Vetenskapsrådet

LUND UNIVERSITY

Collider Cross Talk
CERN, 4 August 2011

# Progress on Event Generation with PYTHIA 8


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Introduction and Overview

Improved Showers and ME matching

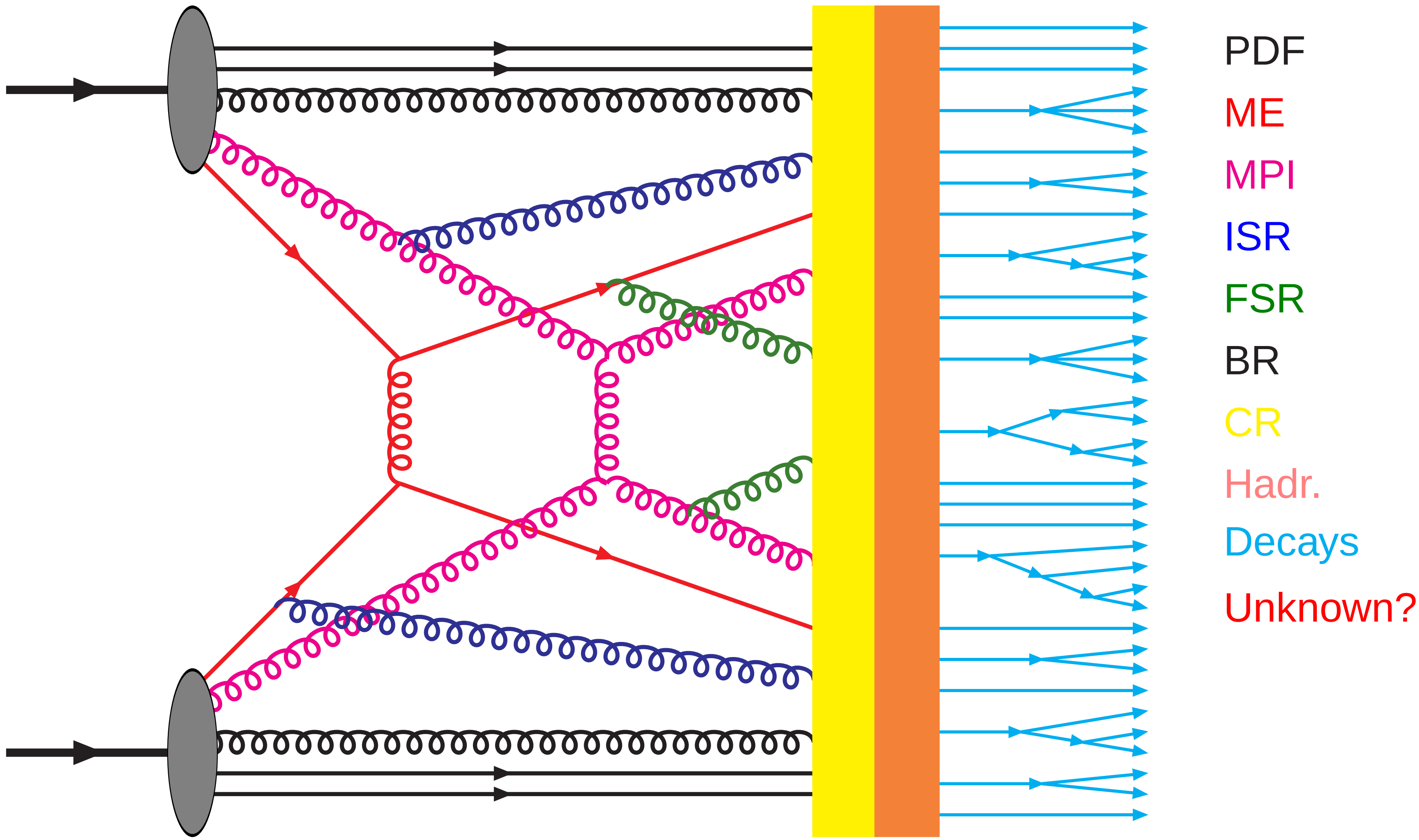
MultiParton Interactions

BSM Physics

Tunes and Comparisons with LHC Data

Summary and Outlook

# Introduction and Overview




General-purpose event generators: PYTHIA, HERWIG, SHERPA

PYTHIA size: ∼80,000 lines (Fortran in PYTHIA 6, C++ in PYTHIA 8)

## PYTHIA 8 ambition

- Meet experimental request for C++ code.
- Housecleaning ⇒ more homogeneous.
- More user-friendly (e.g. settings names).
- Better match to software frameworks (e.g. card files).
- More space for growth.
- Better interfaces to external standards.

## Reality

- Work begun autumn 2004.
- 3 years at CERN ⇒ good progress.
- First release autumn 2007.
- Since then: slower progress, requests lagging behind.
- Usage slowly taking off.

## Team members

Stefan Ask
Richard Corke
Stephen Mrenna
Peter Skands

## Contributors include

O. Alvestad, B. Bellenot,
R. Brun, L. Carloni,
N. Desai, N. Hod,
H. Hoeth, P. Ilten,
T. Kasemets, M. Kirsanov,
B. Lloyd, M. Montull,
A. Morsch, A. Naumann,
S. Navin, P. Newman,
S. Prestel, M. Sutton
MSTW, CTEQ, H1: PDFs
DELPHI, LHCb: D/B BRs
+ several bug reports & fixes

# Key differences between PYTHIA 6.4 and 8.1

Old features definitely removed include, among others:
- independent fragmentation
- mass-ordered showers

Features omitted so far include, among others:
- ep, $\gamma$p and $\gamma\gamma$ beam configurations
- some processes, notably Technicolor

New features, not found in 6.4:
- $\star$ fully interleaved $p_{\perp}$-ordered MPI + ISR + FSR evolution
- $\star$ richer mix of underlying-event processes ($\gamma$, J/$\psi$, DY, ...)
- $\star$ possibility for two selected hard interactions in same event
- $\star$ allow rescattering and $x$-dependent proton size in MPI framework
- $\star$ full hadron–hadron collision machinery for diffractive systems
- $\star$ several new processes, within and beyond SM
- $\star$ possibility to use one PDF set for hard process and another for rest
- $\star$ $\tau$ lepton polarization in production and decay
- $\star$ updated decay data and LO PDF sets

# Interleaved evolution

- Transverse-momentum-ordered parton showers for ISR and FSR
- MPI also ordered in $p_\perp$

Allows interleaved evolution for ISR, FSR and MPI

$$\frac{\mathrm{d}\mathcal{P}}{\mathrm{d}p_\perp} = \left(\frac{\mathrm{d}\mathcal{P}_{\mathrm{MPI}}}{\mathrm{d}p_\perp} + \sum \frac{\mathrm{d}\mathcal{P}_{\mathrm{ISR}}}{\mathrm{d}p_\perp} + \sum \frac{\mathrm{d}\mathcal{P}_{\mathrm{FSR}}}{\mathrm{d}p_\perp}\right) \times \exp\left(-\int_{p_\perp}^{p_{\perp\mathrm{max}}} \left(\frac{\mathrm{d}\mathcal{P}_{\mathrm{MPI}}}{\mathrm{d}p'_\perp} + \sum \frac{\mathrm{d}\mathcal{P}_{\mathrm{ISR}}}{\mathrm{d}p'_\perp} + \sum \frac{\mathrm{d}\mathcal{P}_{\mathrm{FSR}}}{\mathrm{d}p'_\perp}\right) \mathrm{d}p'_\perp\right)$$

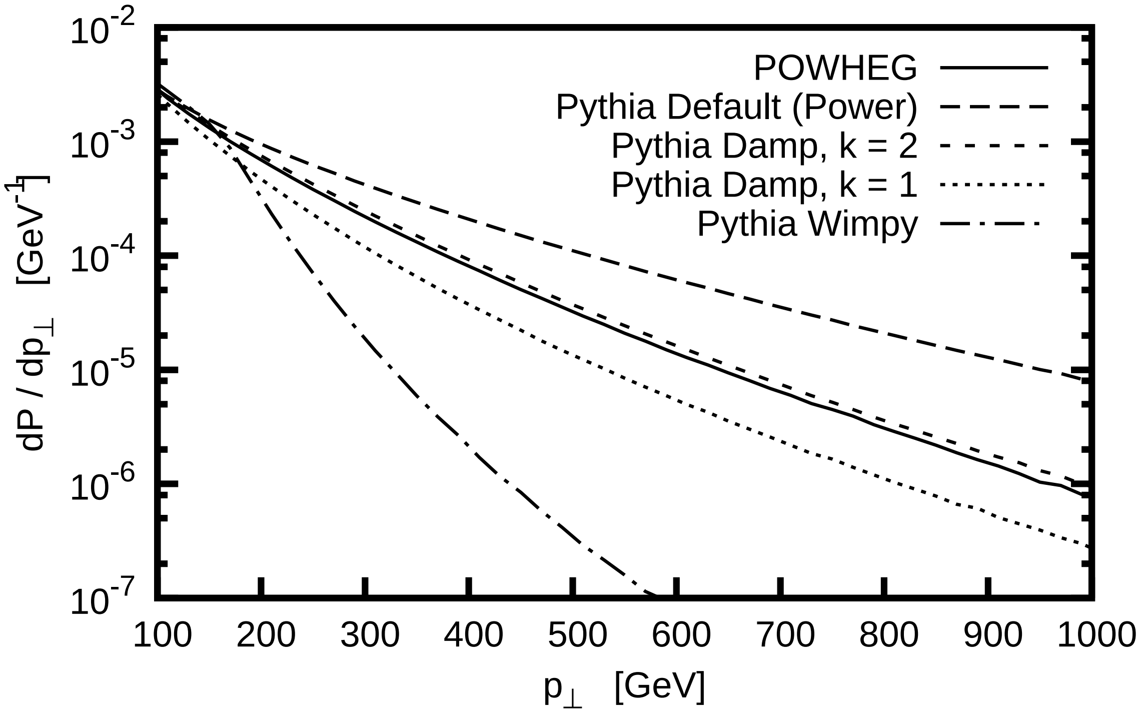
Ordered in decreasing $p_\perp$ using "Sudakov" trick.
Corresponds to increasing "resolution":
smaller $p_\perp$ fill in details of basic picture set at larger $p_\perp$.

Hybrid approach to shower recoils:

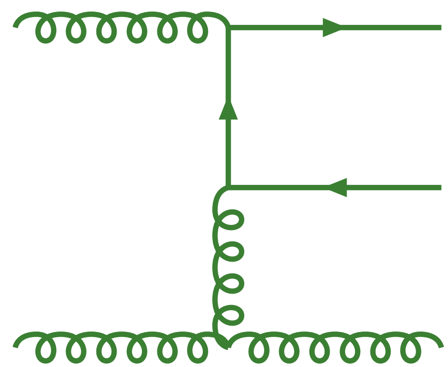
- FSR is dipole: nearest colour-connected neighbour
- ISR is traditional: whole hard-scattering system affected
  (since ISR dipole gives wrong answer e.g. for $p_{\perp Z}$)

# Shower matching to MEs: realistic hard default

Aim: provide better default shower behaviour at large $p_\perp$,
to bridge gap between “power” and “wimpy” showers.




