

Pythia, DIPSY and Angantyr – past, present and future

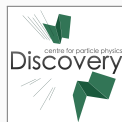
with focus on collectivity and heavy ion physics

Christian Bierlich, bierlich@thep.lu.se

University of Copenhagen

Lund University

March 27 2019, ALICE mini-workshop on MB & MC



Thank you ALICE, for all the beautiful data!

- Tales from the belly of the MC generators: DIPSY and Pythia.
 - Generators with focus on soft QCD.
 - Both capable of colliding HI beams.
 - Extensive use of Lund strings.
- Both are generators without QGP. Idea is:
 - Most observables in pp can be explained without QGP.
 - Collectivity = small effect, added as correction.
 - Long term: What happens when we extrapolate to AA?
- This talk:
 1. Past: The basic formalism, and results \approx 2014
 2. Present: Ropes and shoving in Pythia, Angantyr and possibilities.
 3. Future: Many prospects – three things I work on right now.



- Obviously very well known.

- Obviously very well known.

Information [References \(417\)](#) [Citations \(10001\)](#) [Files](#) [Plots](#)

PYTHIA 6.4 Physics and Manual

Torbjorn Sjostrand (Lund U., Dept. Theor. Phys.), Stephen Mrenna, Peter Z. Skands (Fermilab)

Mar 2006 - 576 pages

JHEP 0605 (2006) 026
DOI: [10.1088/1126-6708/2006/05/026](https://doi.org/10.1088/1126-6708/2006/05/026)
FERMILAB-PUB-06-052-CD-T, LU-TP-06-13
e-Print: [hep-ph/0603175](https://arxiv.org/abs/hep-ph/0603175) | [PDF](#)

- Here: Soft QCD (MPI model) + CR + strings.
- MPIs crucial for high energy pp collisions.

- Several partons from the PDF.
- Hard subcollisions with $2 \rightarrow 2$ ME:

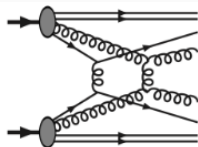


Figure T. Sjöstrand

$$\frac{d\sigma_{2 \rightarrow 2}}{dp_{\perp}^2} \propto \frac{\alpha_s^2(p_{\perp}^2)}{p_{\perp}^4} \rightarrow \frac{\alpha_s^2(p_{\perp}^2 + p_{\perp 0}^2)}{(p_{\perp}^2 + p_{\perp 0}^2)^2}.$$

- Momentum conservation and PDF scaling.
- Ordered emissions: $p_{\perp 1} > p_{\perp 2} > p_{\perp 4} > \dots$ from:

$$\mathcal{P}(p_{\perp} = p_{\perp i}) = \frac{1}{\sigma_{nd}} \frac{d\sigma_{2 \rightarrow 2}}{dp_{\perp}} \exp \left[- \int_{p_{\perp}}^{p_{\perp i-1}} \frac{1}{\sigma_{nd}} \frac{d\sigma}{dp'_{\perp}} dp'_{\perp} \right]$$

- Number distribution narrower than Poissonian (momentum and flavour rescaling).

Color reconnection

- Many partonic subcollisions \Rightarrow Many hadronizing strings.
- But! $N_c = 3$, not $N_c = \infty$ gives interactions.
- Easy to merge low- p_{\perp} systems, hard to merge two hard- p_{\perp} .

$$\mathcal{P}_{\text{merge}} = \frac{(\gamma p_{\perp 0})^2}{(\gamma p_{\perp 0})^2 + p_{\perp}^2}$$

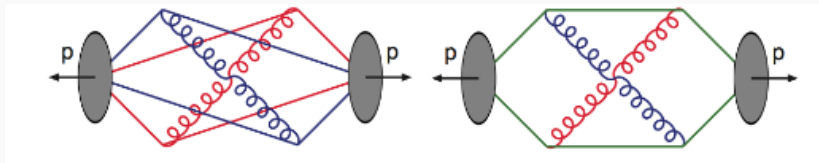


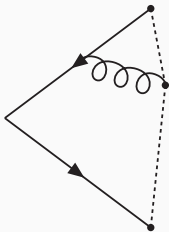
Figure T. Sjöstrand

- Actual merging decided by minimization of:

$$\lambda = \sum_{\text{dipoles}} \log(1 + \sqrt{2}E/m_0)$$

The Lund String (80's: Andersson et al. Z.Phys. C3 (1980) 223, Z.Phys. C20 (1983) 317)

- Non-perturbative phase of final state.
- Confined colour fields \approx *strings* with tension $\kappa \approx 1$ GeV/fm.



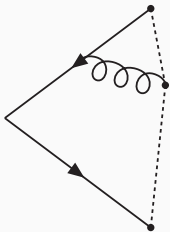
The Lund String (80's: Andersson et al. Z.Phys. C3 (1980) 223, Z.Phys. C20 (1983) 317)

- Non-perturbative phase of final state.
- Confined colour fields \approx strings with tension $\kappa \approx 1$ GeV/fm.
- *Breaking*/tunneling with $\mathcal{P} \propto \exp\left(-\frac{\pi m_{\perp}^2}{\kappa}\right)$ gives hadrons.

Lund symmetric fragmentation function

$$f(z) \propto z^{-1}(1-z)^a \exp\left(\frac{-bm_{\perp}}{z}\right).$$

a and b related to total multiplicity.



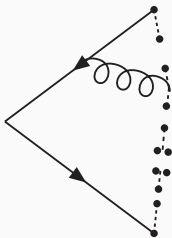
The Lund String (80's: Andersson et al. Z.Phys. C3 (1980) 223, Z.Phys. C20 (1983) 317)

- Non-perturbative phase of final state.
- Confined colour fields \approx *strings* with tension $\kappa \approx 1$ GeV/fm.
- *Breaking*/tunneling with $\mathcal{P} \propto \exp\left(-\frac{\pi m_{\perp}^2}{\kappa}\right)$ gives hadrons.

