Possibility of eD collisions in Pythia 8 / Angantyr

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Pythia 8/Angantyr

Where heavy ions meet proton-proton

- Heavy ion capabilities of PYTHIA8 recently introduced with the Angantyr model (arXiv:1806.10820 [hep-ph]).
- Development focused on High Energy Heavy Ion collisions (PbPb @ LHC etc.).
- Model can be extended to cover EiC/deuteron use cases.

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- This talk:
 - The Angantyr model.
 - 2 Results for pPb and PbPb.
 - Possible inclusion of deuterons.
 - Possible inclusion of γ^* .

Angantyr from helicopter perspective

Extending the PYTHIA8 MPI model to HI.
Only tuning to pp, add Glauber for nuclear geometry.



• Strategy for deuteron and γ^* to be added.

Glauber initial state

- Determine which nucleons are "wounded".
- Geometric picture only relies on pp cross section.



Glauber-Gribov colour fluctuations

- Cross section has EbE colour fluctuations.
- Parametrized in Angantyr, fitted to pp (total, elastic, diffractive).



- Simple model by Białas and Czyz.
- Wounded nucleons contribute equally to multiplicity in η .
- Originally: Emission function $F(\eta)$ fitted to data.



- $rac{dN}{d\eta} = F(\eta)$ (single wounded nucleon
- Angantyr: No fitting to HI data, but include model for emission function.
- Model fitted to reproduce pp case, high \sqrt{s} , can be retuned down to 10 GeV. Christian Bierlich (Lund/NBI) Pythia 8/Angantyr Aug 30, JLab 6 / 17

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 $rac{dN}{d\eta} = F(\eta) + F(-\eta)$ (pp)

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 $\frac{dN}{d\eta} = w_t F(\eta) + F(-\eta) \qquad (pA)$

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- $\frac{dN}{d\eta} = w_t F(\eta) + w_p F(-\eta) \tag{AA}$
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The emission function

- A schematic view of a pD collision. Contains 3 wounded nucleons.
- First two are a normal non-diffractive pp event.
- The second one is modelled as a single diffractive event.
- Generalizes to all pA and AA collisions.



Some results - pPb

• Centrality measures are delicate, but well reproduced.



Some results - pPb

• Multiplicity distributions well reproduced.



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Some results - PbPb

• Multiplicity distributions well reproduced.



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Some results - PbPb

• Spectra to a lesser degree, no collective effects so far.



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- Currently only Woods-Saxon distributed nuclei.
- Using GLISSANDO parametrization for deuteron Hulthen potential could be used.
- ... or external Glauber provider.
- Important test case for the model, a wealth of DAu data available.
- Missing: High energy pD data.

Future II: The electron

- Borrow from dipole models: Treat virtual photon as $q\bar{q}$ pair.
- Possible to construct a photon "Glauber".
- For large Q^2 splitting pertubatively calculable.
- For smaller Q^2 a hadronic component must be added.



• Only the (fluctuating) ep cross section is needed, everything else is in place.

Borrowing from DIPSY (Flensburg et al: arXiv:1103.4321 [hep-ph])

- The approach for cross sections is well tested.
- Wealth of data to test a new implementation on.



Conclusions

- Heavy ion model PYTHIA8 /Angantyr. Available in present version.
- Works well for large nuclei, high energies.
- Lower energies: Need pp data for tuning.
- Directly extendable to deuterons (on the TODO list).
- Virtual photon can be included through dipole formalism.

The DIPSY dipole cascade (Flensburg et al: arXiv:1103.4321 [hep-ph])

- Dipole models have been very succesful for eA collisions.
- DIPSY is a dipole cascade \Rightarrow builds up initial states event by event. Dipole evolution in Impact Parameter Space and rapiditY.
- Been applied to: pp, ep, pA, AA, eA.

$$\frac{dP}{dY} = \frac{3\alpha_s}{2\pi^2} d^2 \vec{z} \frac{(\vec{x} - \vec{y})^2}{(\vec{x} - \vec{z})^2 (\vec{z} - \vec{y})^2}, \ f_{ij} = \frac{\alpha_s^2}{8} \left[\log\left(\frac{(\vec{x}_i - \vec{y}_j)^2 (\vec{y}_i - \vec{x}_j)^2}{(\vec{x}_i - \vec{x}_j)^2 (\vec{y}_i - \vec{y}_j)^2}\right) \right]^2$$



Full event generator using Ariadne FS cascade + Pythia hadronization.

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Cross sections, MPIs and unitarisation

• Eikonal approximation \Rightarrow unitarized amplitude:

$$T \equiv -iA_{el} = 1 - \exp(-\sum f_{ij})$$

• Good–Walker formalism allows for diffraction. $\frac{d\sigma_{tot}}{d^2b} = 2 \langle T \rangle_{t,p}, \quad \frac{d\sigma_{el}}{d^2b} = \langle T \rangle_{t,p}^2, \quad \frac{d\sigma_{SD,(p|t)}}{d^2b} = \left\langle \langle T \rangle_{(t|p)}^2 \right\rangle_{(p|t)} - \left\langle T \right\rangle_{p,t}^2$ $\frac{d\sigma_{DD}}{d^2b} = \left\langle T^2 \right\rangle_{p,t} - \left\langle \langle T \rangle_t^2 \right\rangle_p - \left\langle \langle T \rangle_p^2 \right\rangle_t + \left\langle T \right\rangle_{p,t}^2$

- Includes MPIs by construction.
- Diffractive excitation determined by fluctuations in cascade.
- Not as precise as fine tuned PDFs.
- Not obvious how to include a signal process (everything is fluctuations).