$\mathrm{t\bar{t}}$ production



$$M^2 = m_{\perp\mathrm{t}}^2 = m_\mathrm{t}^2 + p_{\perp\mathrm{t}}^2$$

$$\frac{\mathrm{d}P_\mathrm{ISR}}{\mathrm{d}p_\perp^2} \propto \frac{1}{p_\perp^2}\frac{k^2M^2}{k^2M^2+p_\perp^2}$$ for coloured final state

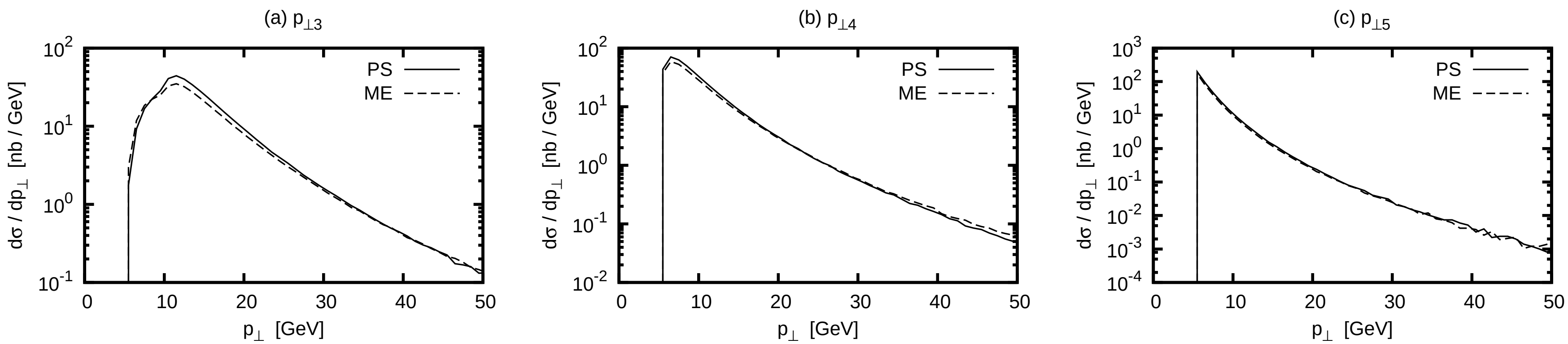
No dampening for uncoloured final state ($\mathrm{W^+W^-}$, . . . , SUSY).

R. Corke & TS, Eur. Phys. J. C69 (2010) 1

# Shower matching to MEs: realistic QCD default

Must avoid doublecounting for QCD jets:
shower starting scale = $p_\perp$ of hard $2 \to 2$ process.

Study how well the parton shower fills the phase space,
as prelude to full matching to $2 \to 3$ real-emission:



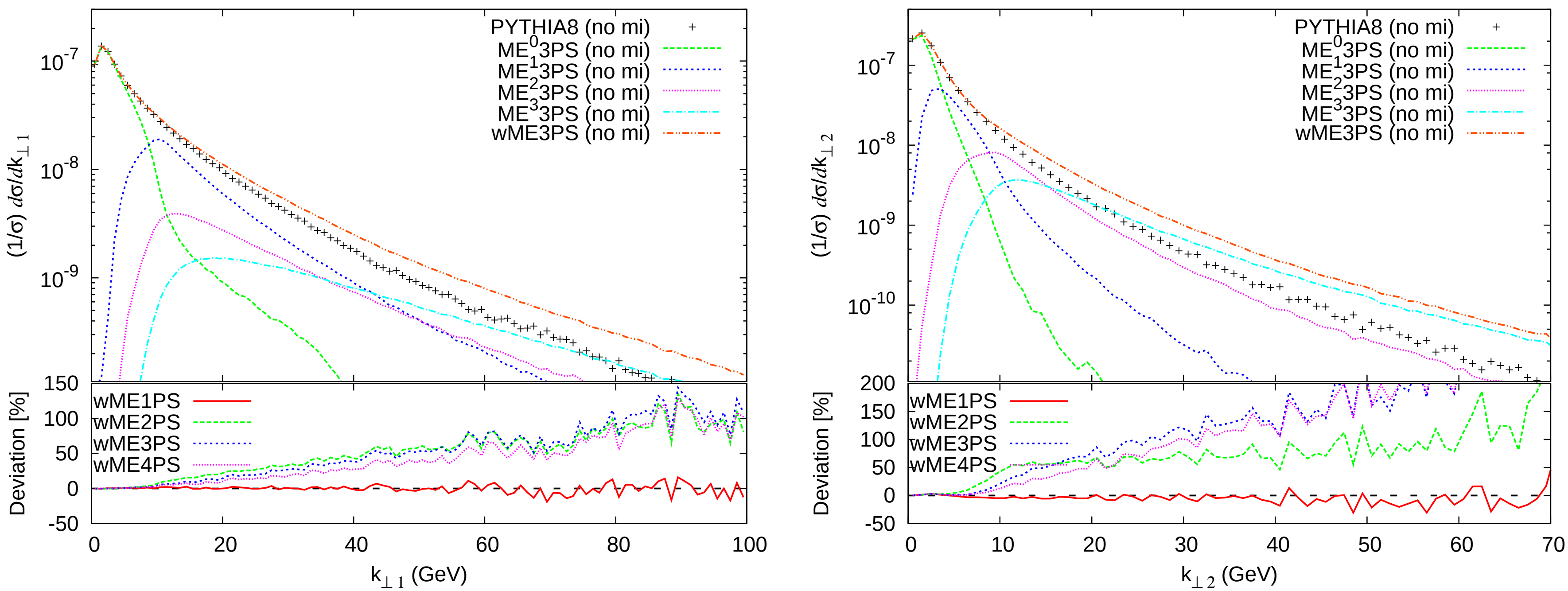

$p_{\perp 3}^{\text{min}} = 5.0$ GeV, $p_{\perp 4}^{\text{min}} = 5.0$ GeV, $R_{\text{sep}} = 0.10$

Obtain good qualitative agreement, best in soft and collinear regions,
but large region of phase space well described, and only corners bad.
No indication for needing a change in starting scale!
R. Corke & TS, JHEP 03 (2011) 032

# Shower matching to MEs: CKKW–L

Leif Lönnblad and Stefan Prestel work to include CKKW–L matching

- for jets, W/Z + jets, WW/WZ/ZZ + jets
- reconstructing ME "history" as viewed by PYTHIA 8 shower
- consistently taking into account effects of MPI




Article to appear any day now; code "soon" in PYTHIA 8;

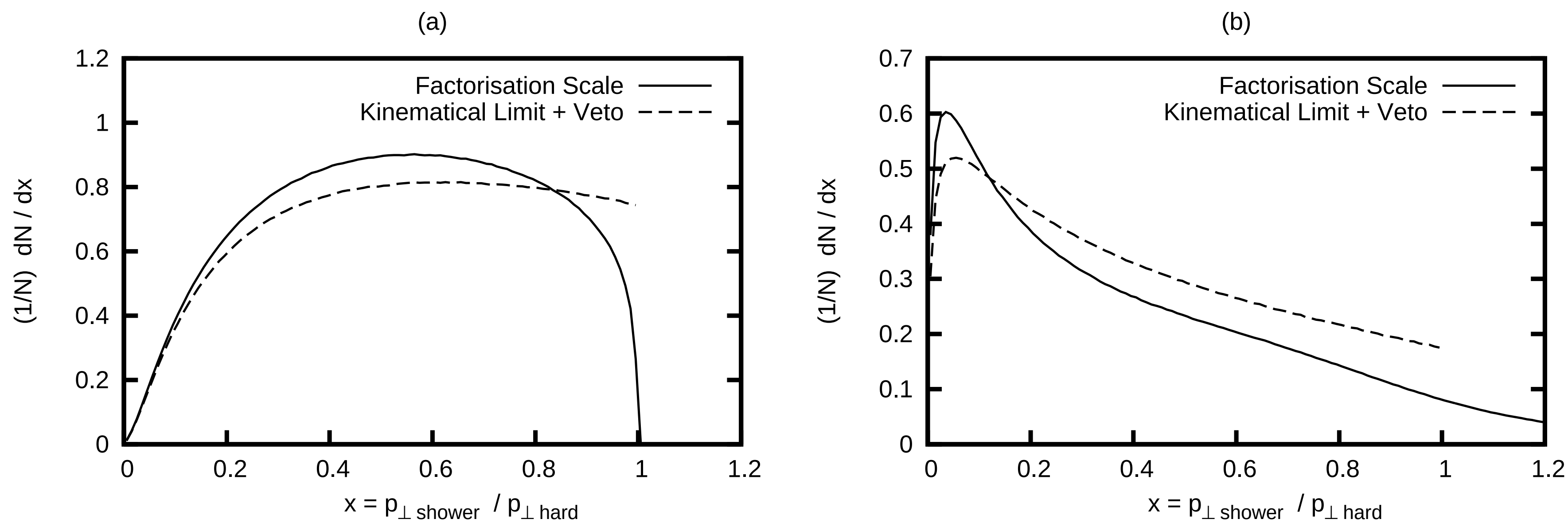
NLO matching natural extension

# Shower matching to MEs: POWHEG

Standard Les Houches interface (LHA, LHEF) specifies startup scale `SCALUP` for showers, so "trivial" to interface any external program, including POWHEG.
Problem: for ISR

$$p_\perp^2 = \mathrm{p}_{\perp\mathrm{evol}}^2 - \frac{\mathrm{p}_{\perp\mathrm{evol}}^4}{p_{\perp\mathrm{evol,max}}^2}$$

i.e. $p_\perp$ decreases for $\theta^* > 90^\circ$ but $\mathrm{p}_{\perp\mathrm{evol}}$ monotonously increasing.
Solution: run "power" shower but kill emissions above the hardest one, by POWHEG's definition.




Available for ISR-dominated, coming for QCD jets with FSR issues.

# Shower matching to MEs: VINCIA

Plug-in replacement to default PYTHIA showers; for now FSR only.

Developed by Walter Giele, David Kosower, Peter Skands, . . .

See talk by Juan Lopez-Villarejo on Wednesday 10 August!

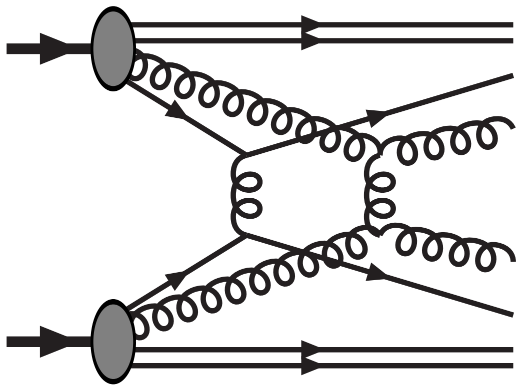
# Shower matching to MEs: the rest

As in PYTHIA 6, it is possible to use external ME generators with their own matching, e.g.

- ALPGEN (with MLM)
- MADGRAPH (with MLM)

# A second hard process

Multiparton interactions key aspect
of PYTHIA since $>$ 20 years.
Central to obtain agreement with data:
Tune A, Professor, Perugia, . . .



Before 8.1 no chance to select character of second interaction.
Now free choice of first process (including LHA/LHEF)
*and* second process combined from list:

- TwoJets (with TwoBJets as subsample)
- PhotonAndJet, TwoPhotons
- Charmonium, Bottomonium (colour octet framework)
- SingleGmZ, SingleW, GmZAndJet, WAndJet
- TopPair, SingleTop

Can be expanded among existing processes as need arises.

By default same phase space cuts as for “first” hard process
$\Longrightarrow$ second can be harder than first.
However, possible to set $\hat{m}$ and $\hat{p_\perp}$ range separately.