Lund symmetric fragmentation function

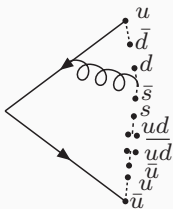
$$f(z) \propto z^{-1}(1-z)^a \exp\left(\frac{-bm_{\perp}}{z}\right).$$

a and b related to total multiplicity.



The Lund String (80's: Andersson et al. Z.Phys. C3 (1980) 223, Z.Phys. C20 (1983) 317)

- Non-perturbative phase of final state.
- Confined colour fields \approx strings with tension $\kappa \approx 1$ GeV/fm.
- *Breaking*/tunneling with $\mathcal{P} \propto \exp\left(-\frac{\pi m_{\perp}^2}{\kappa}\right)$ gives hadrons.



Lund symmetric fragmentation function

$$f(z) \propto z^{-1}(1-z)^a \exp\left(\frac{-bm_{\perp}}{z}\right).$$

a and b related to total multiplicity.

Light flavour determination

$$\rho = \frac{\mathcal{P}_{\text{strange}}}{\mathcal{P}_{u \text{ or } d}}, \xi = \frac{\mathcal{P}_{\text{diquark}}}{\mathcal{P}_{\text{quark}}}$$

Related to κ by Schwinger equation.

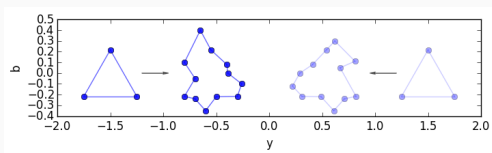
The DIPSY model (Flensburg et al. arXiv:1103.4321 [hep-ph])

- A very different view on MPIs, built on Mueller dipole model
(Mueller and Patel arXiv:hep-ph/9403256).
- Proton structure built up dynamically from dipole splittings:

Model implemented as a MC event generator

Dipole evolution in **I**mpact **P**arameter **S**pace and rapidity **Y**.

$$\frac{dP}{dY} = \frac{3\alpha_s}{2\pi^2} d^2\vec{z} \frac{(\vec{x} - \vec{y})^2}{(\vec{x} - \vec{z})^2(\vec{z} - \vec{y})^2}, \quad f_{ij} = \frac{\alpha_s^2}{8} \left[\log \left(\frac{(\vec{x}_i - \vec{y}_j)^2(\vec{y}_i - \vec{x}_j)^2}{(\vec{x}_i - \vec{x}_j)^2(\vec{y}_i - \vec{y}_j)^2} \right) \right]^2$$



- MPIs are included by construction.
- No PDFs (also: no quarks, no ME \Rightarrow few hard jets).

- Utilize knowledge of string positions – strings (p and q) still overlap.
- Combines into *multiplet* with effective string tension $\tilde{\kappa}$.

Effective string tension from the lattice

$$\kappa \propto C_2 \Rightarrow \frac{\tilde{\kappa}}{\kappa_0} = \frac{C_2(\text{multiplet})}{C_2(\text{singlet})}.$$

- Utilize knowledge of string positions – strings (p and q) still overlap.
- Combines into *multiplet* with effective string tension $\tilde{\kappa}$.

Effective string tension from the lattice

$$\kappa \propto C_2 \Rightarrow \frac{\tilde{\kappa}}{\kappa_0} = \frac{C_2(\text{multiplet})}{C_2(\text{singlet})}.$$

Easily calculable using $SU(3)$ recursion relations

$$\{p, q\} \otimes \vec{3} = \{p+1, q\} \oplus \{p, q+1\} \oplus \{p, q-1\}$$

$$\underbrace{\square \otimes \square \otimes \dots \otimes \square}_{\text{All anti-triplets}} \otimes \underbrace{\square \otimes \square \otimes \dots \otimes \square}_{\text{All triplets}}$$

- Utilize knowledge of string positions – strings (p and q) still overlap.
- Combines into *multiplet* with effective string tension $\tilde{\kappa}$.

Effective string tension from the lattice

$$\kappa \propto C_2 \Rightarrow \frac{\tilde{\kappa}}{\kappa_0} = \frac{C_2(\text{multiplet})}{C_2(\text{singlet})}.$$

Easily calculable using $SU(3)$ recursion relations

$$\{p, q\} \otimes \vec{3} = \{p+1, q\} \oplus \{p, q+1\} \oplus \{p, q-1\}$$

$$\underbrace{\begin{array}{c} \square \\ \square \end{array} \otimes \begin{array}{c} \square \\ \square \end{array} \otimes \dots \otimes \begin{array}{c} \square \\ \square \end{array}}_{\text{All anti-triplets}} \otimes \underbrace{\square \otimes \square \otimes \dots \otimes \square}_{\text{All triplets}}$$

- Transform to $\tilde{\kappa} = \frac{2p+q+2}{4}\kappa_0$ and $2N = (p+1)(q+1)(p+q+2)$.

CR collectivity is short range in rapidity

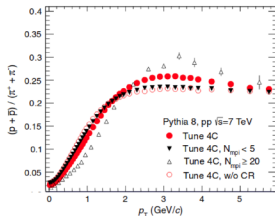
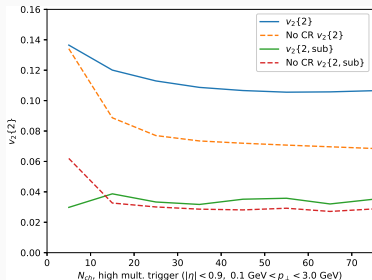
CR = short range in rapidity. Little effect on *inclusive* flavour composition.

CR collectivity is short range in rapidity

CR = short range in rapidity. Little effect on *inclusive* flavour composition.

Quantifying its contribution

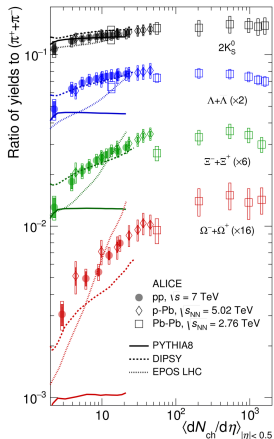
- Moves protons to measured phase space (Velasquez *et al.* PRL 111 (2013) 042001).
- Contribution to radial component, short range in y .



