

Dictionary of American Family Names by Elsdon C. Smith and the Graphical Law

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Abstract

We study Dictionary of American Family Names by Elsdon C. Smith. We draw the natural logarithm of the number of entries, normalised, starting with a letter vs the natural logarithm of the rank of the letter, normalised. We conclude that the Dictionary can be characterised by $BP(4, \beta H = 0.01)$ i.e. a magnetisation curve for the Bethe-Peierls approximation of the Ising model with four nearest neighbours in the presence of external magnetic field, H , with $\beta H = 0.01$. β is $\frac{1}{k_B T}$ where, T is temperature and k_B is the tiny Boltzmann constant.

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I. INTRODUCTION

We go over in this paper to the interesting dictionary, Dictionary of American Family Names by Elsdon C. Smith, [1]. We study magnetic field pattern behind the entries of this dictionary, [1], in this article. We count the entries of this dictionary, one by one, to probe for the magnetic field pattern. We have started considering magnetic field pattern in [2], in the languages we converse with. We have studied there, a set of natural languages, [2] and have found existence of a magnetisation curve under each language. We have termed this phenomenon as the Graphical Law. Then, we moved on to investigate, [3], into dictionaries of five disciplines of knowledge and found the existence of a curve of magnetisation under each discipline. This was followed by finding of the graphical law in references from [4] to [77].

We describe how the graphical law is hidden within Dictionary of American Family Names by Elsdon C. Smith, [1]. in this article. The planning of the paper is as follows. We give an introduction to the standard curves of magnetisation of Ising model in the section II. In the section III, we describe analysis of the entries of Dictionary of American Family Names by Elsdon C. Smith [1]. The section IV is Acknowledgment. The last section is Bibliography.

II. MAGNETISATION

A. Bragg-Williams approximation

Let us consider a coin. Let us toss it many times. Probability of getting head or, tale is half i.e. we will get head and tale equal number of times. If we attach value one to head, minus one to tale, the average value we obtain, after many tossing is zero. Instead let us consider a one-sided loaded coin, say on the head side. The probability of getting head is more than one half, getting tale is less than one-half. Average value, in this case, after many tossing we obtain is non-zero, the precise number depends on the loading. The loaded coin is like ferromagnet, the unloaded coin is like para magnet, at zero external magnetic field. Average value we obtain is like magnetisation, loading is like coupling among the spins of the ferromagnetic units. Outcome of single coin toss is random, but average value we get after long sequence of tossing is fixed. This is long-range order. But if we take a small sequence of tossing, say, three consecutive tossing, the average value we obtain is not fixed,

can be anything. There is no short-range order.

Let us consider a row of spins, one can imagine them as spears which can be vertically up or, down. Assume there is a long-range order with probability to get a spin up is two third. That would mean when we consider a long sequence of spins, two third of those are with spin up. Moreover, assign with each up spin a value one and a down spin a value minus one. Then total spin we obtain is one third. This value is referred to as the value of long-range order parameter. Now consider a short-range order existing which is identical with the long-range order. That would mean if we pick up any three consecutive spins, two will be up, one down. Bragg-Williams approximation means short-range order is identical with long-range order, applied to a lattice of spins, in general. Row of spins is a lattice of one dimension.

Now let us imagine an arbitrary lattice, with each up spin assigned a value one and a down spin a value minus one, with an unspecified long-range order parameter defined as above by $L = \frac{1}{N}\sum_i \sigma_i$, where σ_i is i-th spin, N being total number of spins. L can vary from minus one to one. $N = N_+ + N_-$, where N_+ is the number of up spins, N_- is the number of down spins. $L = \frac{1}{N}(N_+ - N_-)$. As a result, $N_+ = \frac{N}{2}(1 + L)$ and $N_- = \frac{N}{2}(1 - L)$. Magnetisation or, net magnetic moment, M is $\mu\sum_i \sigma_i$ or, $\mu(N_+ - N_-)$ or, μNL , $M_{max} = \mu N$. $\frac{M}{M_{max}} = L$. $\frac{M}{M_{max}}$ is referred to as reduced magnetisation. Moreover, the Ising Hamiltonian,[79], for the lattice of spins, setting μ to one, is $-\epsilon\sum_{n,n}\sigma_i\sigma_j - H\sum_i \sigma_i$, where n.n refers to nearest neighbour pairs. The difference ΔE of energy if we flip an up spin to down spin is, [80], $2\epsilon\gamma\bar{\sigma} + 2H$, where γ is the number of nearest neighbours of a spin. According to Boltzmann principle, $\frac{N_-}{N_+}$ equals $exp(-\frac{\Delta E}{k_B T})$, [81]. In the Bragg-Williams approximation,[82], $\bar{\sigma} = L$, considered in the thermal average sense. Consequently,

$$\ln \frac{1+L}{1-L} = 2 \frac{\gamma\epsilon L + H}{k_B T} = 2 \frac{L + \frac{H}{\gamma\epsilon}}{\frac{T}{\gamma\epsilon/k_B}} = 2 \frac{L + c}{\frac{T}{T_c}} \quad (1)$$

where, $c = \frac{H}{\gamma\epsilon}$, $T_c = \gamma\epsilon/k_B$, [83]. $\frac{T}{T_c}$ is referred to as reduced temperature.