# Multiparton interactions

Regularise cross section with $p_{\perp 0}$ as free parameter

$$\frac{\mathrm{d}\hat{\sigma}}{\mathrm{d}p_{\perp}^2} \propto \frac{\alpha_s^2(p_{\perp}^2)}{p_{\perp}^4} \to \frac{\alpha_s^2(p_{\perp 0}^2 + p_{\perp}^2)}{(p_{\perp 0}^2 + p_{\perp}^2)^2}$$

with energy dependence

$$p_{\perp 0}(E_{\mathrm{CM}}) = p_{\perp 0}^{\mathrm{ref}} \times \left( \frac{E_{\mathrm{CM}}}{E_{\mathrm{CM}}^{\mathrm{ref}}} \right)^{\epsilon}$$

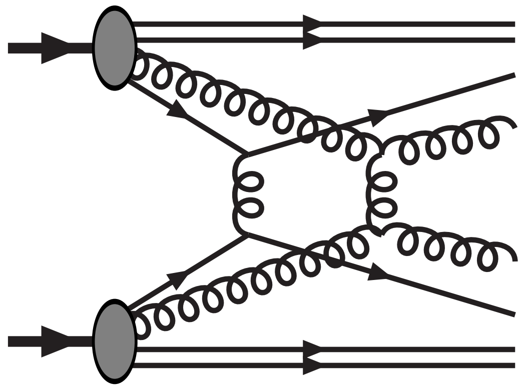
Matter profile in impact-parameter space
gives time-integrated overlap which determines level of activity:
simple Gaussian or more peaked variants

ISR and MPI compete for beam momentum $\longrightarrow$ PDF rescaling
$+$ flavour effects (valence, $\mathrm{q}\overline{\mathrm{q}}$ pair companions, . . . )
$+$ correlated primordial $k_{\perp}$ and colour in beam remnant

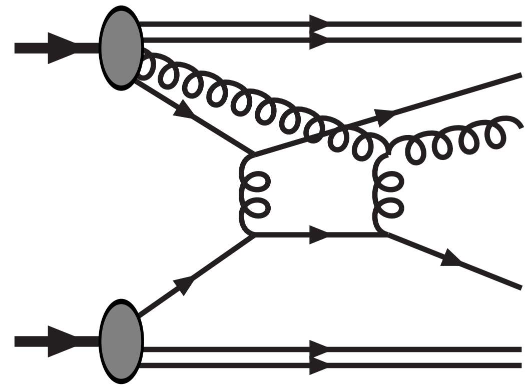
Many partons produced close in space–time $\Rightarrow$ colour rearrangement;
reduction of total string length $\Rightarrow$ steeper $\langle p_{\perp} \rangle (n_{\mathrm{ch}})$

# Rescattering

Often assume that MPI =



. . . but should also include



Same order in $\alpha_s$, $\sim$ same propagators, but

- one PDF weight less $\Rightarrow$ smaller $\sigma$
- one jet less $\Rightarrow$ QCD radiation background $2 \to 3$ larger than $2 \to 4$

$\Rightarrow$ will be tough to find direct evidence.

Rescattering grows with number of “previous” scatterings:

| | Tevatron | | LHC | |
|---|---|---|---|---|
| | Min Bias | QCD Jets | Min Bias | QCD Jets |
| Normal scattering | 2.81 | 5.09 | 5.19 | 12.19 |
| Single rescatterings | 0.41 | 1.32 | 1.03 | 4.10 |
| Double rescatterings | 0.01 | 0.04 | 0.03 | 0.15 |

R. Corke & TS, JHEP 01 (2010) 035

# An $x$-dependent proton size

Normally assume that PDFs factorize in longitudinal and transverse space:

$$f(x,r) = f(x)\,\rho(r)$$

In contradiction with

- intuitive picture of partons spreading out by cascade to lower $x$
- formally BFKL, Balitsky-JIMWLK, Colour Glass Condensate, ...
- Mueller's dipole cascade (Lund program: DIPSY; study by Avsar)
- Froissart-Martin $\sigma_{\mathrm{tot}} \propto \ln^2 s$ by Gribov theory related to $r_{\mathrm{p}} \propto \ln(1/x)$
- generalized parton distributions, ...

For now address only inelastic nondiffrative events,
$\sigma_{\mathrm{tot}} = \sigma_{\mathrm{el}} + \sigma_{\mathrm{SD}} + \sigma_{\mathrm{DD}} + \sigma_{\mathrm{ND}}$, with ansatz:

$$\rho(r,x) \propto \frac{1}{a^3(x)} \exp\left(-\frac{r^2}{a^2(x)}\right) \quad \text{with} \quad a(x) = a_0\left(1 + a_1 \ln\frac{1}{x}\right)$$

$a_1 \approx 0.15$ tuned to **rise** of $\sigma_{\mathrm{ND}}$ in Donnachie & Landshoff + Schuler & TS
$a_0$ tuned to **value** of $\sigma_{\mathrm{ND}}$, given PDF, $p_{\perp 0}$, ...

R. Corke & TS, JHEP 05 (2011) 009

Convolution of two incoming protons gives impact parameter shape

$$\tilde{\mathcal{O}}(b;x_1,x_2) = \frac{1}{\pi}\frac{1}{a^2(x_1)+a^2(x_2)}\,\exp\left(-\frac{b^2}{a^2(x_1)+a^2(x_2)}\right)$$

Define $\overline{n}(b)$ as average number of interactions for passage at $b$

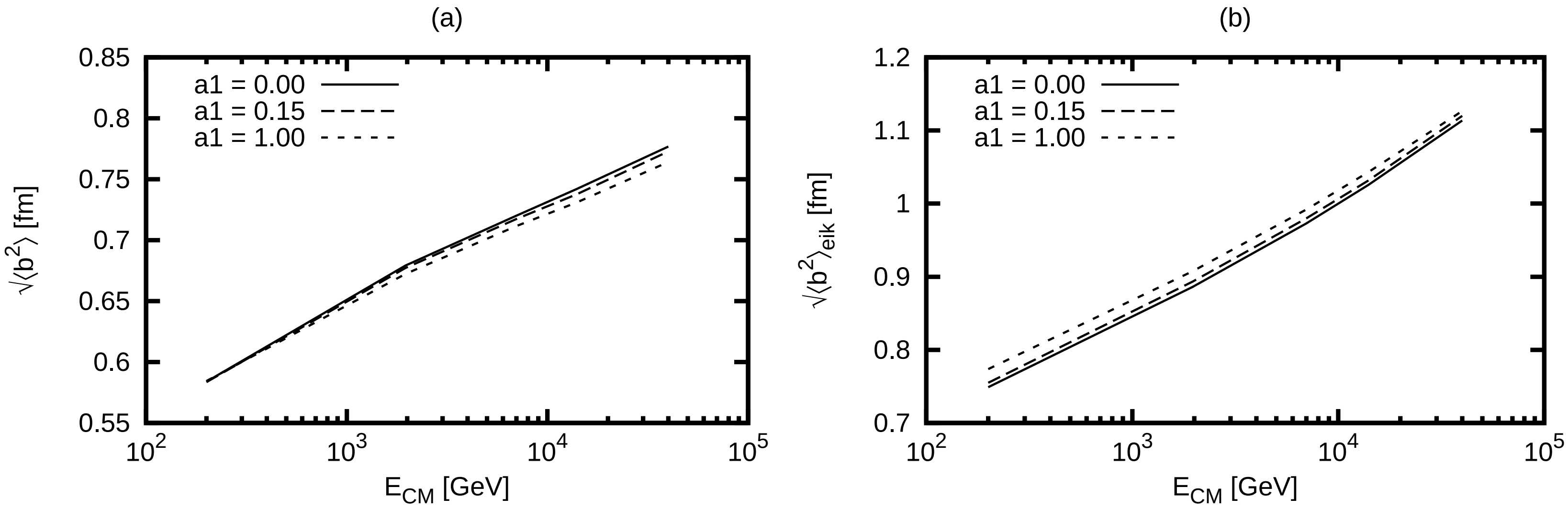
$$\overline{n}(b) = \sum_{i,j}\iiint \mathrm{d}x_1\,\mathrm{d}x_2\,\mathrm{d}p_\perp^2\, f_i(x_1,p_\perp^2)\, f_j(x_2,p_\perp^2)\, \left.\frac{\mathrm{d}\hat{\sigma}_{ij}}{\mathrm{d}p_\perp^2}\right|_{\mathrm{reg}} \tilde{\mathcal{O}}(b;x_1,x_2)$$

such that

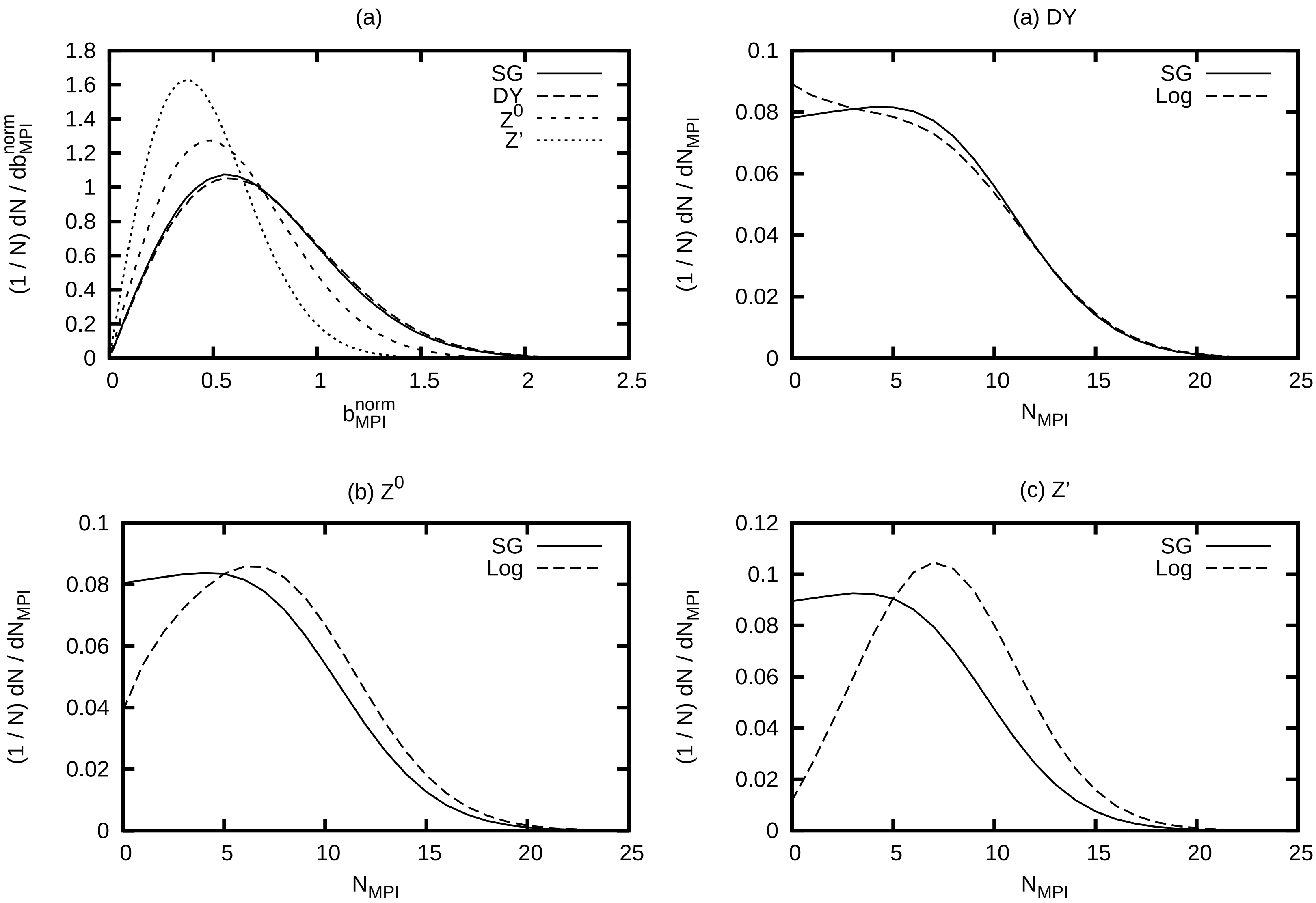
$$\sigma_{\mathrm{hard}} = \int \overline{n}(b)\,\mathrm{d}^2b$$

$$\sigma_{\mathrm{ND}} = \int P_{\mathrm{int}}\,\mathrm{d}^2b = \int \left(1-e^{-\overline{n}(b)}\right)\mathrm{d}^2b$$

which gives a $\sqrt{\langle b^2\rangle}$ for $\sigma_{\mathrm{hard}}$ and $\sigma_{\mathrm{ND}}$:

