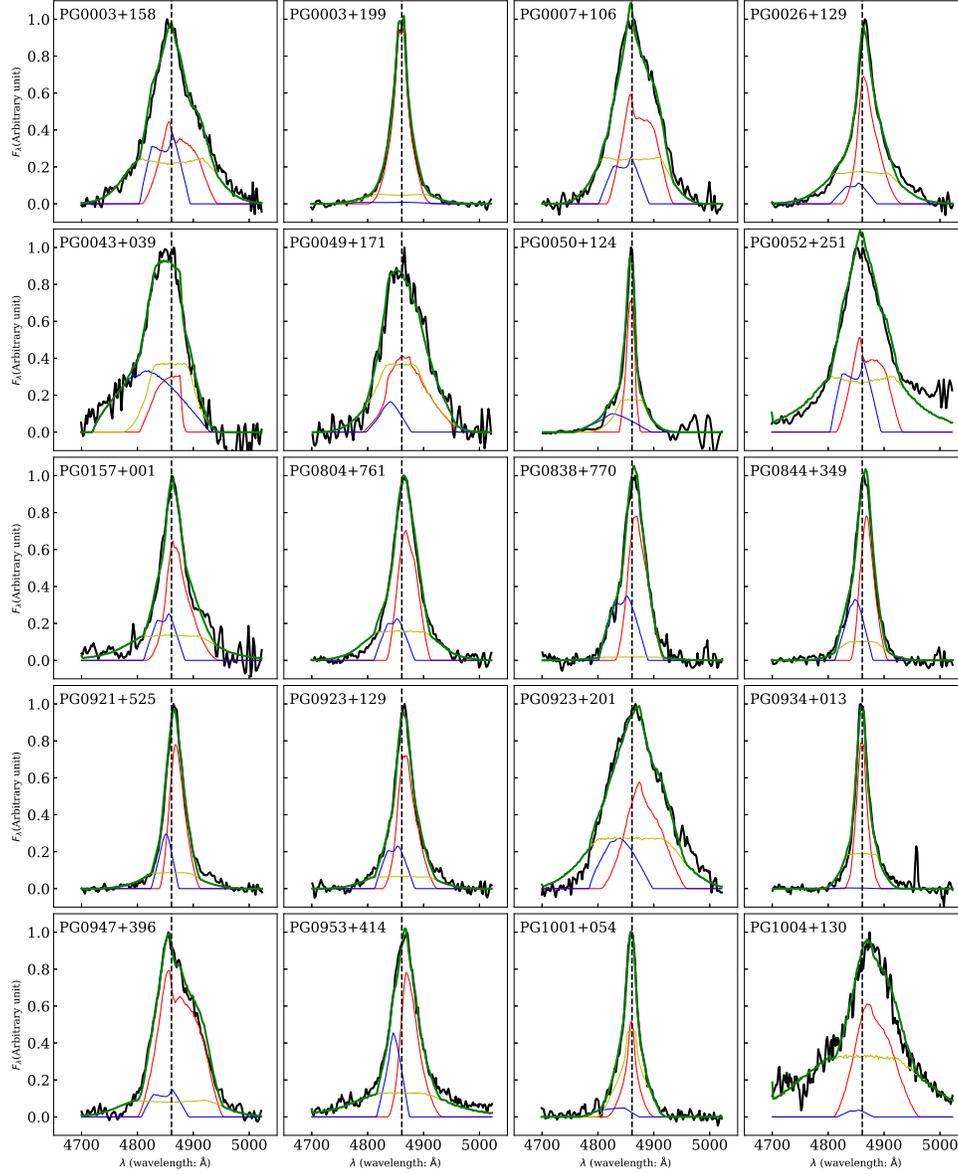
Classifications of resultant fittings

Model	Number	Objects (total of 87 PG quasars)
B	1	1354+213
B+C	1	2251+113
A+B	12	0003+199, 0934+013, 1004+130, 1012+008, 1103-006, 1211+143, 1351+236 1351+640, 1512+370, 1534+580, 1617+175, 2233+134
A+B+C	73	0003+158, 0007+106, 0026+129, 0043+039, 0049+171, 0050+124, 0052+251 0157+001, 0804+761, 0838+770, 0844+349, 0921+525, 0923+129, 0923+201 0947+396, 0953+414, 1001+054, 1011-040, 1022+519, 1048-090, 1048+342 1049-006, 1100+772, 1114+445, 1115+407, 1116+215, 1119+120, 1121+422 1126-041, 