Plot of L vs $\frac{T}{T_c}$ or, reduced magnetisation vs. reduced temperature is used as reference curve. In the presence of magnetic field, $c \neq 0$, the curve bulges outward. Bragg-Williams is a Mean Field approximation. This approximation holds when number of neighbours interacting with a site is very large, reducing the importance of local fluctuation or, local order, making the long-range order or, average degree of freedom as the only degree of freedom of the lattice. To have a feeling how this approximation leads to matching between experimental and Ising

model prediction one can refer to FIG.12.12 of [80]. W. L. Bragg was a professor of Hans Bethe. Rudolf Peierls was a friend of Hans Bethe. At the suggestion of W. L. Bragg, Rudolf Peierls following Hans Bethe improved the approximation scheme, applying quasi-chemical method.

B. Bethe-peierls approximation in presence of four nearest neighbours, in absence of external magnetic field

In the approximation scheme which is improvement over the Bragg-Williams, [79],[80],[81],[82],[83], due to Bethe-Peierls, [84], reduced magnetisation varies with reduced temperature, for γ neighbours, in absence of external magnetic field, as

$$\frac{\ln \frac{\gamma}{\gamma-2}}{\ln \frac{factor-1}{factor^{\frac{\gamma-1}{\gamma}} - factor^{\frac{1}{\gamma}}}} = \frac{T}{T_c}; factor = \frac{\frac{M}{M_{max}} + 1}{1 - \frac{M}{M_{max}}}. \quad (2)$$

$\ln \frac{\gamma}{\gamma-2}$ for four nearest neighbours i.e. for $\gamma = 4$ is 0.693. For a snapshot of different kind of magnetisation curves for magnetic materials the reader is urged to give a google search "reduced magnetisation vs reduced temperature curve". In the following, we describe data s generated from the equation(1) and the equation(2) in the table, I, and curves of magnetisation plotted on the basis of those data s. BW stands for reduced temperature in Bragg-Williams approximation, calculated from the equation(1). BP(4) represents reduced temperature in the Bethe-Peierls approximation, for four nearest neighbours, computed from the equation(2). The data set is used to plot fig.1. Empty spaces in the table, I, mean corresponding point pairs were not used for plotting a line.

C. Bethe-peierls approximation in presence of four nearest neighbours, in presence of external magnetic field

In the Bethe-Peierls approximation scheme , [84], reduced magnetisation varies with reduced temperature, for γ neighbours, in presence of external magnetic field, as

$$\frac{\ln \frac{\gamma}{\gamma-2}}{\ln \frac{factor-1}{e^{\frac{2\beta H}{\gamma}} factor^{\frac{\gamma-1}{\gamma}} - e^{-\frac{2\beta H}{\gamma}} factor^{\frac{1}{\gamma}}}} = \frac{T}{T_c}; factor = \frac{\frac{M}{M_{max}} + 1}{1 - \frac{M}{M_{max}}}. \quad (3)$$

BW	BW(c=0.01)	BP(4, $\beta H = 0$)	reduced magnetisation
0	0	0	1
0.435	0.439	0.563	0.978
0.439	0.443	0.568	0.977
0.491	0.495	0.624	0.961
0.501	0.507	0.630	0.957
0.514	0.519	0.648	0.952
0.559	0.566	0.654	0.931
0.566	0.573	0.7	0.927
0.584	0.590	0.7	0.917
0.601	0.607	0.722	0.907
0.607	0.613	0.729	0.903
0.653	0.661	0.770	0.869
0.659	0.668	0.773	0.865
0.669	0.676	0.784	0.856
0.679	0.688	0.792	0.847
0.701	0.710	0.807	0.828
0.723	0.731	0.828	0.805
0.732	0.743	0.832	0.796
0.756	0.766	0.845	0.772
0.779	0.788	0.864	0.740
0.838	0.853	0.911	0.651
0.850	0.861	0.911	0.628
0.870	0.885	0.923	0.592
0.883	0.895	0.928	0.564
0.899	0.918		0.527
0.904	0.926	0.941	0.513
0.946	0.968	0.965	0.400
0.967	0.998	0.965	0.300
0.987		1	0.200
0.997		1	0.100
1	1	1	0

TABLE I. Reduced magnetisation vs reduced temperature data s for Bragg-Williams approximation, in absence of and in presence of magnetic field, $c = \frac{H}{\gamma\epsilon} = 0.01$, and Bethe-Peierls approximation in absence of magnetic field, for four nearest neighbours .