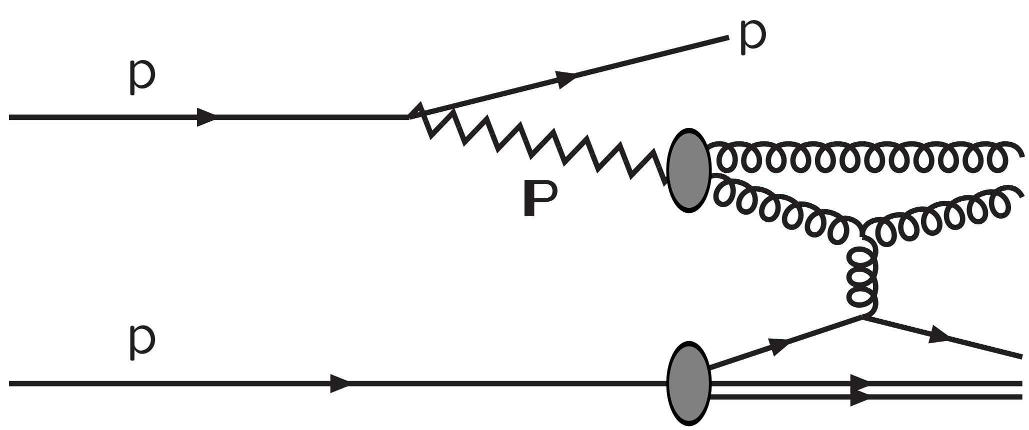


Consequence: collisions at large $x$ will have to happen at small $b$,
and hence further large-to-medium-$x$ MPIs are enhanced,
while low-$x$ partons are so spread out that it plays less role.



# Diffraction

Ingelman-Schlein: Pomeron as hadron with partonic content
Diffractive event = (Pomeron flux) $\times$ ($\mathbb{P}$p collision)




Used e.g. in
POMPYT
POMWIG
PHOJET

1) $\sigma_{\mathrm{SD}}$ and $\sigma_{\mathrm{DD}}$ taken from existing parametrization or set by user.

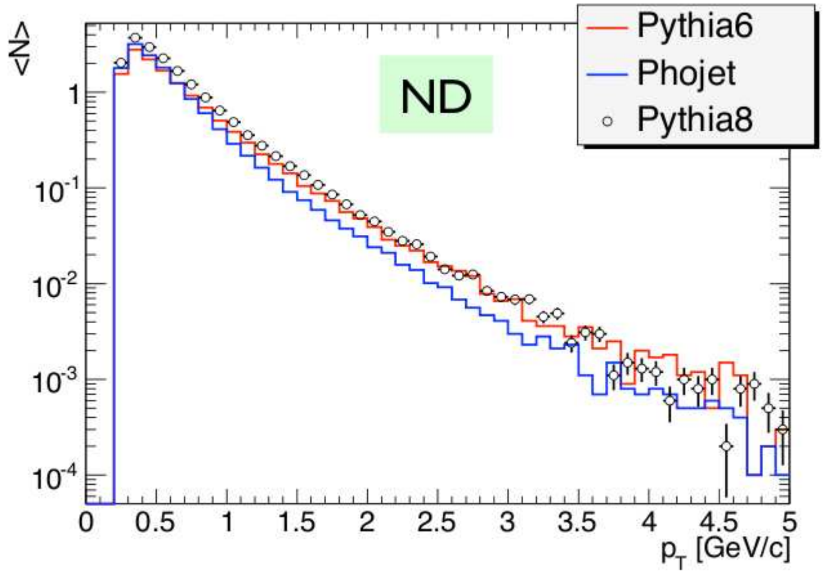
2) Shape of Pomeron distribution inside a proton, $f_{\mathbb{P}/\mathrm{p}}(x_{\mathbb{P}}, t)$ gives diffractive mass spectrum and scattering $p_{\perp}$ of proton.

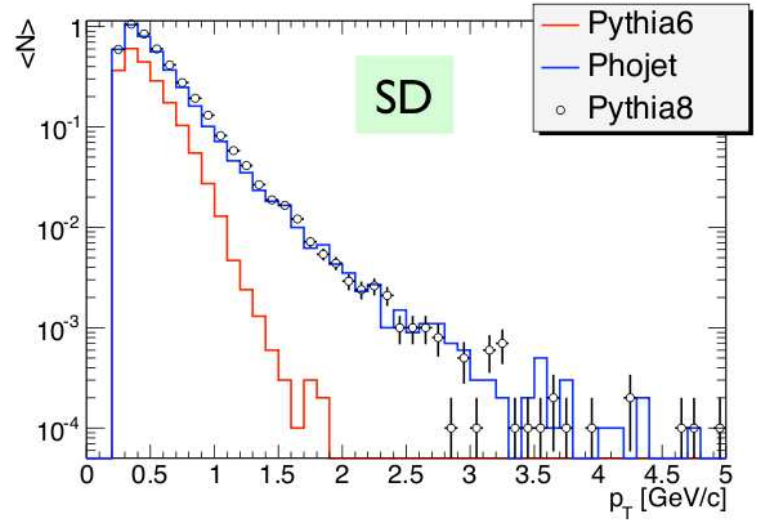
3) At low masses retain old framework, with longitudinal string(s). Above 10 GeV begin smooth transition to $\mathbb{P}$p handled with full pp machinery: multiple interactions, parton showers, beam remnants, . . . .

4) Choice between 5 Pomeron PDFs. Free parameter $\sigma_{\mathbb{P}\mathrm{p}}$ needed to fix $\langle n_{\mathrm{interactions}} \rangle = \sigma_{\mathrm{jet}}/\sigma_{\mathbb{P}\mathrm{p}}$.

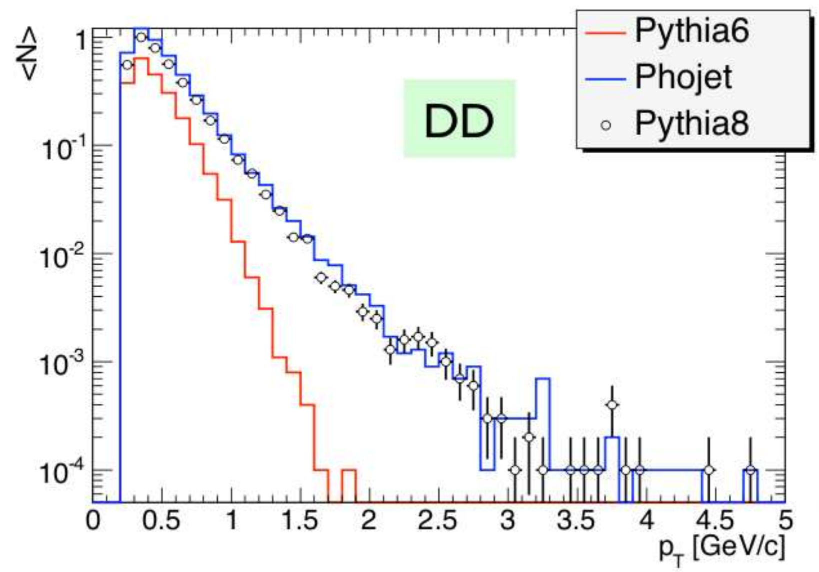
5) Framework needs testing and tuning, e.g. of $\sigma_{\mathbb{P}\mathrm{p}}$.

# $p_T$ Distributions ($\sqrt{s}$=0.9 TeV)




- Softer $p_T$ spectrum in **Pythia6** due to lack of high mass diffraction
- **Pythia8** and **Phojet** agree quite well

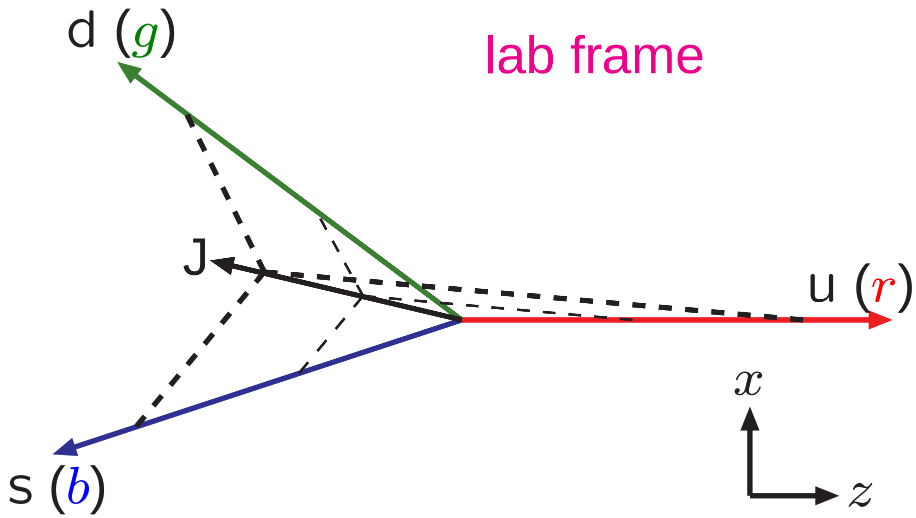


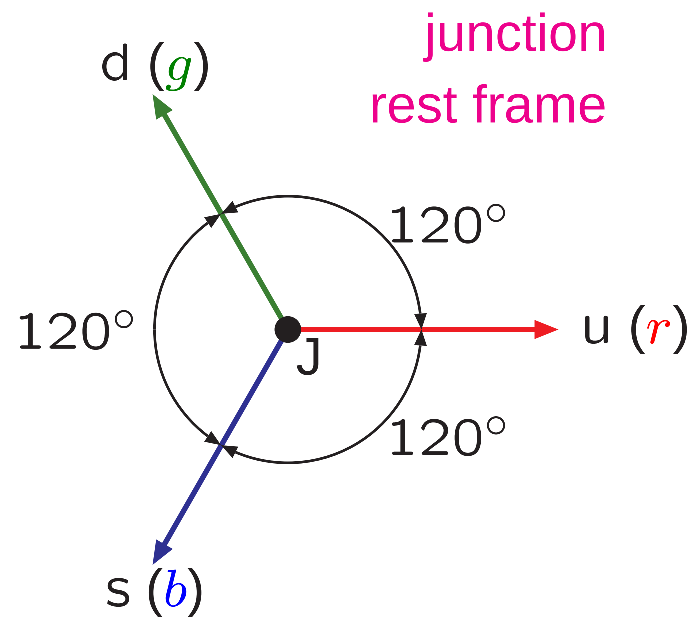


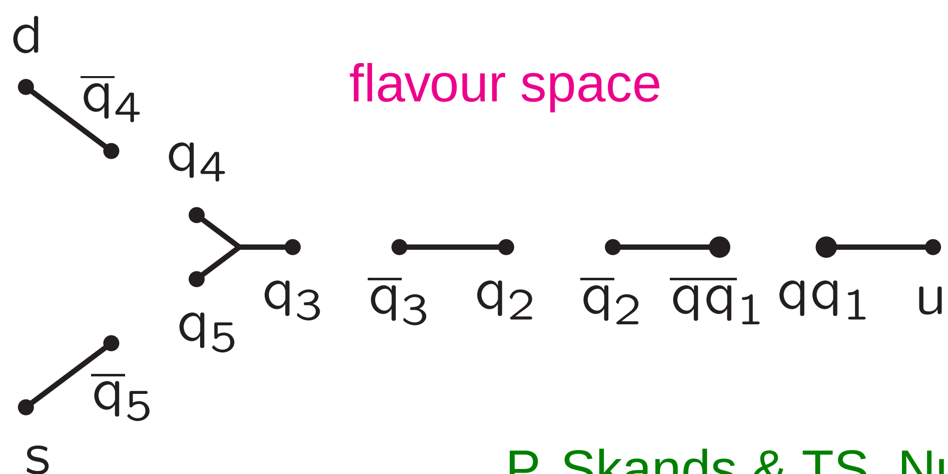

Beate Heinemann, MB/UE Working Group (also Sparsh Navin)