1149-110, 1151+117, 1202+281, 1216+069, 1226+023, 1229+204 1244+026, 1259+593, 1302-102, 1307+085, 1309+355, 1310-108, 1322+659 1341+258, 1352+183, 1402+261, 1404+226, 1411+442, 1415+451, 1416-129 1425+267, 1426+015, 1427+480, 1435-067, 1440+356, 1444+407, 1448+273 1501+106, 1519+226, 1535+547, 1543+489, 1545+210, 1552+085, 1612+261 1613+658, 1626+554, 1700+518, 1704+608, 2112+059, 2130+099, 2209+184 2214+139, 2304+042, 2308+098

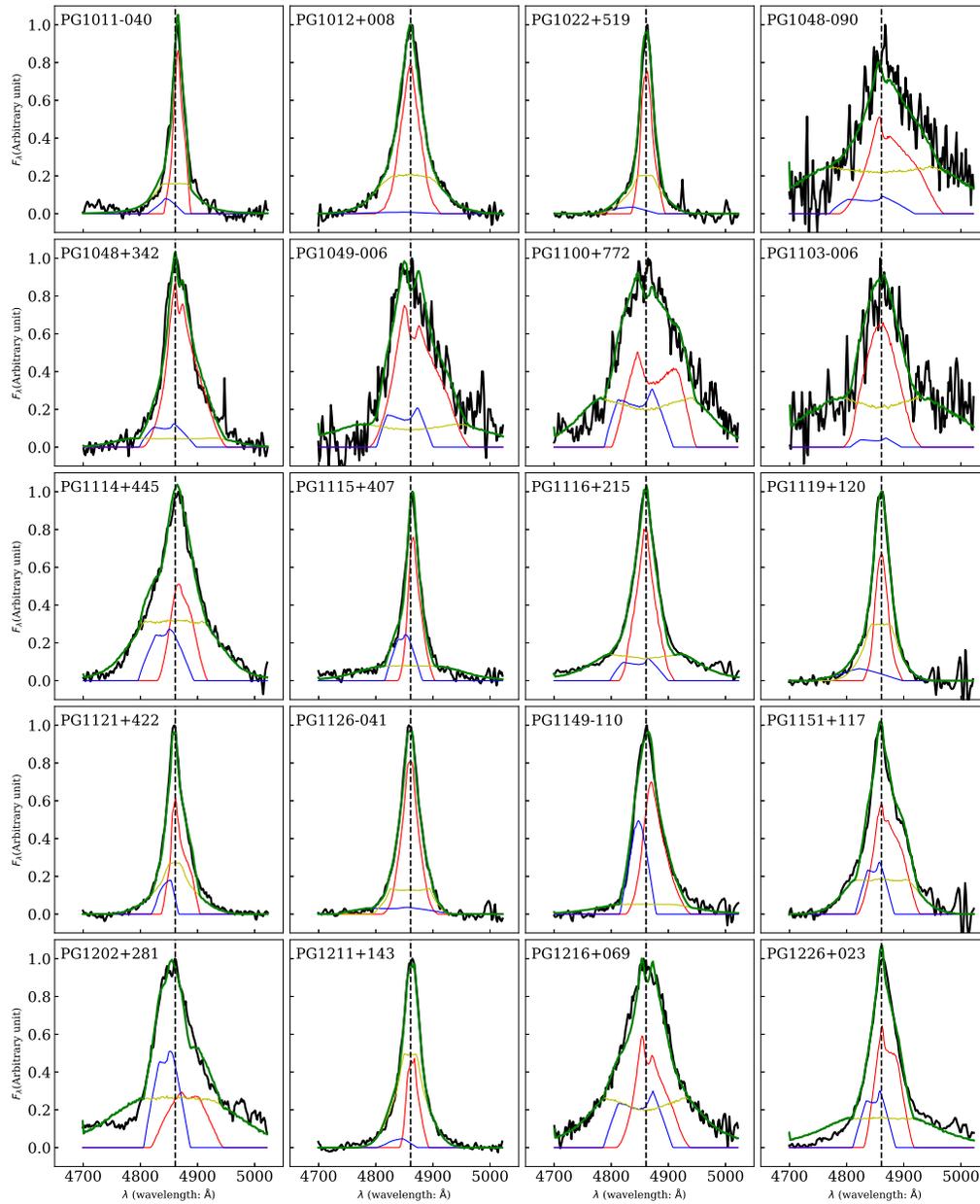
B: $f_{A,C} \leq 5\%$; **B+C:** $f_A \leq 5\%$; **A+B:** $f_C \leq 5\%$; **A+B+C:** $f_{A,B,C} > 5\%$.