Derivation of this formula Ala [84] is given in the appendix of [7].

$\ln \frac{\gamma}{\gamma-2}$ for four nearest neighbours i.e. for $\gamma = 4$ is 0.693. For four neighbours,

$$\frac{0.693}{\ln \frac{factor-1}{e^{-\frac{2\beta H}{\gamma}} factor^{\frac{\gamma-1}{\gamma}} - e^{-\frac{2\beta H}{\gamma}} factor^{\frac{1}{\gamma}}}} = \frac{T}{T_c}; factor = \frac{\frac{M}{M_{max}} + 1}{1 - \frac{M}{M_{max}}}. \quad (4)$$

In the following, we describe data s in the table, II, generated from the equation(4) and curves of magnetisation plotted on the basis of those data s. BP(m=0.03) stands for reduced temperature in Bethe-Peierls approximation, for four nearest neighbours, in presence of a variable external magnetic field, H, such that $\beta H = 0.06$. calculated from the equation(4). BP(m=0.025) stands for reduced temperature in Bethe-Peierls approximation, for four nearest neighbours, in presence of a variable external magnetic field, H, such that $\beta H = 0.05$. calculated from the equation(4). BP(m=0.02) stands for reduced temperature

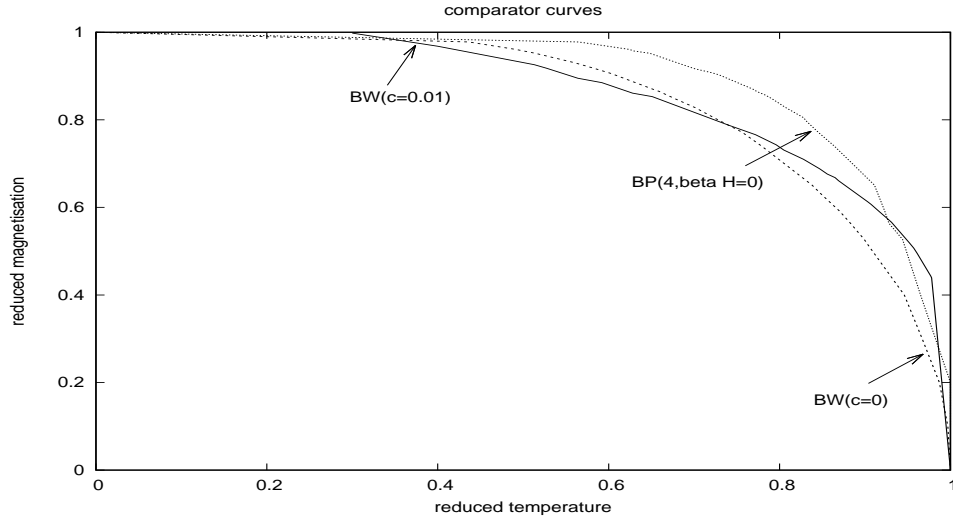


FIG. 1. Reduced magnetisation vs reduced temperature curves for Bragg-Williams approximation, in absence(dark) of and presence(inner in the top) of magnetic field, $c = \frac{H}{\gamma\epsilon} = 0.01$, and Bethe-Peierls approximation in absence of magnetic field, for four nearest neighbours (outer in the top).

in Bethe-Peierls approximation, for four nearest neighbours, in presence of a variable external magnetic field, H , such that $\beta H = 0.04$. calculated from the equation(4). BP(m=0.01) stands for reduced temperature in Bethe-Peierls approximation, for four nearest neighbours, in presence of a variable external magnetic field, H , such that $\beta H = 0.02$. calculated from the equation(4). BP(m=0.005) stands for reduced temperature in Bethe-Peierls approximation, for four nearest neighbours, in presence of a variable external magnetic field, H , such that $\beta H = 0.01$. calculated from the equation(4). The data set is used to plot fig.2. Empty spaces in the table, II, mean corresponding point pairs were not used for plotting a line.