# BSM Physics 1: $R$-parity violation

Encountered in $R$-parity violating SUSY decays $\tilde{\chi}_1^0 \to$ uds,
or when 2 valence quarks kicked out of proton beam




More complicated
(but ≈solved) with
gluon emission and
massive quarks

P. Skands & TS, Nucl. Phys. B659 (2003) 243

# BSM Physics 2: $R$-hadrons

What if coloured (SUSY) particle like $\tilde{\mathrm{g}}$ or $\tilde{\mathrm{t}}_1$ is long-lived?

⋆ Formation of $R$-hadrons

| | | |
|---|---|---|
| $\tilde{\mathrm{g}}\mathrm{q}\overline{\mathrm{q}}$ | $\tilde{\mathrm{t}}_1\overline{\mathrm{q}}$ | "mesons" |
| $\tilde{\mathrm{g}}\mathrm{qqq}$ | $\tilde{\mathrm{t}}_1\mathrm{qq}$ | "baryons" |
| $\tilde{\mathrm{g}}\mathrm{g}$ | | "glueballs" |

⋆ Conversion between $R$-hadrons

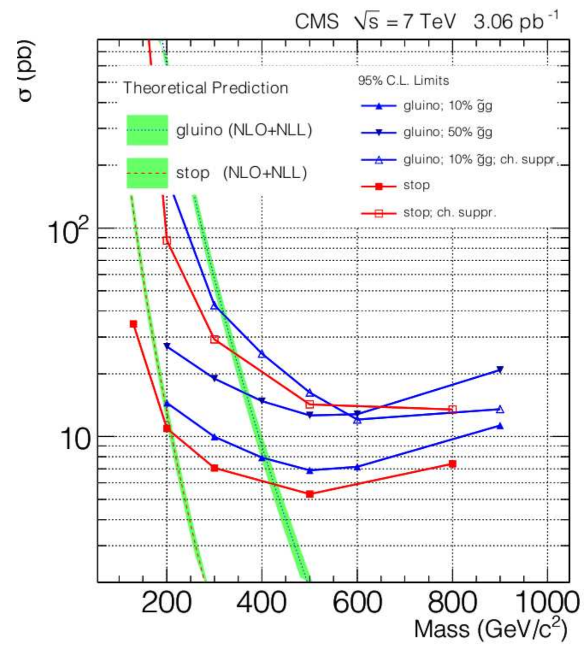
by "low-energy" interactions with matter:

$\tilde{\mathrm{g}}\mathrm{u}\overline{\mathrm{d}} + \mathrm{p} \to \tilde{\mathrm{g}}\mathrm{uud} + \pi^+$ irreversible

⋆ Displaced vertices if finite lifetime, or else

⋆ punch-through: $\sigma \approx \sigma_{\mathrm{had}}$ but

$\Delta E \lesssim 1$ GeV $\ll E_{\mathrm{kin},R}$




A.C. Kraan, Eur. Phys. J. C37 (2004) 91;

M. Fairbairn et al., Phys. Rep. 438 (2007) 1

CMS, arXiv:1101.1645

Partly event generation, partly detector simulation.

Public add-on in PYTHIA 6, now integrated part of PYTHIA 8.

Can also be applied to non-SUSY long-lived "hadrons".

# BSM Physics 3: Hidden Valley (Secluded Sector)

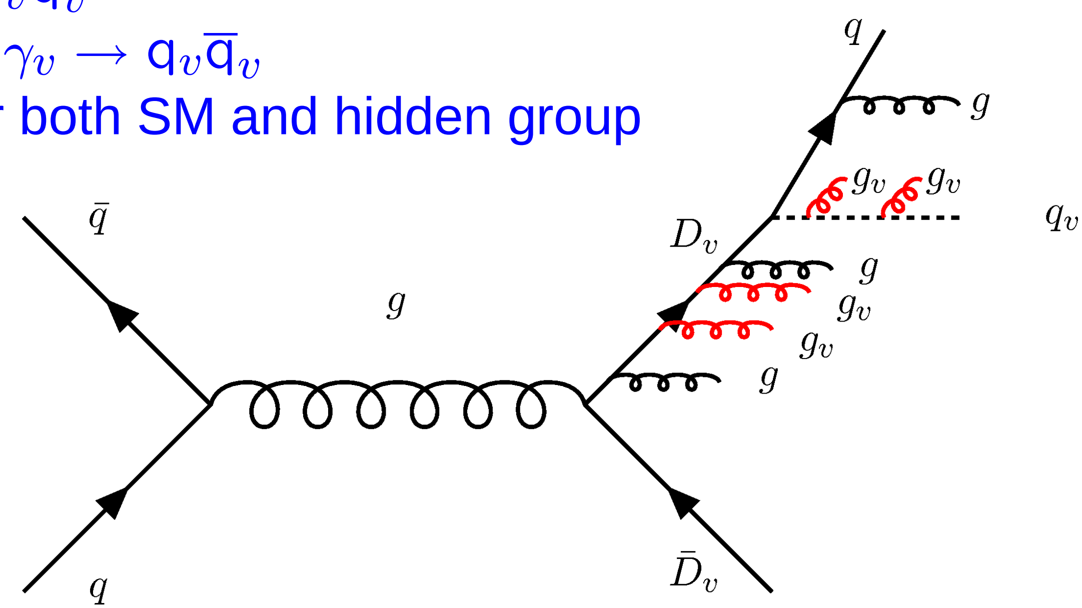
What if new gauge groups at low energy scales, hidden by potential barrier or weak couplings? (M. Strassler & K. Zurek, . . . )

Complete framework implemented in PYTHIA:

⋆ New gauge group either Abelian $U(1)$ or non-Abelian $SU(N)$

⋆ 3 alternative production mechanisms
1) massive $\mathrm{Z}'$: $\mathrm{q}\overline{\mathrm{q}} \to \mathrm{Z}' \to \mathrm{q}_v\overline{\mathrm{q}}_v$
2) kinetic mixing: $\mathrm{q}\overline{\mathrm{q}} \to \gamma \to \gamma_v \to \mathrm{q}_v\overline{\mathrm{q}}_v$
3) massive $F_v$ charged under both SM and hidden group

⋆ Interleaved shower in QCD, QED and HV sectors: add $\mathrm{q}_v \to \mathrm{q}_v\gamma_v$ (and $F_v$) or $\mathrm{q}_v \to \mathrm{q}_v\mathrm{g}_v$, $\mathrm{g}_v \to \mathrm{g}_v\mathrm{g}_v$, which gives recoil effects also in visible sector



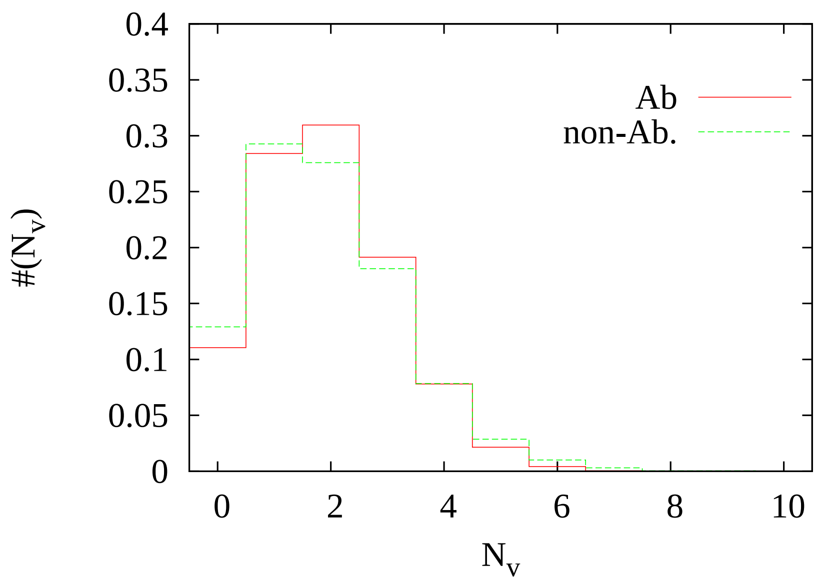

L. Carloni & TS, JHEP 09 (2010) 105;
L. Carloni, J. Rathsman & TS, JHEP 04 (2011) 091

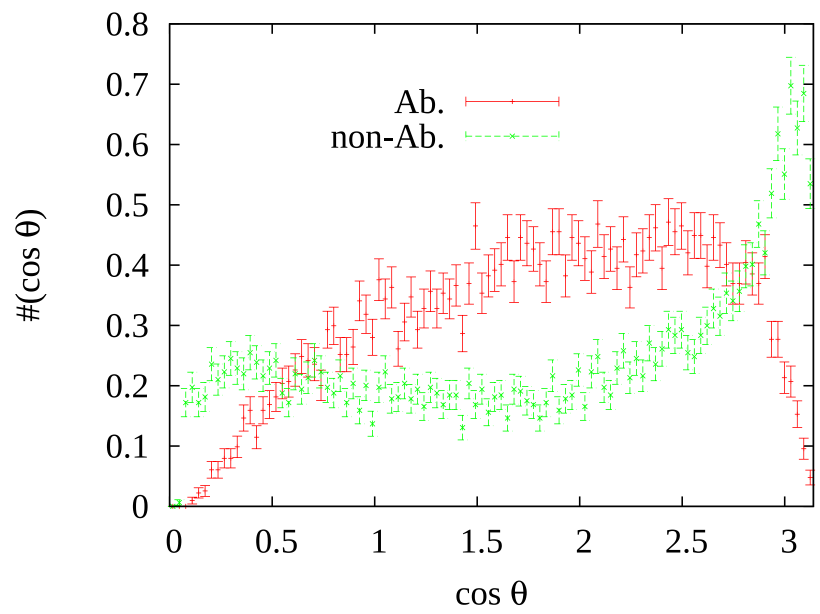
⋆ Hidden Valley particles may remain invisible, or . . .