Contribution to $v_2\{2\}$ disappears: CR not long range.

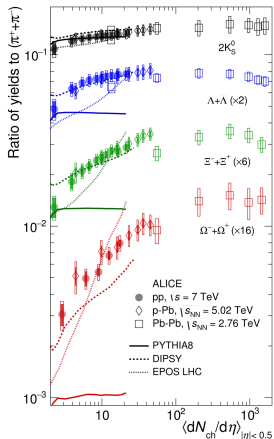
Strangeness enhancement

- A game of *density*.
- Good description of strangeness enhancement.

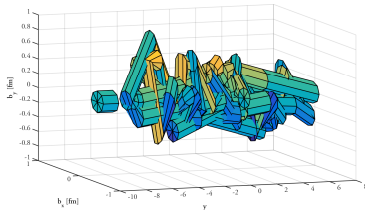


Strangeness enhancement

- A game of *density*.
- Good description of strangeness enhancement.



- DIPSY can make use of its impact parameter picture.





String shoving (CB, Gustafson, Lönnblad: 1612.05132, 1710.09725)

- Strings = interacting vortex lines.
- For $t \rightarrow \infty$, profile known from IQCD (Cea et al.: PRD89 (2014) no.9, 094505):

String showing (CB, Gustafson, Lönnblad: 1612.05132, 1710.09725)

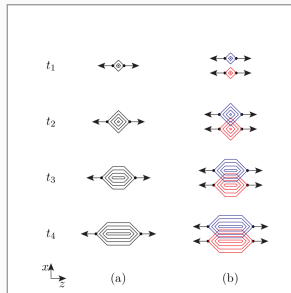
- Strings = interacting vortex lines.
- For $t \rightarrow \infty$, profile known from IQCD (Cea et al.: PRD89 (2014) no.9, 094505):

$$\mathcal{E}(r_{\perp}) = C \exp(-r_{\perp}^2/2R^2)$$

$$E_{int}(d_{\perp}) = \int d^2 r_{\perp} \mathcal{E}(\vec{r}_{\perp}) \mathcal{E}(\vec{r}_{\perp} - \vec{d}_{\perp})$$

$$f(d_{\perp}) = \frac{dE_{int}}{dd_{\perp}} = \frac{g\kappa d_{\perp}}{R^2} \exp\left(-\frac{d_{\perp}^2(t)}{4R^2}\right).$$

- Dominated by electric field $\rightarrow g = 1$.



String showing (CB, Gustafson, Lönnblad: 1612.05132, 1710.09725)

- Strings = interacting vortex lines.
- For $t \rightarrow \infty$, profile known from IQCD (Cea et al.: PRD89 (2014) no.9, 094505):

$$\mathcal{E}(r_{\perp}) = C \exp(-r_{\perp}^2/2R^2)$$

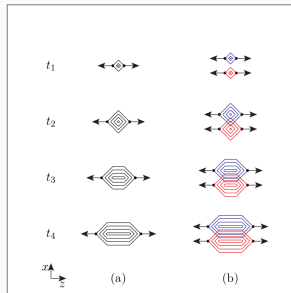
$$E_{int}(d_{\perp}) = \int d^2 r_{\perp} \mathcal{E}(\vec{r}_{\perp}) \mathcal{E}(\vec{r}_{\perp} - \vec{d}_{\perp})$$

$$f(d_{\perp}) = \frac{dE_{int}}{dd_{\perp}} = \frac{g\kappa d_{\perp}}{R^2} \exp\left(-\frac{d_{\perp}^2(t)}{4R^2}\right).$$

- Dominated by electric field $\rightarrow g = 1$.
- Reality:

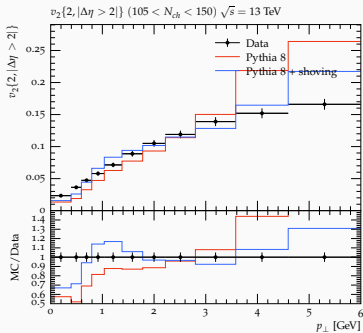
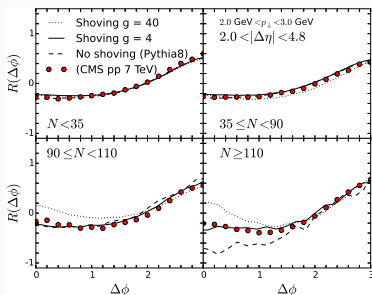
Type 1 Energy to destroy vacuum.

Type 2 Energy in current.



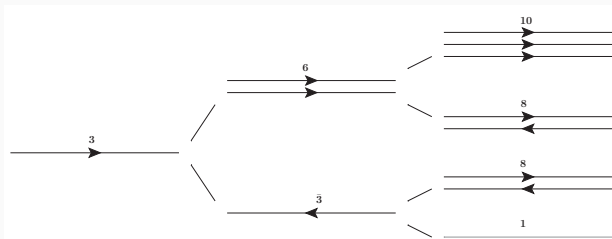
Some Results: shoving

- Reproduces the pp ridge with suitable choice of g parameter.
- Improved description of $v_2\{2|\Delta\eta| > 2, (p_\perp)$ at high multiplicity.