BP(m=0.03)	BP(m=0.025)	BP(m=0.02)	BP(m=0.01)	BP(m=0.005)	reduced magnetisation
0	0	0	0	0	1
0.583	0.580	0.577	0.572	0.569	0.978
0.587	0.584	0.581	0.575	0.572	0.977
0.647	0.643	0.639	0.632	0.628	0.961
0.657	0.653	0.649	0.641	0.637	0.957
0.671	0.667		0.654	0.650	0.952
	0.716			0.696	0.931
0.723	0.718	0.713	0.702	0.697	0.927
0.743	0.737	0.731	0.720	0.714	0.917
0.762	0.756	0.749	0.737	0.731	0.907
0.770	0.764	0.757	0.745	0.738	0.903
0.816	0.808	0.800	0.785	0.778	0.869
0.821	0.813	0.805	0.789	0.782	0.865
0.832	0.823	0.815	0.799	0.791	0.856
0.841	0.833	0.824	0.807	0.799	0.847
0.863	0.853	0.844	0.826	0.817	0.828
0.887	0.876	0.866	0.846	0.836	0.805
0.895	0.884	0.873	0.852	0.842	0.796
0.916	0.904	0.892	0.869	0.858	0.772
0.940	0.926	0.914	0.888	0.876	0.740
	0.929			0.877	0.735
	0.936			0.883	0.730
	0.944			0.889	0.720
	0.945				0.710
	0.955			0.897	0.700
	0.963			0.903	0.690
	0.973			0.910	0.680
				0.909	0.670
	0.993			0.925	0.650
		0.976	0.942		0.651
	1.00				0.640
		0.983	0.946	0.928	0.628
		1.00	0.963	0.943	0.592
			0.972	0.951	0.564
			0.990	0.967	0.527
			1.00	0.964	0.513
				1.00	0.500
					0.400
					0.300
					0.200
					0.100
					0

TABLE II. Bethe-Peierls approx. in presence of little external magnetic fields

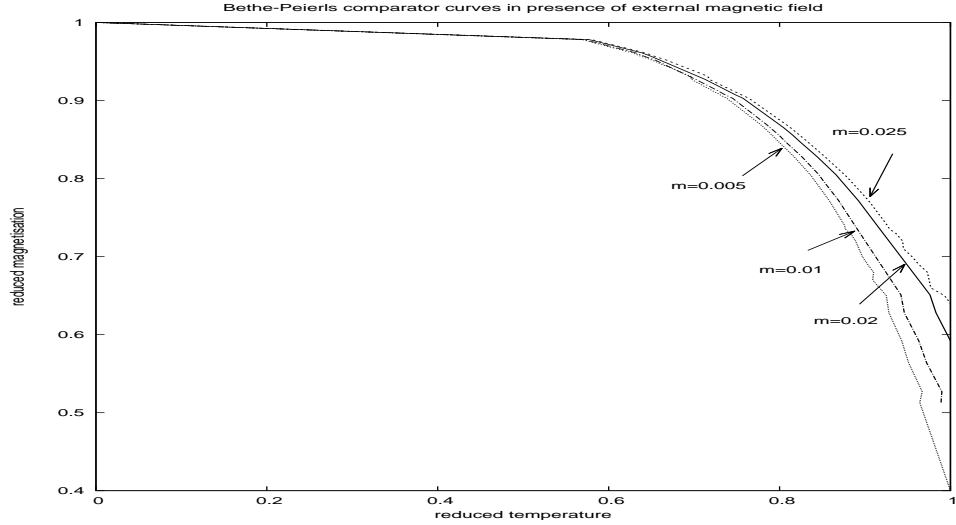


FIG. 2. Reduced magnetisation vs reduced temperature curves for Bethe-Peierls approximation in presence of little external magnetic fields, for four nearest neighbours, with $\beta H = 2m$.

A	B	C	D	E	F	G	H	I	J	K	L	M	N	O	P	Q	R	S	T	U	V	W	X	Y	Z
322	1192	722	578	247	487	627	913	65	213	606	643	1084	242	184	604	33	615	1316	439	37	163	696	6	65	97

TABLE III. The entries of Dictionary of American Family Names by Elsdon C. Smith along the English letters