⋆ Broken $U(1)$: $\gamma_v$ acquire mass, radiated $\gamma_v$s decay back
$\gamma_v \to \gamma \to \mathrm{f\bar{f}}$ with BRs as photon (⇒ lepton pairs!)

⋆ $SU(N)$: hadronization in hidden sector, with full string fragmentation,
permitting up to 8 different $\mathrm{q}_v$ flavours and 64 $\mathrm{q}_v\bar{\mathrm{q}}_v$ mesons,
but for now assumed degenerate in mass, so only distinguish
– off-diagonal, flavour-charged, stable & invisible
– diagonal, can decay back $\mathrm{q}_v\bar{\mathrm{q}}_v \to \mathrm{f\bar{f}}$

Even when tuned to same average activity, hope to separate $U(1)$ and $SU(N)$:

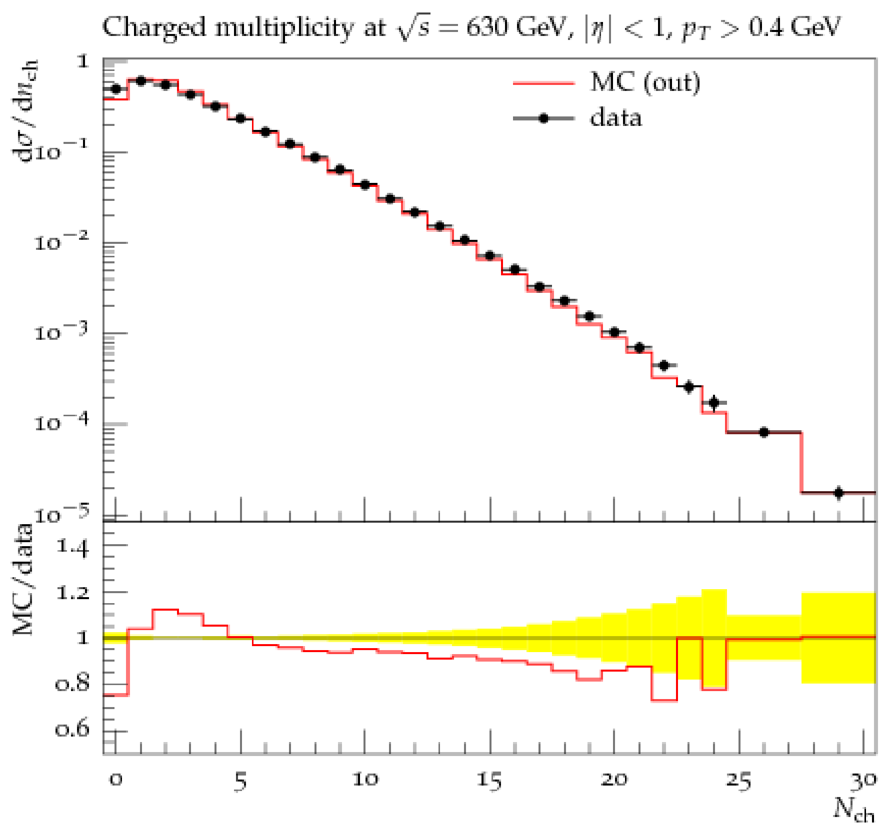


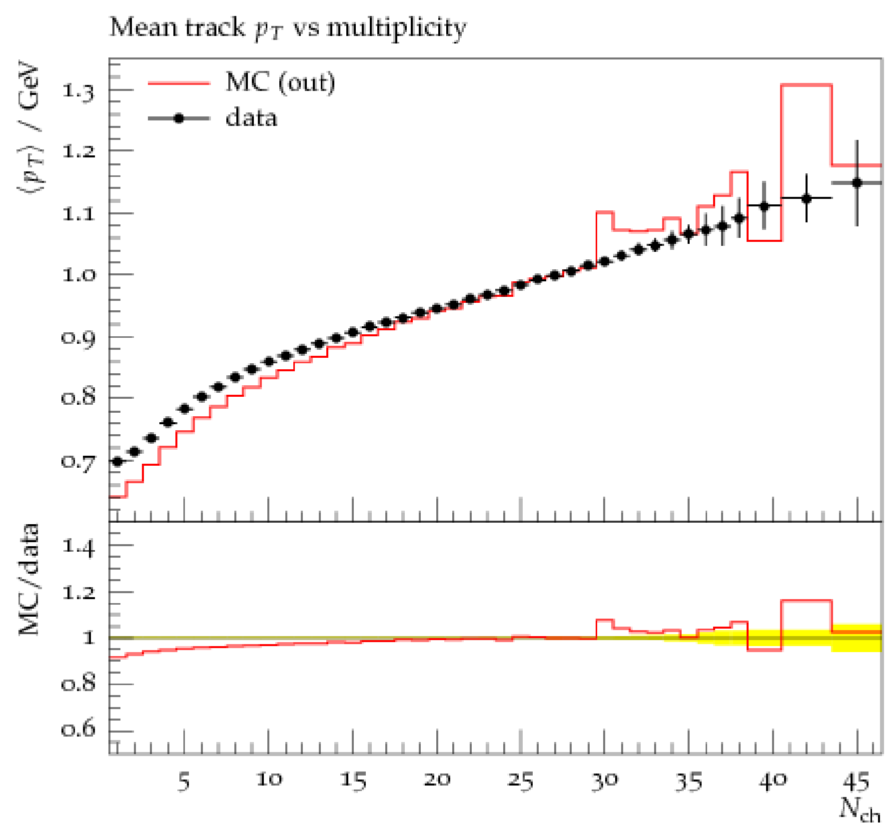



# Tuning prospects

Tuning to $e^+e^-$ closely related to $p_\perp$-ordered PYTHIA 6.4; Rivet+Professor (H. Hoeth) ⇒ FSR & hadronization OK (?)

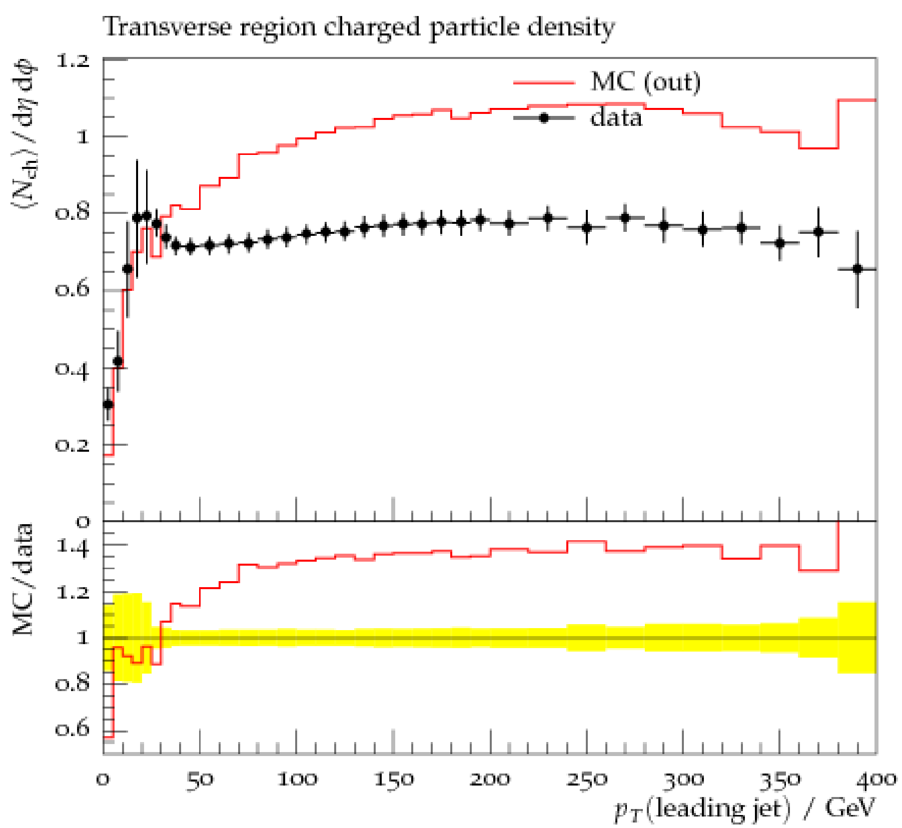
First tuning to MB data by P. Skands:

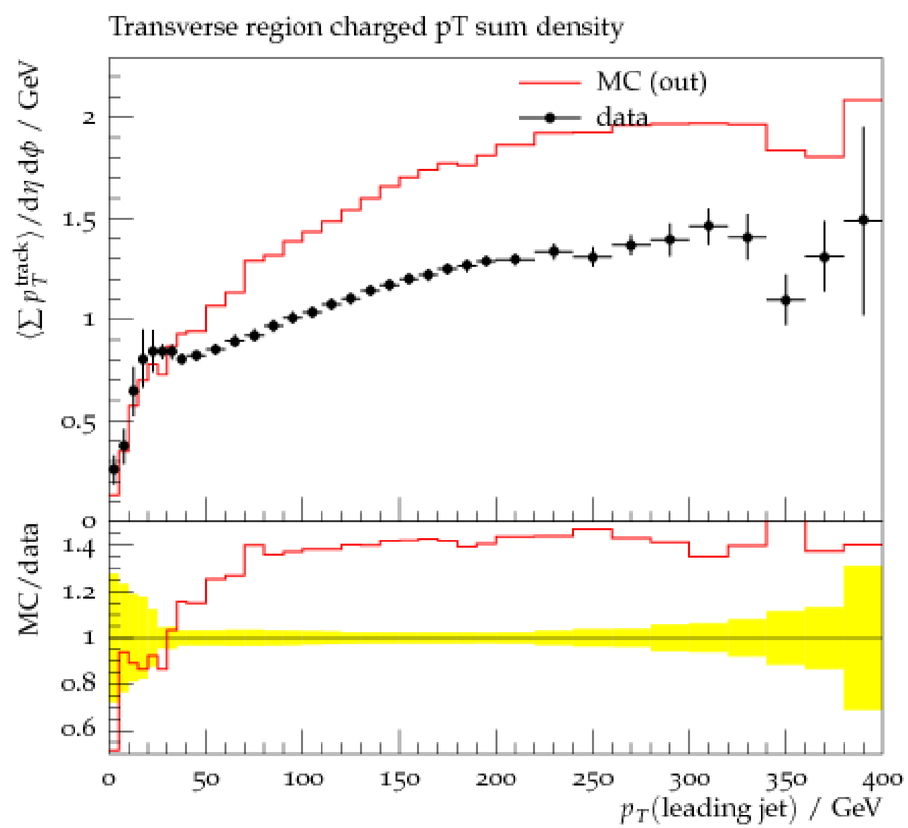



⇒ MPI & colour reconnection OK (?)