If it is less than 5%, we consider that component to be negligible. There are only two profiles with B and (B+C) components, only 12 objects with (A+B), while the majority of 73 objects employ (A+B+C).

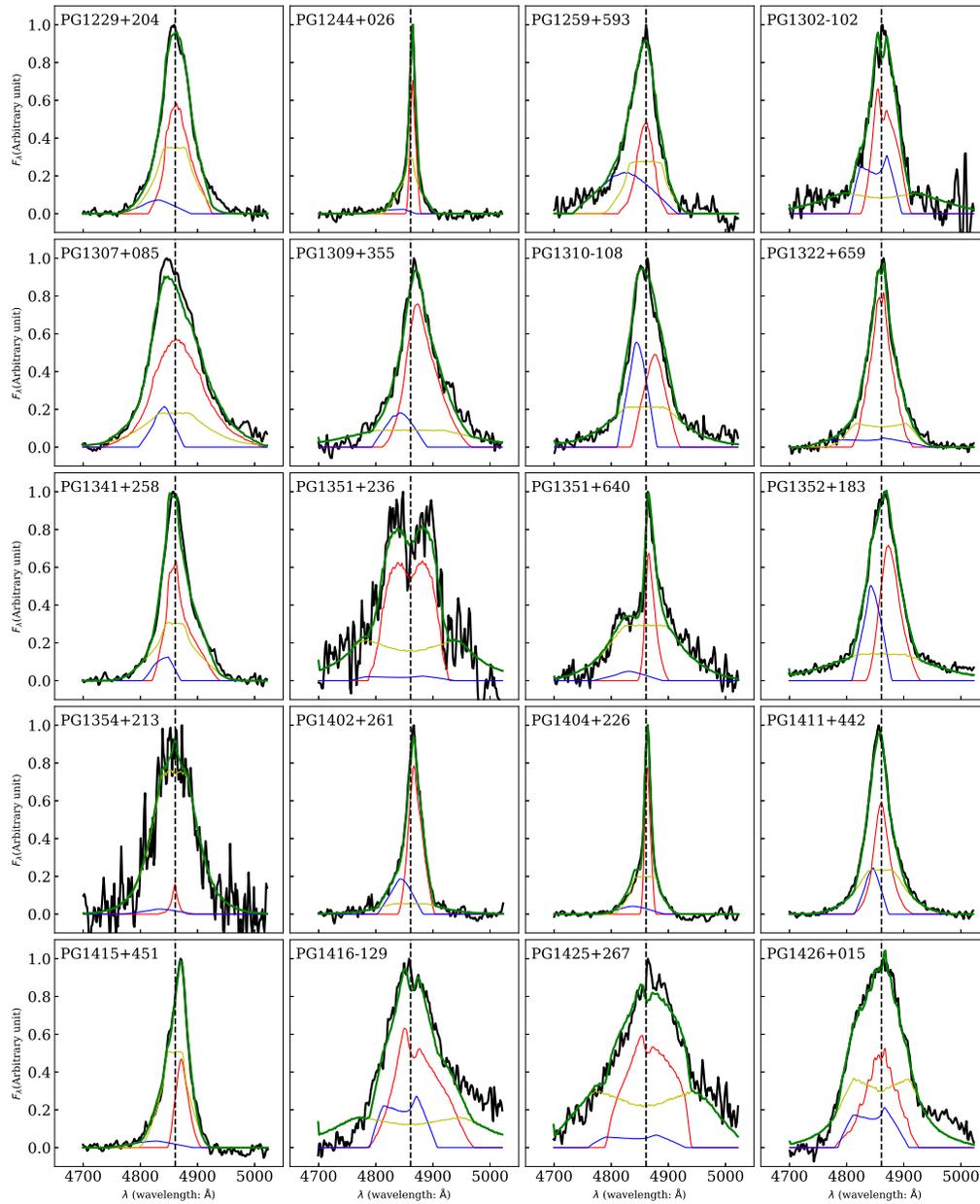
Generally, type B clouds are necessary in all objects, type A appears in 85/87 objects, and the type C in 74/87.