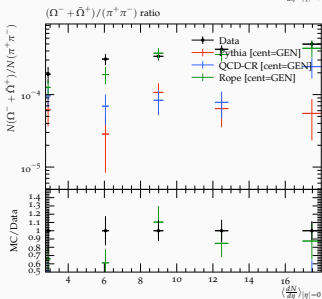
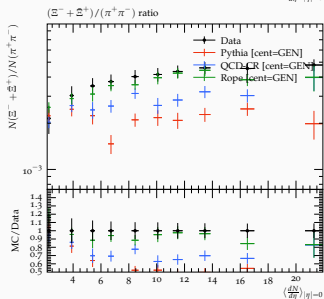
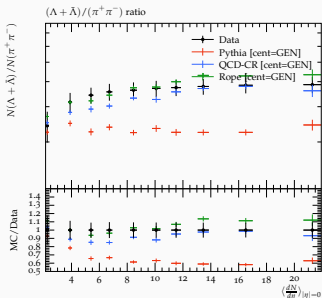
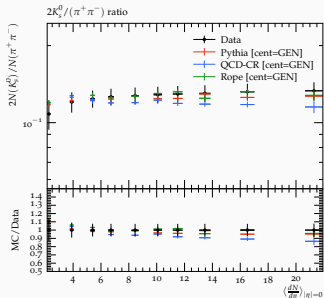


Ropes in Pythia

- The rope framework ported from DIPSY to Pythia.
- Requires space–time picture.
- Here: Overlapping 2D Gaussians (p mass distribution).
- All lower multiplets handled by CR.

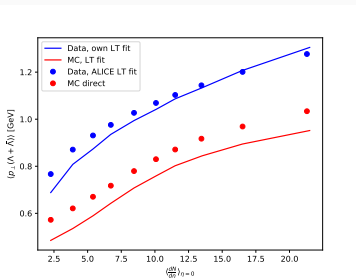
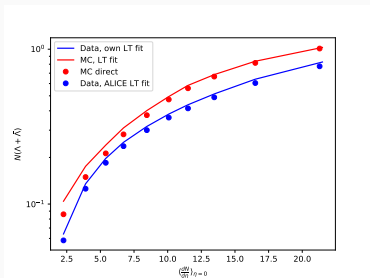


Results



An aside about Levy–Tsallis fits

- Extrapolated spectra are difficult to compare to!
- For Pythia: Yields matches the fit, $\langle p_{\perp} \rangle$ not.

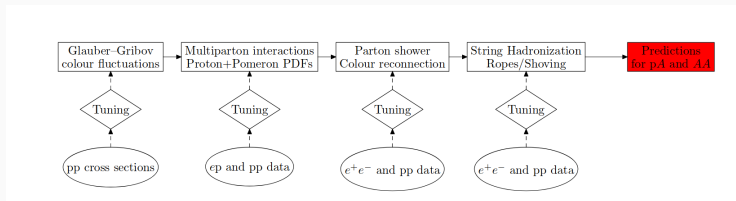


Take home message

MC: Don't rely on fits for average quantities when the spectrum is off.

Pythia still has problems describing this. Shoving could improve matters.

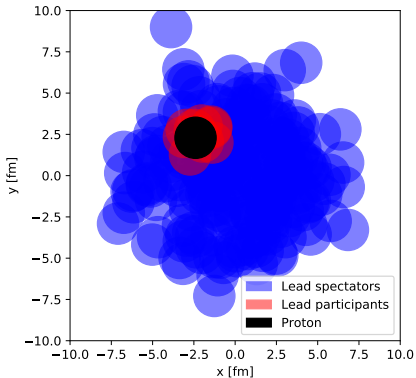
- Extending the Pythia MPI model to Heavy Ions.
 1. Only tuning to pp, add Glauber for nuclear geometry.



- Focus on cross section fluctuations (Glauber-Gribov) and correct handling of diffractive excitation.

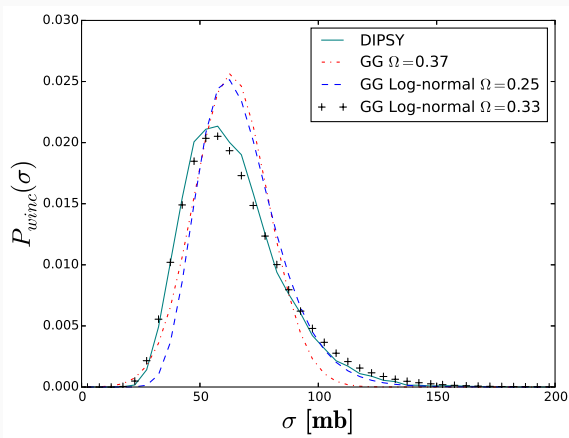
Glauber initial state

- Determine which nucleons are "wounded".
- Geometric picture only relies on pp cross section.



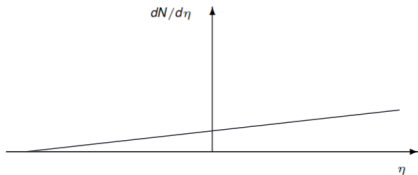
Glauber–Gribov colour fluctuations

- Cross section has EbE colour fluctuations.
- Parametrized in Angantyr, fitted to pp (total, elastic, diffractive).



Particle production: Wounded nucleons

- Simple model by Białas and Czyz.
- Wounded nucleons contribute equally to multiplicity in η .
- Originally: Emission function $F(\eta)$ fitted to data.

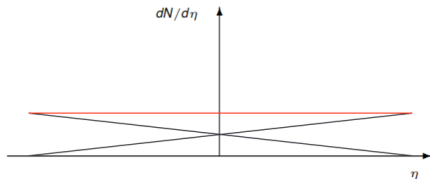


$$\frac{dN}{d\eta} = F(\eta) \quad (\text{single wounded nucleon})$$

- Angantyr: No fitting to HI data, but include model for emission function.
- Model fitted to reproduce pp case, high \sqrt{s} , can be retuned down to 10 GeV.

Particle production: Wounded nucleons

- Simple model by Białas and Czyz.
- Wounded nucleons contribute equally to multiplicity in η .
- Originally: Emission function $F(\eta)$ fitted to data.

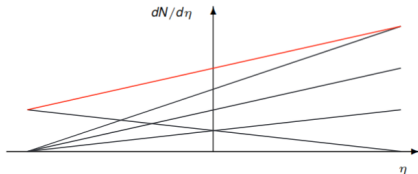


$$\frac{dN}{d\eta} = F(\eta) + F(-\eta) \quad (\text{pp})$$

- Angantyr: No fitting to HI data, but include model for emission function.
- Model fitted to reproduce pp case, high \sqrt{s} , can be retuned down to 10 GeV.