But Rivet+Professor (H. Hoeth) shows it fails miserably for UE (Rick Field's transverse flow as function of jet $p_\perp$):

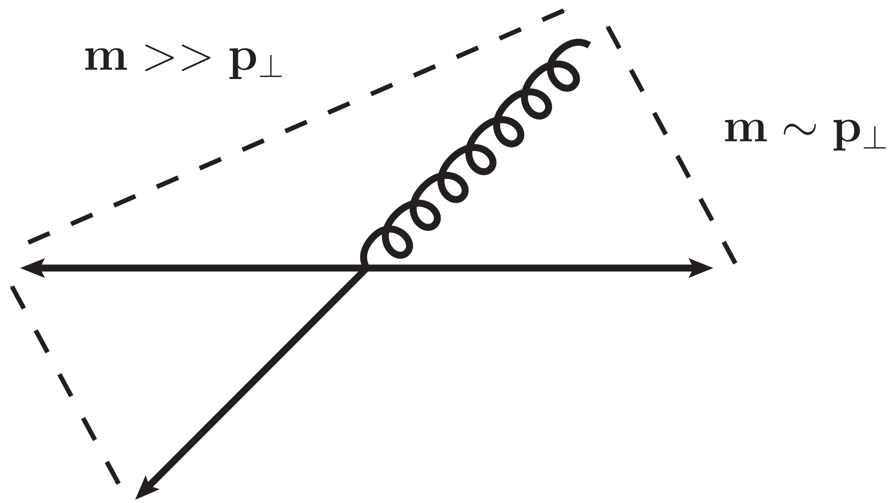


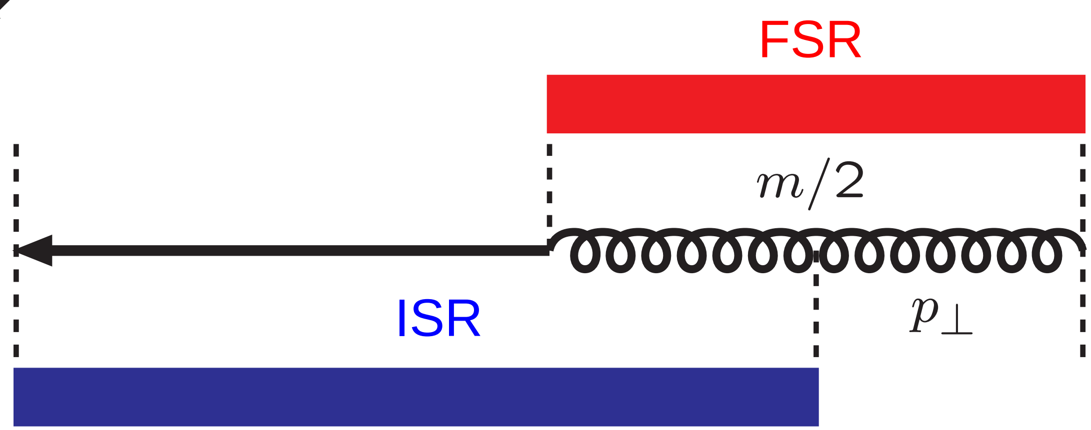

No universal tune MB + UE!

**Where did we go wrong?**

Final-state parton may have colour partner in the initial state.
How to subdivide FSR and ISR in this kind of dipole?
Large mass $\longrightarrow$ large rapidity range for emission:

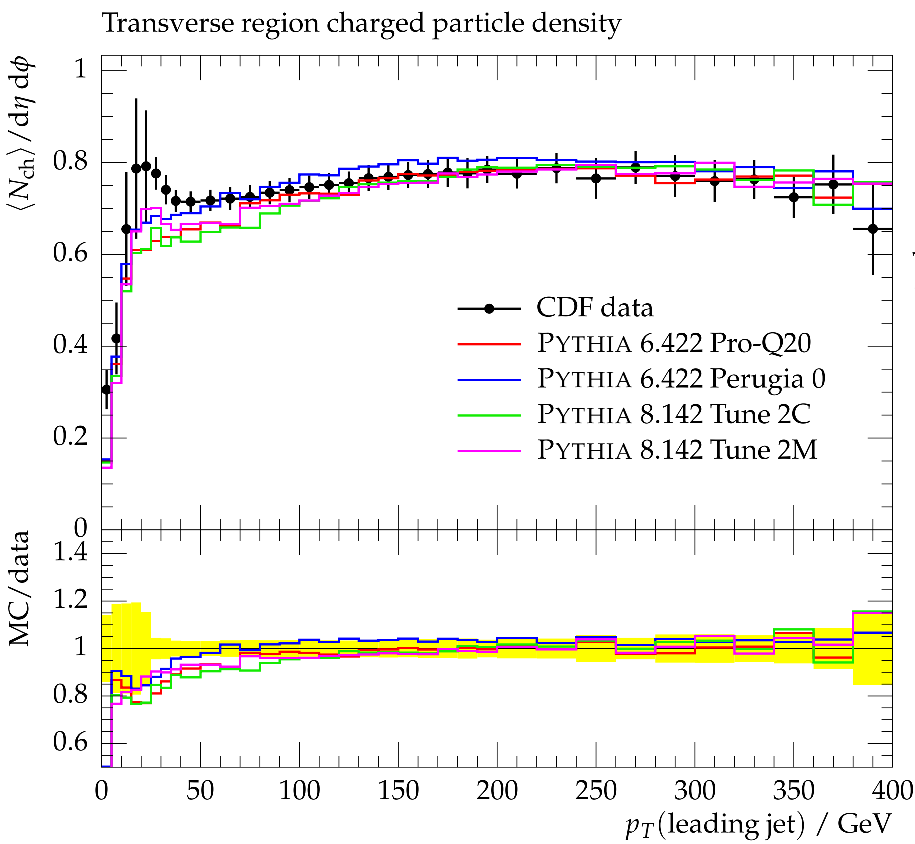


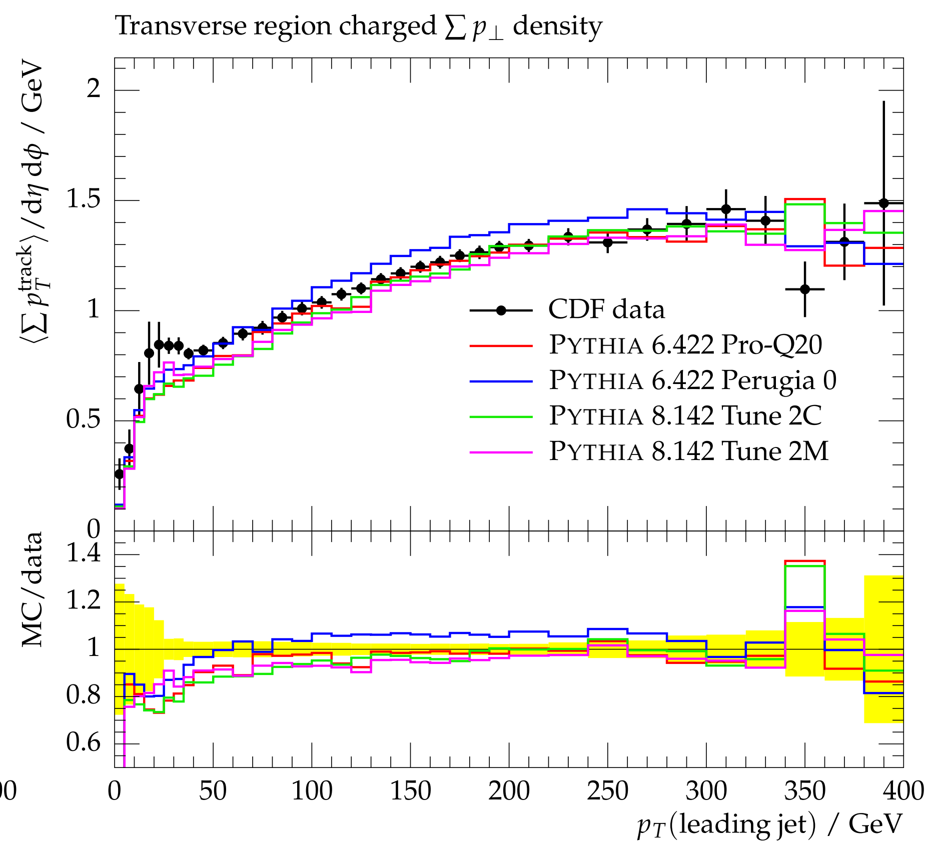

In dipole rest frame
(think rapidity space)




Solution: suppress final-state radiation in double-counted region

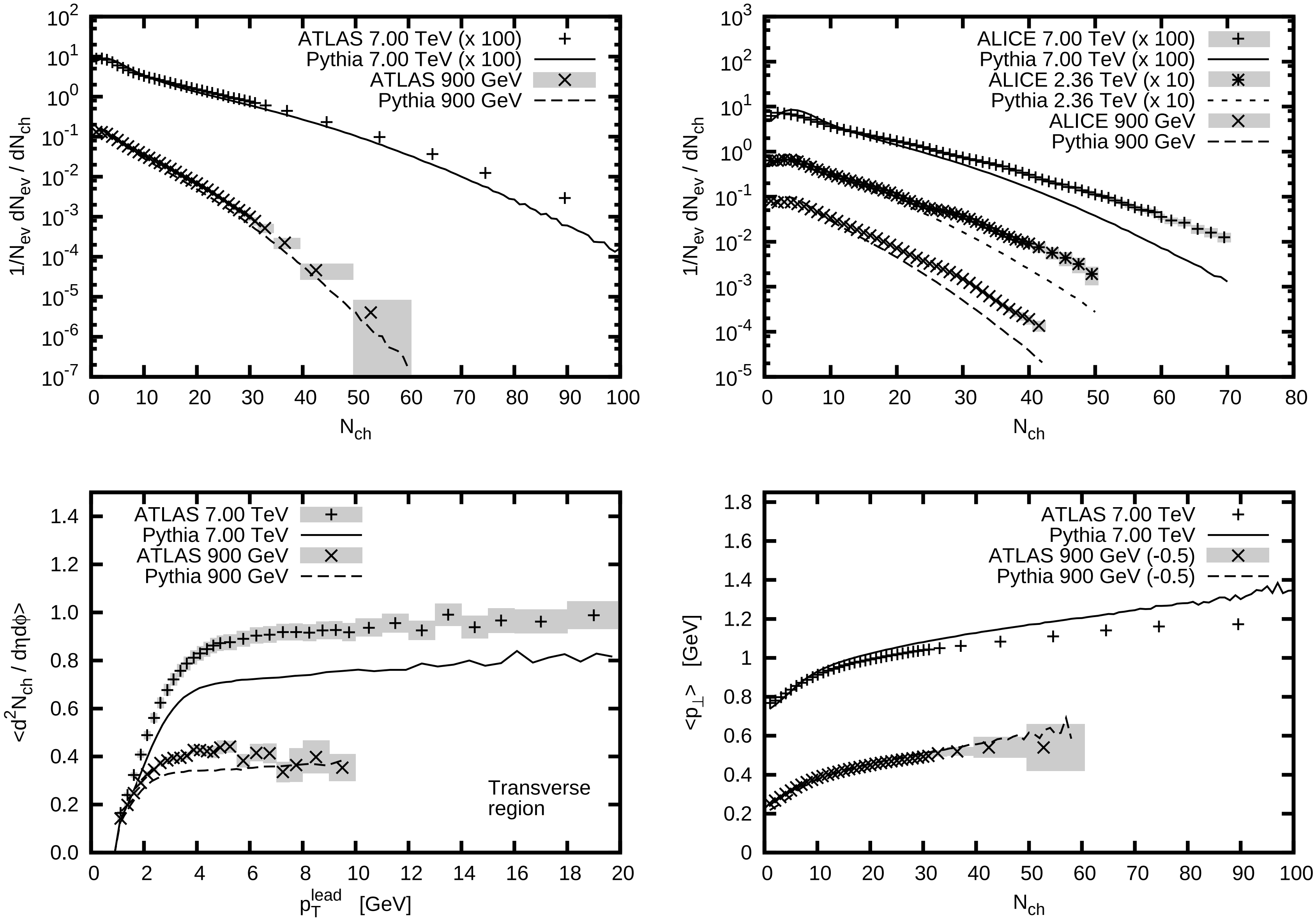
Is a simultaneous MB/UE Tevatron tune now possible?
Tunes 2C and 2M done “by hand” (= using Rivet, but not Professor), using the CTEQ6L1 and MRST LO** PDF sets, respectively, to MB data ($n_{\text{ch}}$, $\langle p_\perp \rangle(n_{\text{ch}})$, . . . .
Compare against Pro-Q20 and Perugia 0 (PYTHIA 6)