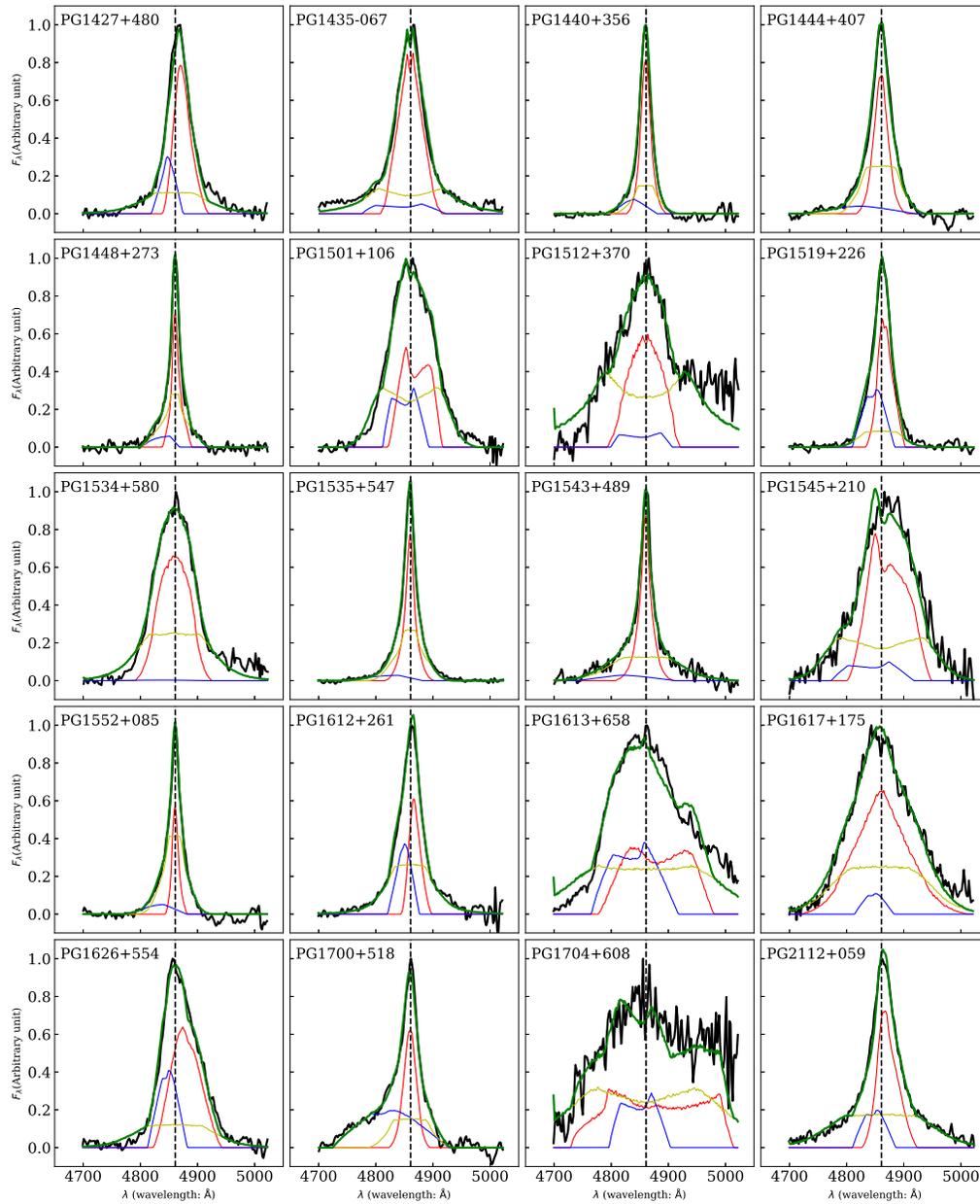
Supplementary Figure 1: Comparisons of our model fits with observed $H\beta$ profiles of PG quasars. Solid black lines are the observed profiles, red lines are