Particle production: Wounded nucleons

- Simple model by Białas and Czyz.
- Wounded nucleons contribute equally to multiplicity in η .
- Originally: Emission function $F(\eta)$ fitted to data.

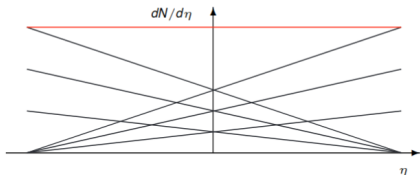


$$\frac{dN}{d\eta} = w_t F(\eta) + F(-\eta) \quad (\text{pA})$$

- Angantyr: No fitting to HI data, but include model for emission function.
- Model fitted to reproduce pp case, high \sqrt{s} , can be retuned down to 10 GeV.

Particle production: Wounded nucleons

- Simple model by Białas and Czyz.
- Wounded nucleons contribute equally to multiplicity in η .
- Originally: Emission function $F(\eta)$ fitted to data.

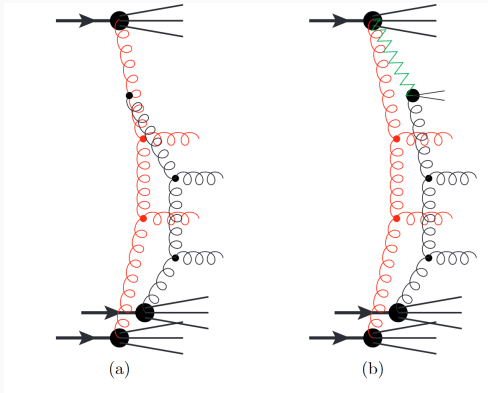


$$\frac{dN}{d\eta} = w_t F(\eta) + w_p F(-\eta) \quad (AA)$$

- Angantyr: No fitting to HI data, but include model for emission function.
- Model fitted to reproduce pp case, high \sqrt{s} , can be retuned down to 10 GeV.

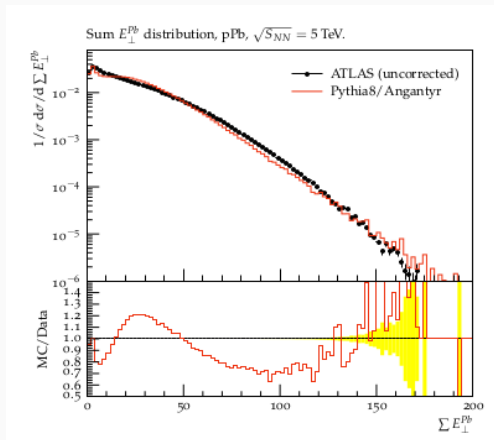
The emission function

- A schematic view of a pD collision. Contains 3 wounded nucleons.
- First two are a normal non-diffractive pp event.
- The second one is modelled as a single diffractive event.
- Generalizes to all pA and AA collisions.



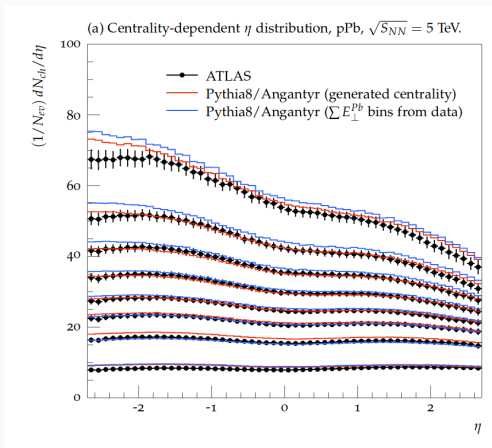
Some results - pPb

- Centrality measures are delicate, but well reproduced.



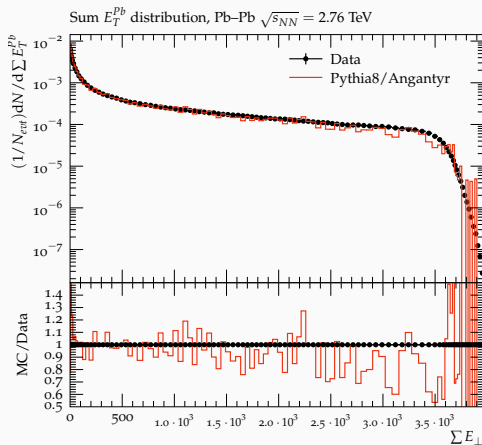
Some results - pPb

- Multiplicity distributions well reproduced.



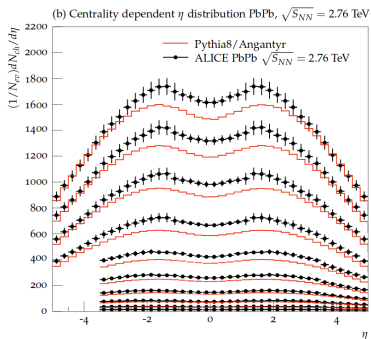
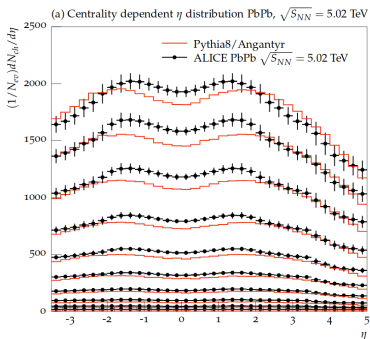
Some results - PbPb

- Centrality measure well reproduced.



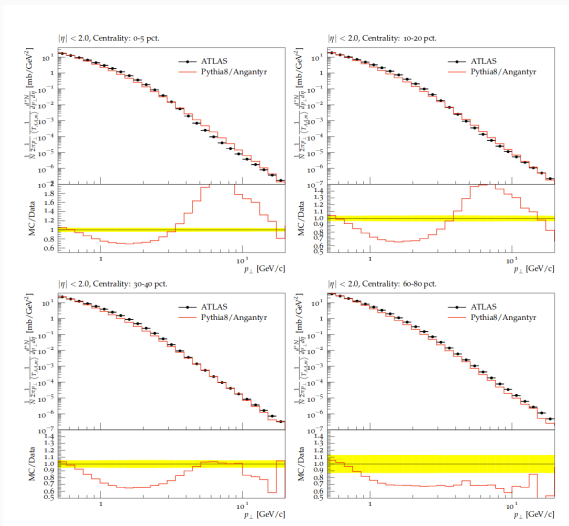
Some results - PbPb

- Multiplicity distributions well reproduced.
- Also XeXe (prediction) including up-tick.