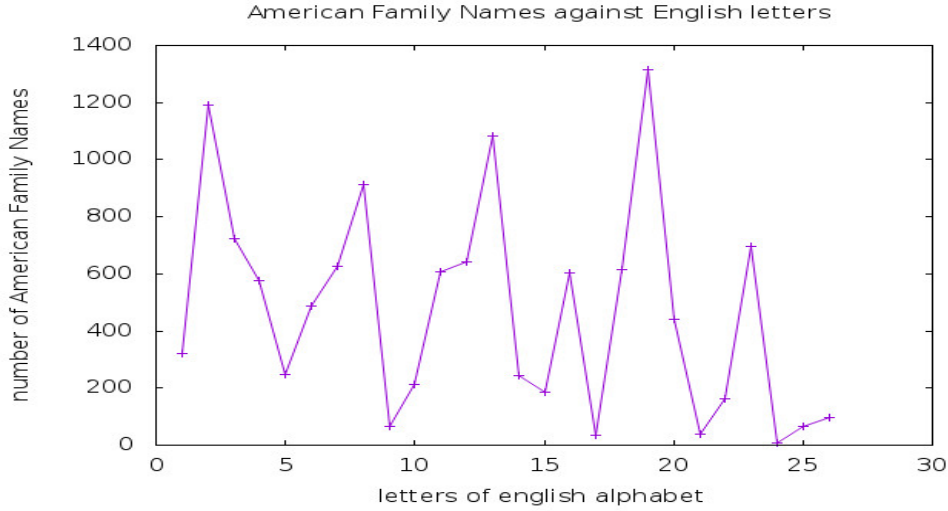


FIG. 3. The vertical axis is the number of entries of Dictionary of American Family Names by Elsdon C. Smith, [1], and the horizontal axis is the respective letters. Letters are represented by the sequence number in the alphabet or, dictionary sequence,[1].

III. ANALYSIS OF THE ENTRIES OF DICTIONARY OF AMERICAN FAMILY NAMES BY ELSDON C. SMITH

Counting one by one the entries of Dictionary of American Family Names by Elsdon C. Smith, [1], starting with different letters, leads us to the table, III. Highest number of entries, one thousand three hundred sixteen, starts with the letter S followed by entries numbering one thousand one hundred ninety two beginning with B, one thousand eighty four with the letter M etc. To visualise we plot the number of entries against respective letters in the dictionary sequence, [1], in the figure fig.3.

For the purpose of exploring graphical law, we assort the letters according to the number of entries, in the descending order, denoted by f and the respective rank, denoted by k . k is a positive integer starting from one. The lowest value of f is six. Hence we attach a limiting f equal to one. The corresponding rank, k , denoted as k_{lim} is twenty six. As a result both

k	lnk	lnk/ lnk_{lim}	f	lnf	lnf/ lnf_{max}	lnf/ lnf_{nmax}	lnf/ lnf_{2nmax}	lnf/ lnf_{3nmax}	lnf/ lnf_{4nmax}	lnf/ lnf_{5nmax}
1	0	0	1316	7.182	1	Blank	Blank	Blank	Blank	Blank
2	0.69	0.212	1192	7.083	0.986	1	Blank	Blank	Blank	Blank
3	1.10	0.337	1084	6.988	0.973	0.987	1	Blank	Blank	Blank
4	1.39	0.426	913	6.817	0.949	0.962	0.976	1	Blank	Blank
5	1.61	0.494	722	6.582	0.916	0.929	0.942	0.966	1	Blank
6	1.79	0.549	696	6.545	0.911	0.924	0.937	0.960	0.994	1
7	1.95	0.598	643	6.466	0.900	0.913	0.925	0.949	0.982	0.988
8	2.08	0.638	627	6.441	0.897	0.909	0.922	0.945	0.979	0.984
9	2.20	0.675	615	6.422	0.894	0.907	0.919	0.942	0.976	0.981
10	2.30	0.706	606	6.407	0.892	0.905	0.917	0.940	0.973	0.979
11	2.40	0.736	604	6.404	0.892	0.904	0.916	0.939	0.973	0.978
12	2.48	0.761	578	6.360	0.886	0.898	0.910	0.933	0.966	0.972
13	2.56	0.785	487	6.188	0.862	0.874	0.886	0.908	0.940	0.945
14	2.64	0.810	439	6.084	0.847	0.859	0.871	0.892	0.924	0.930
15	2.71	0.831	322	5.775	0.804	0.815	0.826	0.847	0.877	0.882
16	2.77	0.850	247	5.509	0.767	0.778	0.788	0.808	0.837	0.842
17	2.83	0.868	242	5.489	0.764	0.775	0.785	0.805	0.834	0.839
18	2.89	0.887	213	5.361	0.746	0.757	0.767	0.786	0.814	0.819
19	2.94	0.902	184	5.215	0.726	0.736	0.746	0.765	0.792	0.797
20	3.00	0.920	163	5.094	0.709	0.719	0.729	0.747	0.774	0.778
21	3.04	0.933	97	4.575	0.637	0.646	0.655	0.671	0.695	0.699
22	3.09	0.948	65	4.174	0.581	0.589	0.597	0.612	0.634	0.638
23	3.14	0.963	37	3.611	0.503	0.510	0.517	0.530	0.549	0.552
24	3.18	0.975	33	3.497	0.487	0.494	0.500	0.513	0.531	0.534
25	3.22	0.988	6	1.792	0.250	0.253	0.256	0.263	0.272	0.274
26	3.26	1	1	0	0	0	0	0	0	0