from type A clouds, the yellow from type B and the blue from type C and the green is the total of the three cloud types. The dashed lines are $\lambda_0 = 4861 \text{ \AA}$.



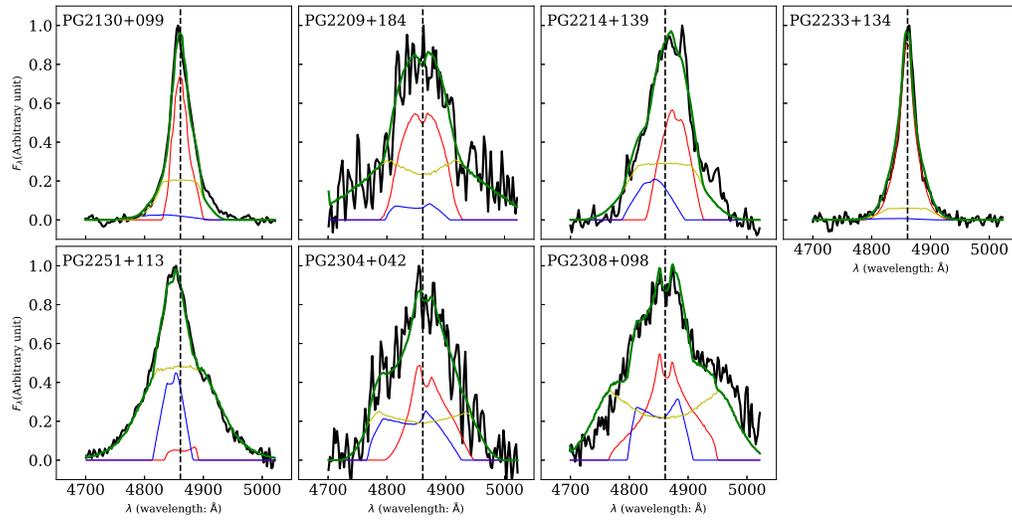