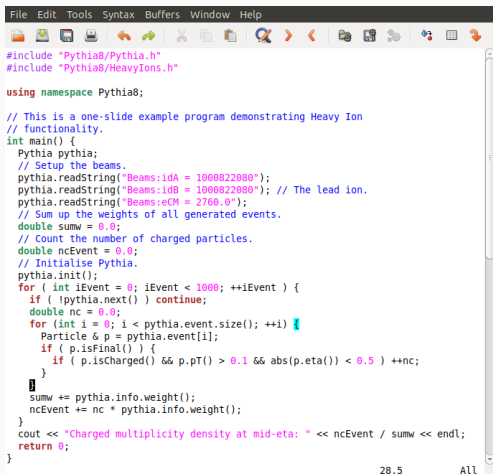


Some results - PbPb

- Spectra to a lesser degree, no collective effects so far.



Easy to use!



```
File Edit Tools Syntax Buffers Window Help
#include "Pythia8/Pythia.h"
#include "Pythia8/HeavyIons.h"

using namespace Pythia8;

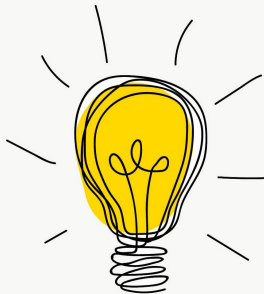
// This is a one-slide example program demonstrating Heavy Ion
// functionality.
int main() {
    Pythia pythia;
    // Setup the beams.
    pythia.readString("Beams:idA = 1000822080");
    pythia.readString("Beams:idB = 1000822080"); // The lead ion.
    pythia.readString("Beams:eCM = 2760.0");
    // Sum up the weights of all generated events.
    double sumw = 0.0;
    // Count the number of charged particles.
    double ncEvent = 0.0;
    // Initialise Pythia.
    pythia.init();
    for ( int iEvent = 0; iEvent < 1000; ++iEvent ) {
        if ( !pythia.next() ) continue;
        double nc = 0.0;
        for ( int i = 0; i < pythia.event.size(); ++i ) {
            Particle & p = pythia.event[i];
            if ( p.isFinal() ) {
                if ( p.ischarged() && p.pT() > 0.1 && abs(p.eta()) < 0.5 ) ++nc;
            }
            sumw += pythia.info.weight();
            ncEvent += nc * pythia.info.weight();
        }
        cout << "Charged multiplicity density at mid-eta: " << ncEvent / sumw << endl;
        return 0;
    }
}
```

28,5 All

- Fully integrated with Pythia.
- Internal or external ME's.
- Support for several nuclei.
- C++, Python interface distributed w. Pythia.
- Output: Rivet, HepMC, ROOT6 trees.

Two ideas

- MC implementation of models allows development of observables.
- Crucial: Physics can be "switched off", in a way it cannot in nature.



IDEA

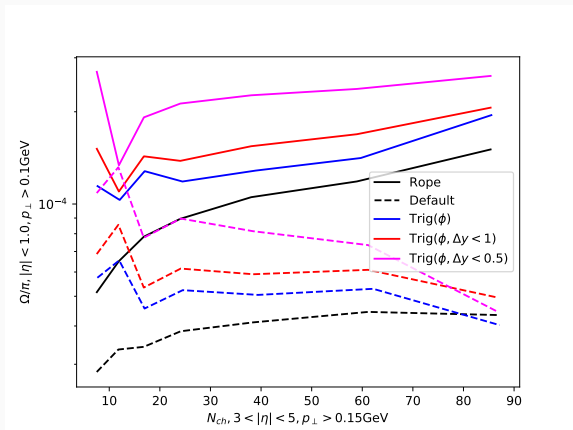
Can we get a better handle on strangeness enhancement? (CB, in

prep.)

- We can do better than inclusive rates.
- Accessing longitudinal (rapidity) structure: Correlation measurements.
- Consider ropes in a ϕ -triggered event.
 1. Even in e^+e^- we bias to more strange production,
 2. In pp we can assess the difference wrt. default strings.
 3. Moving closer to the ϕ production rapidity gives larger string tension.
- Statistics hungry analysis – something for HL-LHC?

Preliminary: pp @ 13 TeV (Pythia8 + ropes)

- Input for discussion:
 1. Sensible measurement?
 2. What does thermal models say?
 3. Can we remove the neighbor bias? (require neighbor etc.)



What about shoving and jets? (CB: 1901.07447)

- String dynamics ought to be universal.
- Consider now:
 1. Events with a Z -boson present.
 2. Events with Z +jet.
- $Z \rightarrow l^+ l^-$ not affected by shoving.
- Provides kinematics handle.

What about shoving and jets? (CB: 1901.07447)

- String dynamics ought to be universal.
- Consider now:
 1. Events with a Z -boson present.
 2. Events with Z +jet.
- $Z \rightarrow l^+l^-$ not affected by shoving.
- Provides kinematics handle.

Common statement:

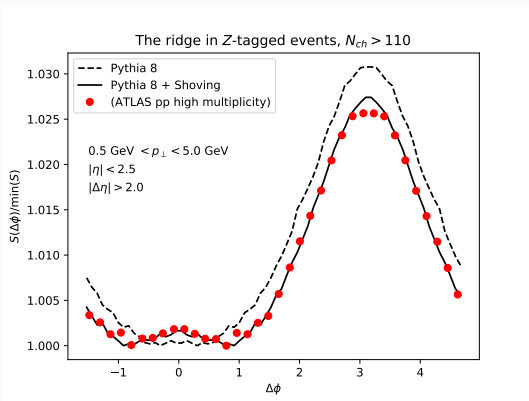
- ◇ FS interactions \rightarrow flow should also affect jets.
- ◇ The shoving model provides a framework to study such effects.
- ◇ This does not mean that shoving is the full story.

