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Intrinsic spin Hall torque in a moiré Chern magnet

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SUPPLEMENTARY MATERIALS



Fig. S1. NanoSQUID imaging at $\nu = -1$ on monolayer MoTe₂/monolayer WSe₂ device. (A) Optical image of a monolayer MoTe₂/monolayer WSe₂ device. (B) NanoSQUID imaging revealed qualitatively similar magnetism and anomalous Hall resistances in this device. This device was rather disordered and did not show quantization of anomalous Hall resistance.



Fig. S2. Phase diagram from transport data (a) Two terminal resistance of device (measured between contacts *abcd* and *j*) plotted as a function of electron density *n* and displacement field *D*, where $D = \frac{V_{TG} - V_{BG}}{d_t + d_b}$, $n = \epsilon \epsilon_0 (V_{TG}/d_t + V_{BG}/d_b) + n_0$. Here $d_t = 2.7$ nm is the thickness of the top hBN layer, $d_b = 12.1$ nm is the thickness of the bottom hBN layer, $\epsilon \approx 3$ is the relative dielectric constant of hBN, and n_0 is the offset charge carrier density, 5.9×10^{12} cm⁻². At low displacement fields $\nu = -1$ and $\nu = -2$ both host topologically trivial insulating states, with the $\nu = -1$ insulator presumably interaction-driven. In the electron-doping regime large contact resistances hamper transport measurements. (b) ΔR_{xy} measurement in the region inside the green dotted line in **a**. A finite ΔR_{xy} appears near $\nu = -1$ in a narrow range of *D*. (c) ΔR_{xy} measurement in the region inside the green dotted line in **b**. Precise quantization of ΔR_{xy} obtains over a range of displacement fields at $\nu = -1$. (d) Symmetrized R_{xx} measurement in the region inside the dotted line in **b**. (d) Linecut of ΔR_{xy} along the dotted line in **c** illustrating the appearance of a QAH effect.



Fig. S3. AC bottom gate conversion to magnetization. (a) AC bottom gate magnetometry produces $\delta B_z/\delta V_{BG}(x, y, V_{BG})$. (b) We integrate $\delta B_z/\delta V_{BG}(x, y, V_{BG})$ with respect to V_{BG} to obtain $B_z(x, y, V_{BG})$ (c) For each value of V_{BG} we invert B(x, y) to produce $m_z(x, y, V_{BG})$, which is presented as a video of the out-of-plane magnetization as V_{BG} enters and then leaves the QAH regime in the supplementary data.



Fig. S4. Repeated cooldowns (a) Magnetic structure measured at 28 mT, $V_{TG} = -3.734$ V, $V_{BG} = 7.818$ V at a height of 125 nm with nanoSQUID diameter $\phi = 159$ nm. (b) Magnetic structure measured at 36 mT, $V_{TG} = -3.756$ V, $V_{BG} = 7.943$ V during a subsequent cooldown at a height of 175 nm with nanoSQUID diameter $\phi = 113$ nm.



Fig. S5. Spin diffusion length (a) Spin Hall effect near $\nu = -1$ (b) Fit of the left edge. Green regions are masked from the fit, blue is fit to an exponential decay function, $B(x) = A \cdot e^{-\frac{x}{\lambda}} - c$. Approximate spin diffusion length λ is 700 nm. (c) Fit of the right edge, approximate spin diffusion length λ is 400 nm.



Fig. S6. Additional transport properties of the magnetic Chern insulator (a) Magnetic hysteresis loop in QAH regime. At $\pm 500 \text{ mT}$ quantization reaches $1.000 \pm 0.006 h/e^2$ and $-0.996 \pm 0.005 h/e^2$. (b) Close to $B = 0 \text{ R}_{xy} \approx 0.9 h/e^2$. Coercive fields are less than 1 mT at the measurement temperature of 1.6K. (c) Dependence of the degeneracy of the Chern band on B reveals the Chern number of the ground state at finite field, which is -1 in this system. (d) Linecut shows R_{xx} reaches $-93 \Omega \pm 115 \Omega$.



Fig. S7. Current-switching phase diagram (a) Bottom gate modulation magnetometry phase diagram taken at the point at which measurements for Fig. 4g were performed. The green dotted outline appears in that figure. $\delta V_{BG} = 35 \text{ mV}$ was used. (b) Current-induced magnetic domain switching signal δB_I as a function of top and bottom gate voltages with $I_{SD} = 130 \text{ nA}$, $\delta I_{SD} = 30 \text{ nA}$. (c) The same measurement with $I_{SD} = 290 \text{ nA}$, $\delta I_{SD} = 30 \text{ nA}$. Magnetic switching appears over broad regions of the magnetic phase diagram, although currents required to effect domain switching vary with gate voltages.