Supplementary Figure 1 *Continued.*



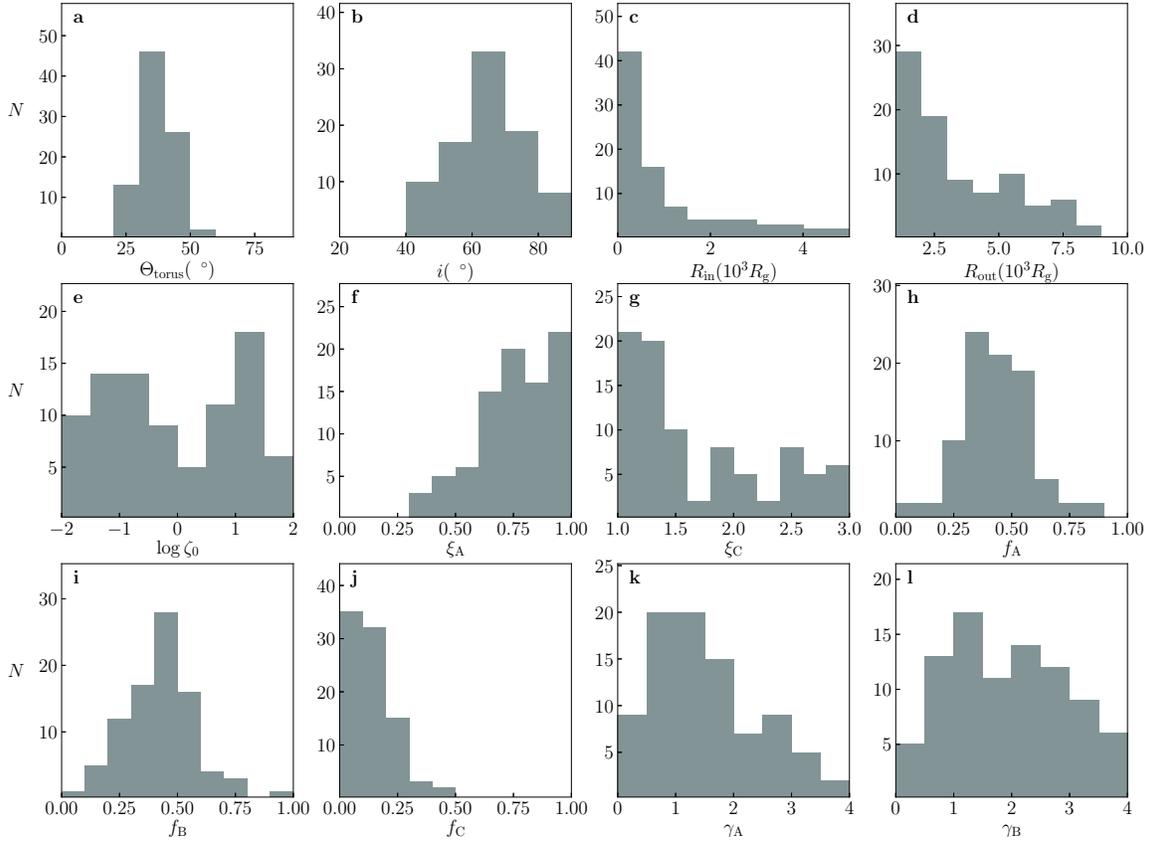
Supplementary Figure 1 *Continued.*



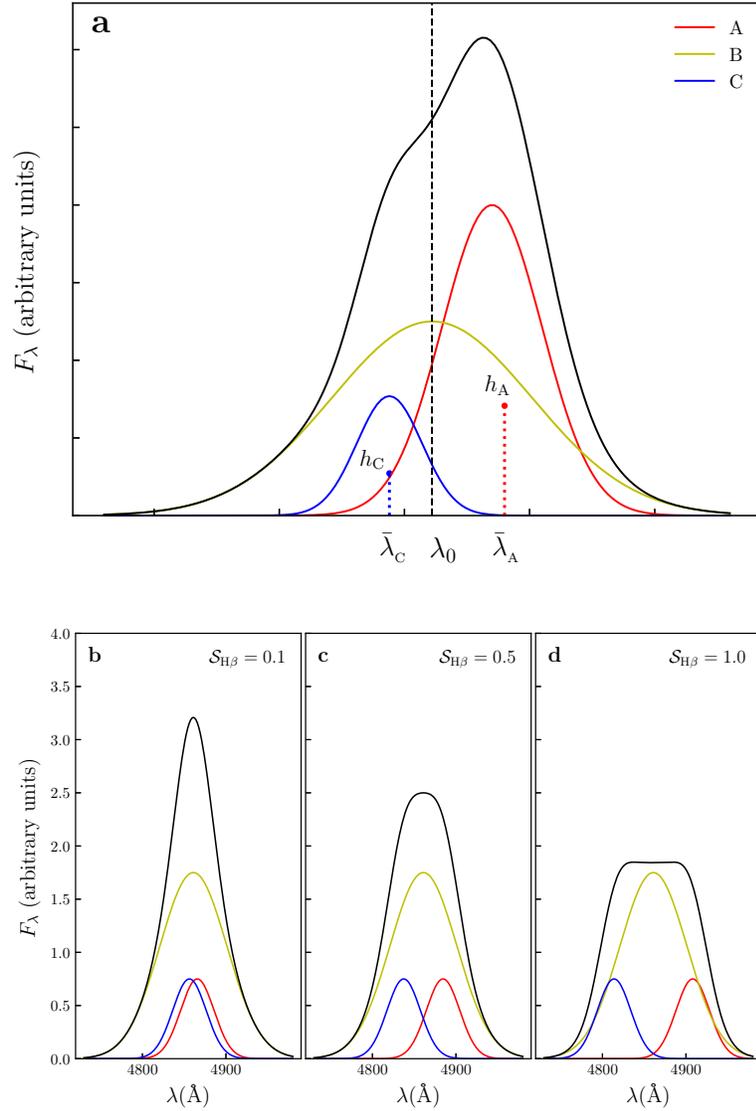
Supplementary Figure 1 *Continued*.