Try just a Z -boson

- The presence of a Z should not change the physics.
- It *can* introduce kinematical biases.
- Recently measured by ATLAS ([ATLAS-CONF-2017-068](#)).

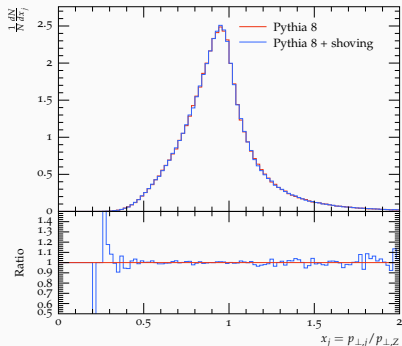
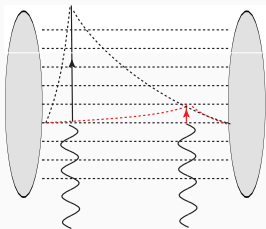
Try just a Z-boson

- The presence of a Z should not change the physics.
- It *can* introduce kinematical biases.
- Recently measured by ATLAS ([ATLAS-CONF-2017-068](#)).



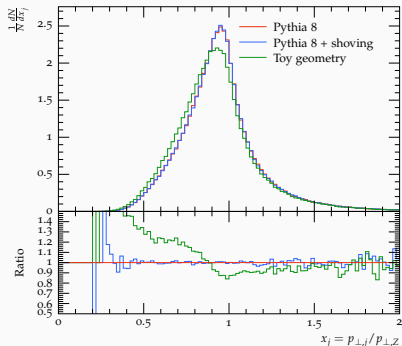
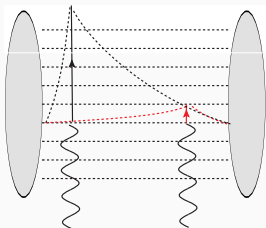
What is the effect of shoving?

- Nothing! Surprised?
- Of course not – the effect is geometrically suppressed.
- Toy geometry: Let jet hadronize "inside".
- Mimic the effect in AA collisions.

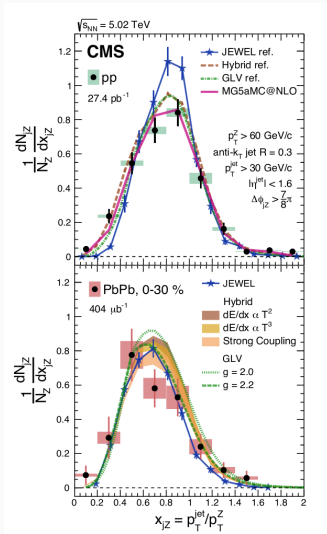


What is the effect of shoving?

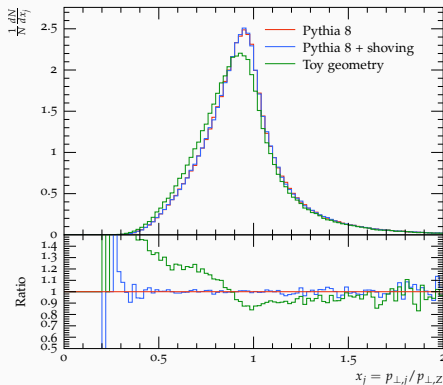
- Nothing! Surprised?
- Of course not – the effect is geometrically suppressed.
- Toy geometry: Let jet hadronize "inside".
- Mimic the effect in AA collisions.



Qualitative similarities (CMS: 1702.01060)



- Need better observables.
- Soft modifications on jet edge (large R).

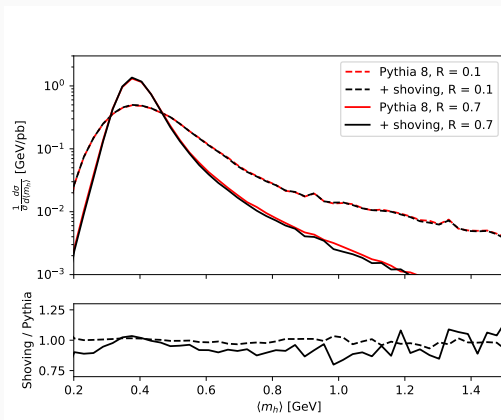


Hadrochemistry

- Hadrochemistry indirectly affected through basic string equations.
- Study inclusive quantities: Average hadron mass and total jet charge: $\langle m_h \rangle = \frac{1}{N_p} \sum_i^{N_p} m_{h,i}$, $Q_j = \sum_i^{N_p} q_{h,i}$

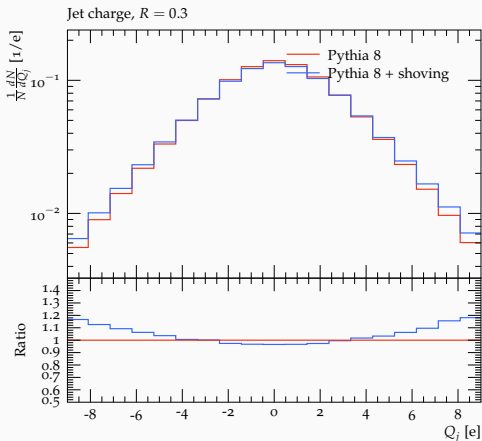
Hadrochemistry

- Hadrochemistry indirectly affected through basic string equations.
- Study inclusive quantities: Average hadron mass and total jet charge:
charge: $\langle m_h \rangle = \frac{1}{N_p} \sum_i^{N_p} m_{h,i}$, $Q_j = \sum_i^{N_p} q_{h,i}$