TABLE IV. The entries of Dictionary of American Family Names by Elsdon C. Smith: ranking, natural logarithm, normalisations

$\frac{lnf}{lnf_{max}}$ and $\frac{lnk}{lnk_{lim}}$ varies from zero to one. Then we tabulate in the adjoining table, IV and plot $\frac{lnf}{lnf_{max}}$ against $\frac{lnk}{lnk_{lim}}$ in the figure fig.4. We then ignore the letter with the highest of entries, tabulate in the adjoining table, IV and redo the plot, normalising the $lnfs$ with next-to-maximum $lnf_{nextmax}$, and starting from $k = 2$ in the figure fig.5. This program then we repeat up to $k = 6$, resulting in figures up to fig.9.

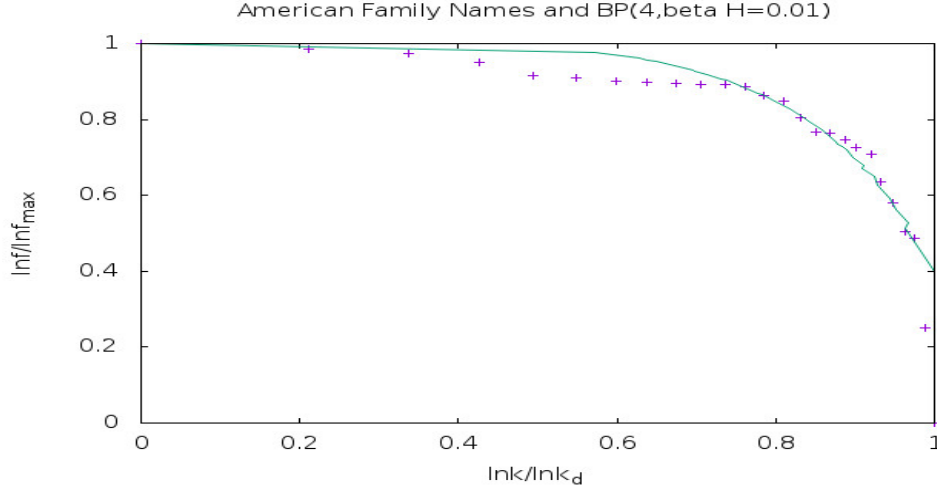


FIG. 4. The vertical axis is $\frac{\ln f}{\ln f_{max}}$ and the horizontal axis is $\frac{\ln k}{\ln k_{lim}}$. The + points represent the entries of Dictionary of American Family Names by Elsdon C. Smith with the fit curve being the Bethe-Peierls curve, $BP(4, \beta H = 0.01)$, in the presence of four nearest neighbours and little external magnetic field, $m = 0.005$ or, $\beta H = 0.01$.

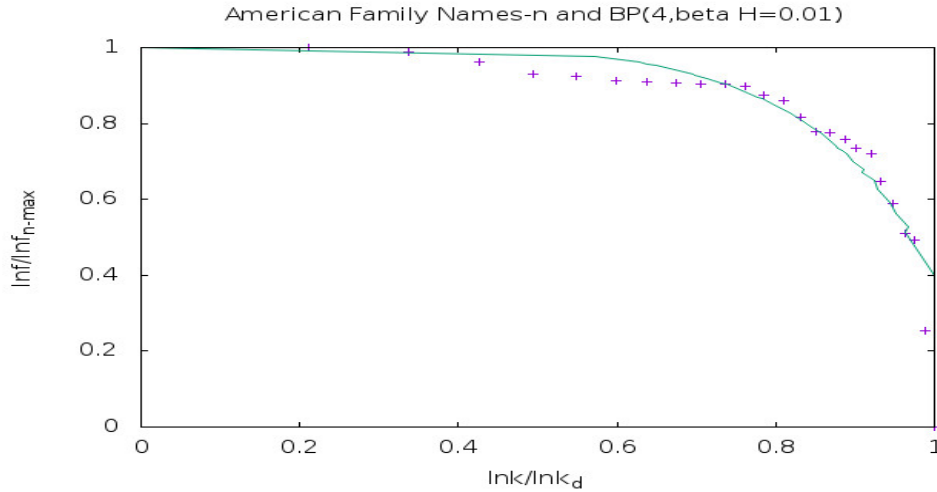


FIG. 5. The vertical axis is $\frac{\ln f}{\ln f_{n-max}}$ and the horizontal axis is $\frac{\ln k}{\ln k_{lim}}$. The + points represent the entries of Dictionary of American Family Names by Elsdon C. Smith with the fit curve being the Bethe-Peierls curve, $BP(4, \beta H = 0.01)$, in the presence of four nearest neighbours and little external magnetic field, $m = 0.005$ or, $\beta H = 0.01$.