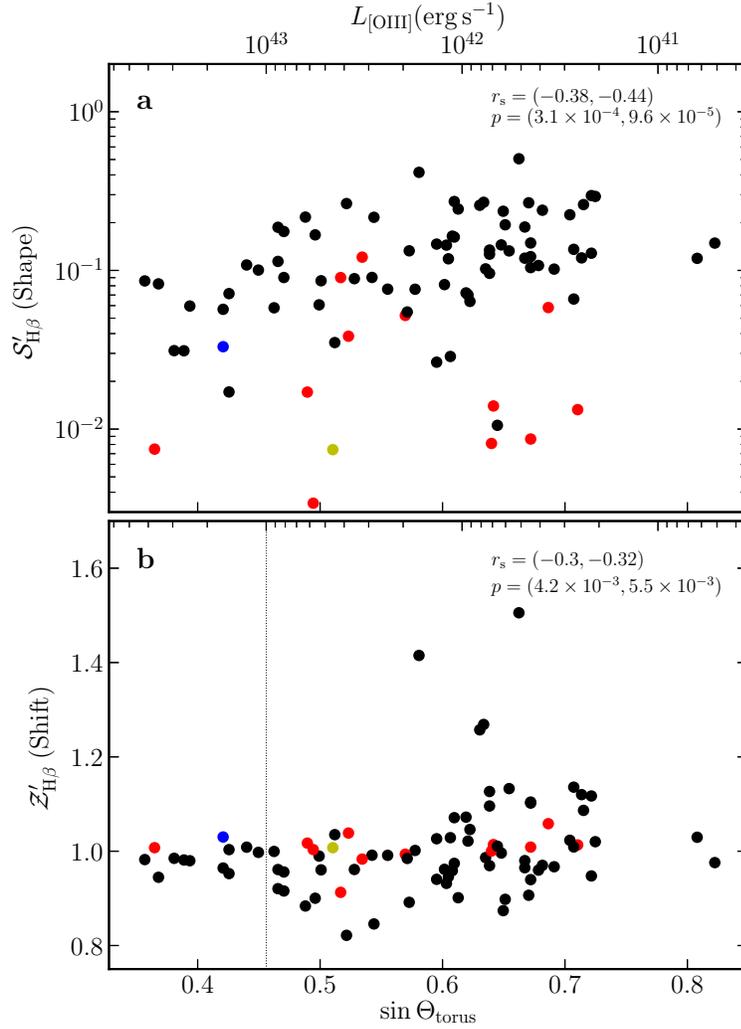
Supplementary Figure 1 *Continued.*



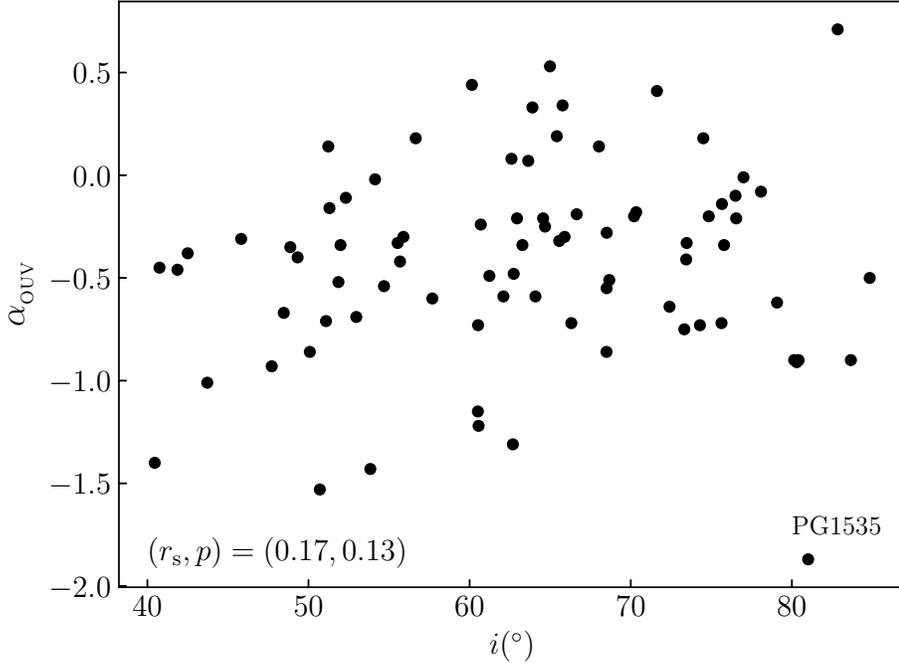
Supplementary Figure 2: **Distributions of the parameters obtained by fitting $H\beta$ profiles.** Note $f_C = 1 - (f_A + f_B)$, the average fraction of type C clouds is $\langle f_C \rangle \approx 0.1$ and $f_C \ll (f_B, f_A)$ holds for most objects of the PG sample. There are a couple of objects with $\xi_C \gtrsim 2$, potentially implying acceleration of type C clouds driven by radiation pressure, but this possibility needs to be explored by numerical simulations.



Supplementary Figure 3: **Illustration of physical meanings of parameters describing profiles.** Panel *a* shows a cartoon of decomposed profiles for physical meanings of the three parameters defined by Equation (16). The barycentre wavelength and fluxes are obtained by including flux-weight. As the first order approximation, the two parameters of $(h, \bar{\lambda})$ can completely describe individual profiles of the decomposed components and hence provide reasonable asymmetries, shapes and shifts of the total spectra. Panels *b-d* illustrate the relationship between $S_{H\beta}$ and real shapes of profiles (from a triangular to boxy).



Supplementary Figure 4: **Correlations between $S'_{H\beta}$ and $Z'_{H\beta}$ versus $L_{[OIII]}$.** Panel a shows a consistent correlation with panel b in Figure 4. There are four outliers with $Z'_{H\beta} \gtrsim 1.2$ deviating from the correlation in panel b. For objects with $L_{[OIII]} \gtrsim 10^{43} \text{ erg s}^{-1}$, $Z'_{H\beta} \approx 1$ remains, which have smaller torus angles ($\Theta_{\text{torus}} \lesssim 30^\circ$).



Supplementary Figure 5: **A comparison of PG quasar inclinations from the present model with α_{OUV} as an orientation².** This shows that they are consistent with each others. The Spearman coefficient and null-probability are indicated in the plot except for PG 1535. Error bars of α_{OUV} are not given by Ref.², and i is given in Supplementary Table 1.

1. Reyes, R. *et al.* Space density of optically selected type 2 quasars. *Astron. J.* **136**, 2373-2390 (2008).
2. Baskin, A. & Laor, A. What controls the C IV line profile in active galactic nuclei? *Mon. Not. R. Astron. Soc.* **356**, 1029–1044 (2005).