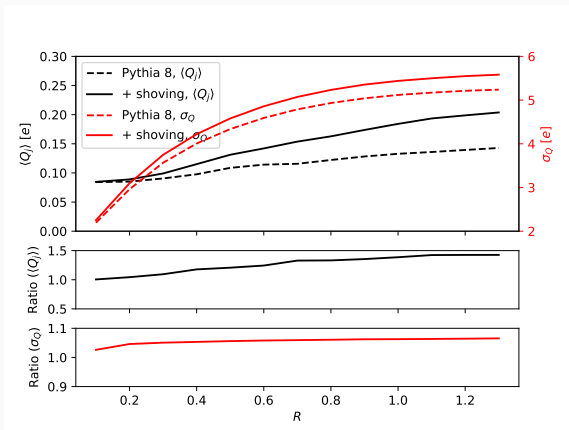
Hadrochemistry

- Hadrochemistry indirectly affected through basic string equations.
- Study inclusive quantities: Average hadron mass and total jet charge:
charge: $\langle m_h \rangle = \frac{1}{N_p} \sum_i^{N_p} m_{h,i}$, $Q_j = \sum_i^{N_p} q_{h,i}$



Hadrochemistry

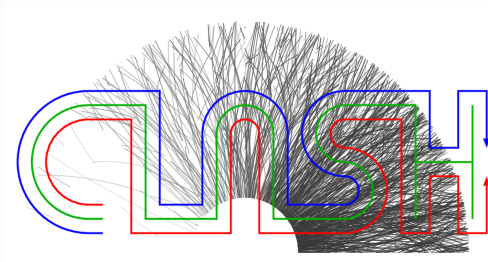
- Hadrochemistry indirectly affected through basic string equations.
- Study inclusive quantities: Average hadron mass and total jet charge: $\langle m_h \rangle = \frac{1}{N_p} \sum_i^{N_p} m_{h,i}$, $Q_j = \sum_i^{N_p} q_{h,i}$





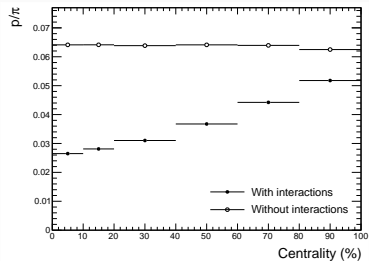
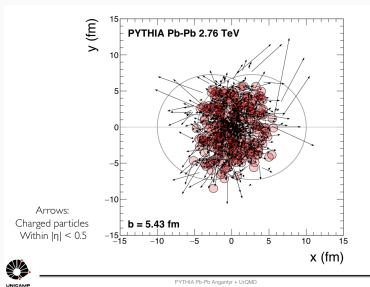
Obvious extensions

- Extension of ropes and shoving to pA and AA an obvious venue.
- Very active, still few results.
- In the framework of CLASH, see www.hep.lu.se/clash



Final state interactions in Angantyr (w. D. Chinellato & A. da Silva)

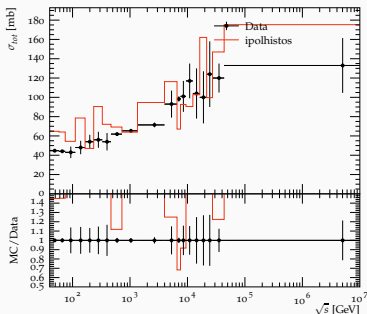
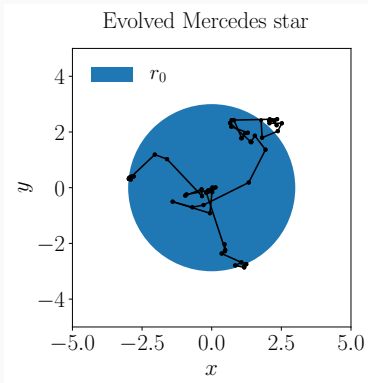
- Hadronic interactions in final state with URQMD.
- Hardon vertices from string model (Ferreres-Solé & Sjöstrand: 1808.04619).



- Definitely necessary ingredient for ropes in AA collisions.
- Many interesting prospects, resonances, effects on jets etc.
- Maybe possible to investigate pp as well.

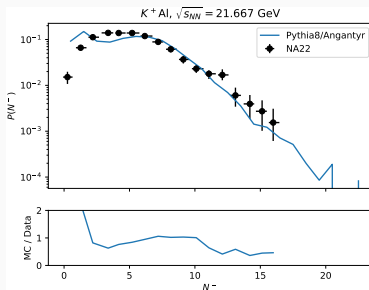
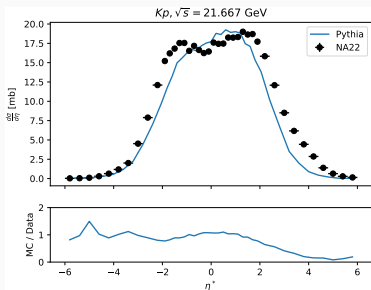
Improving the Pythia space–time picture (w. C.O. Rasmussen)

- Mueller dipole QCD (re: DIPSY) interesting features for space time model.
 1. Perturbative calculation (+ non-perturbative corrections).
 2. Structure parameters fitted to cross sections.
- Goal: Mueller dipoles \rightarrow space time information to Pythia MPIs.



The cosmic connection (w. A. Fedynitch, J. Koskinen & I. Storehaug)

- Uncertainties on cosmic data \leftarrow uncertainty on hadronic MCs.
- ... this is in turn limited by lack of good data.
- Use Pythia/Angantyr for cosmic data.
- Relies heavily on data from NA-22 days.



Opportunity for ALICE

Particle production with PID in pO at high energies.
Valuable cross-collaboration output.

Instead of a summary: The experimental wishlist

- Strangeness in pp:
 1. The ϕ is a sensitive probe.
 2. Triggered ratios \rightarrow higher resolution.
 3. p_{\perp} of heavy hadrons continues to puzzle us.
- Correlations:
 1. Continued efforts on precise flow measurements & SC's important for geometry. TH is lacking behind.
 2. Z-triggered jets a window to SS jet modifications.
 3. Shoving gives effects on jet chemistry.
- AA particle production:
 1. Centrality measures unfolded.
 2. Strong case for Oxygen collisions also from cosmic ray community.

Thank you!