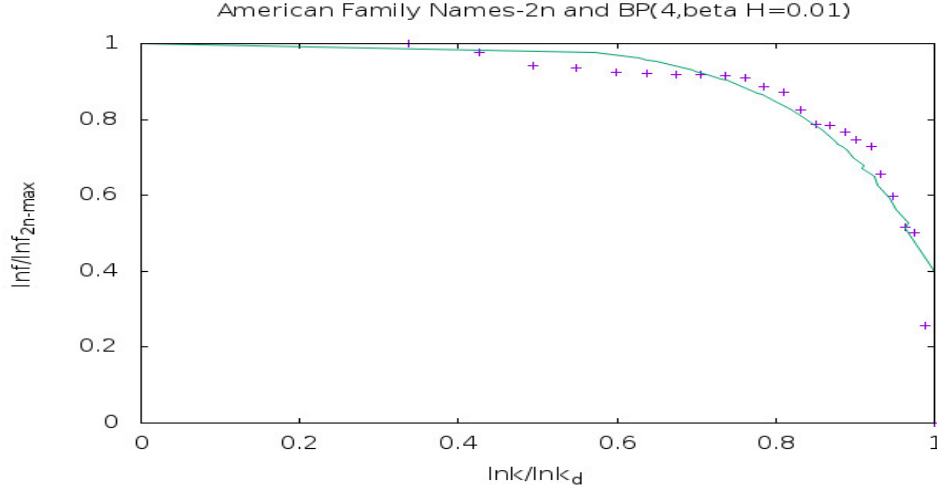


FIG. 6. The vertical axis is $\frac{\ln f}{\ln f_{2n-max}}$ and the horizontal axis is $\frac{\ln k}{\ln k_{lim}}$. The + points represent the entries of Dictionary of American Family Names by Elsdon C. Smith with the fit curve being the Bethe-Peierls curve, BP(4, $\beta H = 0.01$), in the presence of four nearest neighbours and little external magnetic field, $m = 0.005$ or, $\beta H = 0.01$.

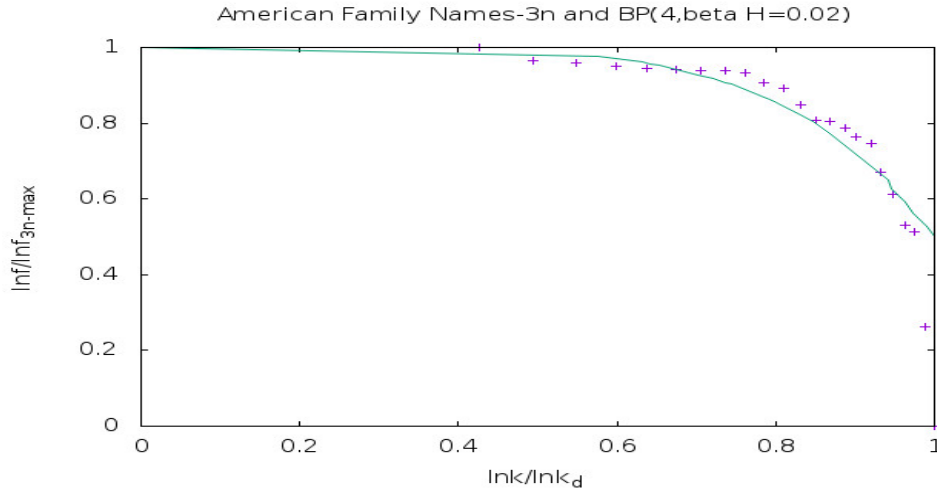


FIG. 7. The vertical axis is $\frac{\ln f}{\ln f_{3n-max}}$ and the horizontal axis is $\frac{\ln k}{\ln k_{lim}}$. The + points represent the entries of Dictionary of American Family Names by Elsdon C. Smith with the fit curve being the Bethe-Peierls curve, BP(4, $\beta H = 0.02$), in the presence of four nearest neighbours and little external magnetic field, $m = 0.01$ or, $\beta H = 0.02$.

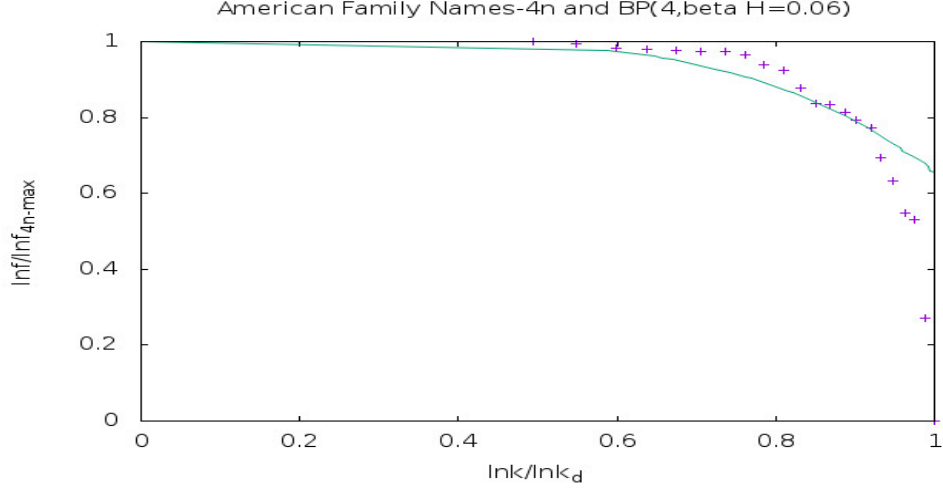


FIG. 8. The vertical axis is $\frac{\ln f}{\ln f_{4n-max}}$ and the horizontal axis is $\frac{\ln k}{\ln k_{lim}}$. The + points represent the entries of Dictionary of American Family Names by Elsdon C. Smith with the fit curve being the Bethe-Peierls curve, BP(4, $\beta H = 0.06$), in the presence of four nearest neighbours and little external magnetic field, $m = 0.03$ or, $\beta H = 0.06$.

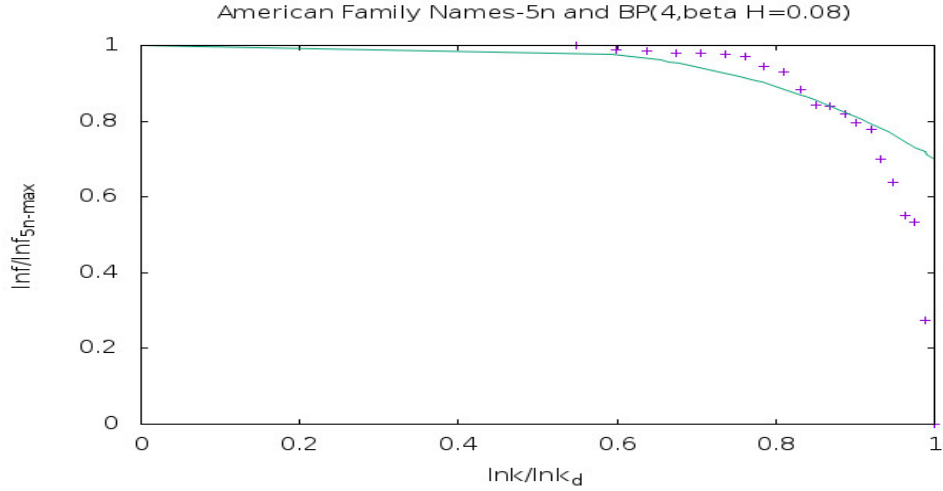


FIG. 9. The vertical axis is $\frac{\ln f}{\ln f_{5n-max}}$ and the horizontal axis is $\frac{\ln k}{\ln k_{lim}}$. The + points represent the entries of Dictionary of American Family Names by Elsdon C. Smith with the fit curve being the Bethe-Peierls curve, BP(4, $\beta H = 0.08$), in the presence of four nearest neighbours and little external magnetic field, $m = 0.04$ or, $\beta H = 0.08$.

A. conclusion

From the figures (fig.4-fig.9), we observe that there is a curve of magnetisation, behind the entries of Dictionary of American Family Names by Elsdon C. Smith, [1], This is the magnetisation curve, $BP(4, \beta H = 0.01)$, of the Ising model, in the presence of four nearest neighbours and little external magnetic field, $m = 0.005$ or, $\beta H = 0.01$. Moreover, the associated correspondence is,

$$\frac{\ln f}{\ln f_{n-max}} \longleftrightarrow \frac{M}{M_{max}}, \quad \ln k \longleftrightarrow T.$$

k corresponds to temperature in an exponential scale, [85].

This result should be seen in the background of the references, [10] and [6] respectively.

IV. ACKNOWLEDGMENT

The author would like to thank the State Central Library, West Bengal, Calcutta, to allow us to use Dictionary of American Family Names by Elsdon C. Smith, [1], in its premises. We have used gnuplot for plotting the figures in this paper.

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