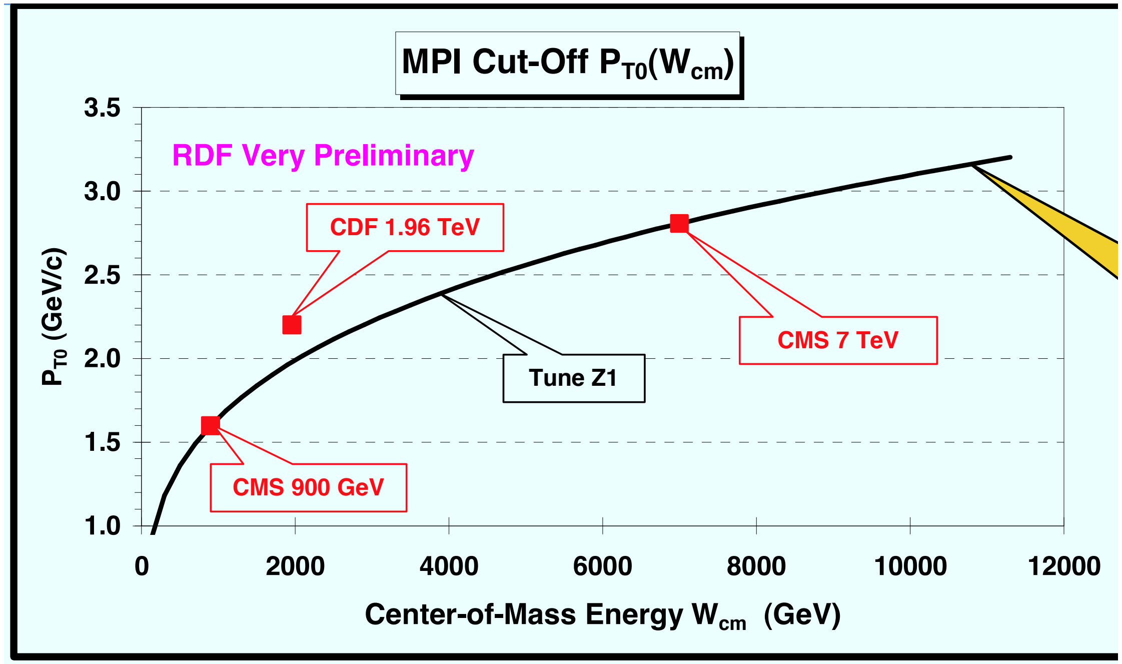

Now generally comparably good description to old tunes at the Tevatron.

# Tune 2C applied to LHC data does not do so good:



# Tension between Tevatron and LHC data?

Rick Field: if $p_{\perp 0}(E_{\mathrm{CM}}) \propto E_{\mathrm{CM}}^{\epsilon}$ tuned to LHC data,
then gives too much UE activity at Tevatron
($\Rightarrow$ need higher $p_{\perp 0}$ to compensate)




Pick some key LHC data sets, use Tune 2C as starting point:

- slightly dampen diffractive cross section (ATLAS)
- only vary MPI and colour reconnection parameters

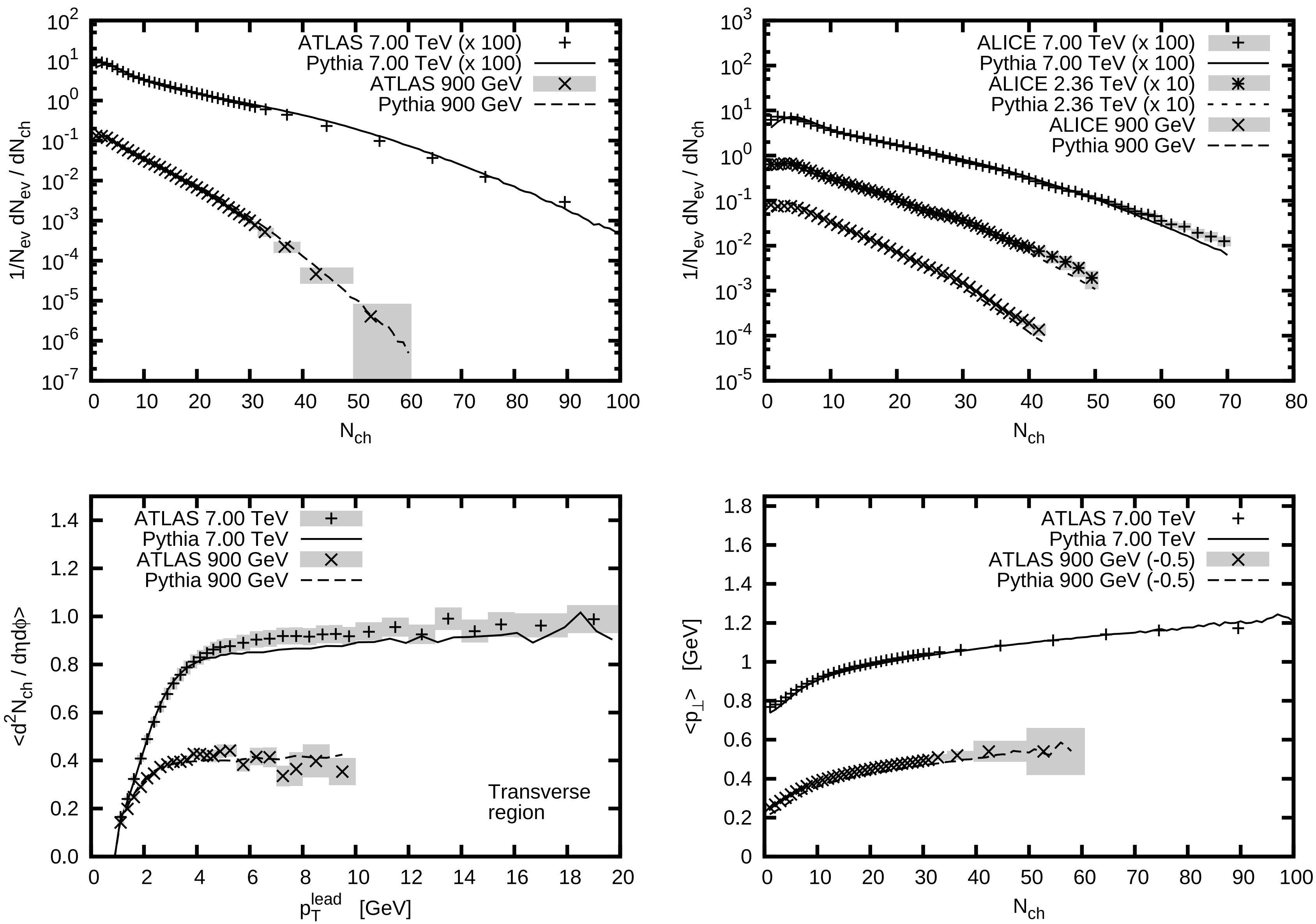
# New tunes

. . . while waiting for more LHC data conveniently implemented in Rivet.

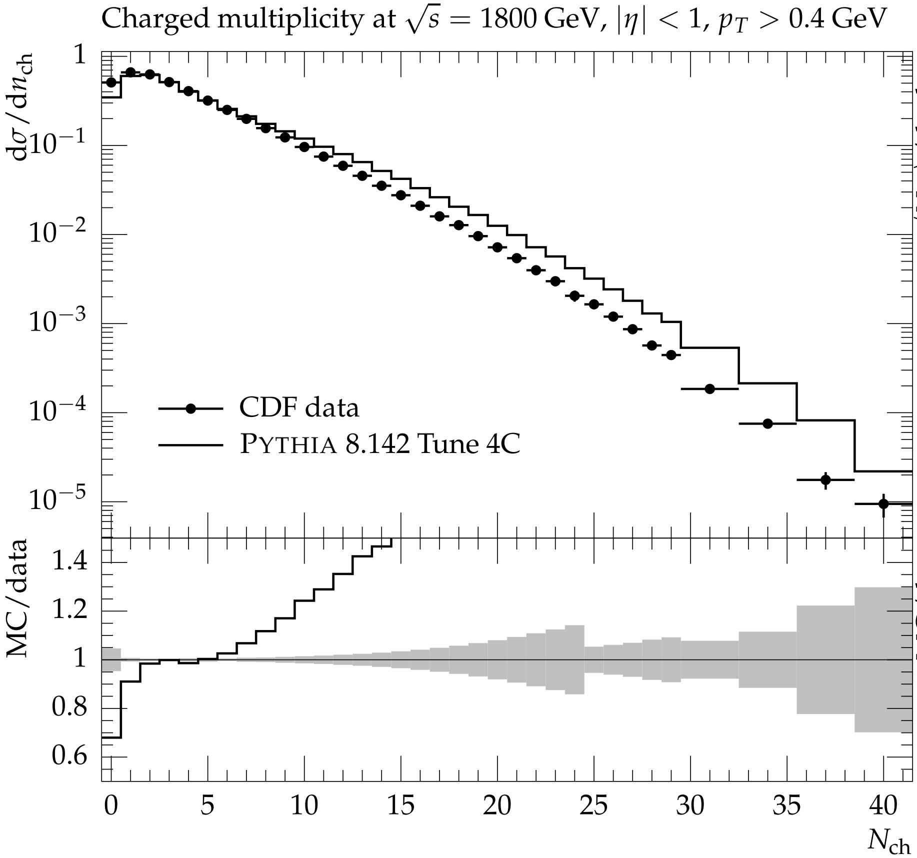
| **Parameter** | **2C** | **2M** | **4C** | **4Cx** |
|---|---|---|---|---|
| `SigmaProcess:alphaSvalue` | 0.135 | 0.1265 | 0.135 | 0.135 |
| `SpaceShower:rapidityOrder` | on | on | on | on |
| `SpaceShower:alphaSvalue` | 0.137 | 0.130 | 0.137 | 0.137 |
| `SpaceShower:pT0Ref` | 2.0 | 2.0 | 2.0 | 2.0 |
| `MultipleInteractions:alphaSvalue` | 0.135 | 0.127 | 0.135 | 0.135 |
| `MultipleInteractions:pT0Ref` | 2.320 | 2.455 | 2.085 | 2.15 |
| `MultipleInteractions:ecmPow` | 0.21 | 0.26 | 0.19 | 0.19 |
| `MultipleInteractions:bProfile` | 3 | 3 | 3 | 4 |
| `MultipleInteractions:expPow` | 1.60 | 1.15 | 2.00 | N/A |
| `MultipleInteractions:a1` | N/A | N/A | N/A | 0.15 |
| `BeamRemnants:reconnectRange` | 3.0 | 3.0 | 1.5 | 1.5 |
| `SigmaDiffractive:dampen` | off | off | on | on |
| `SigmaDiffractive:maxXB` | N/A | N/A | 65 | 65 |
| `SigmaDiffractive:maxAX` | N/A | N/A | 65 | 65 |
| `SigmaDiffractive:maxXX` | N/A | N/A | 65 | 65 |

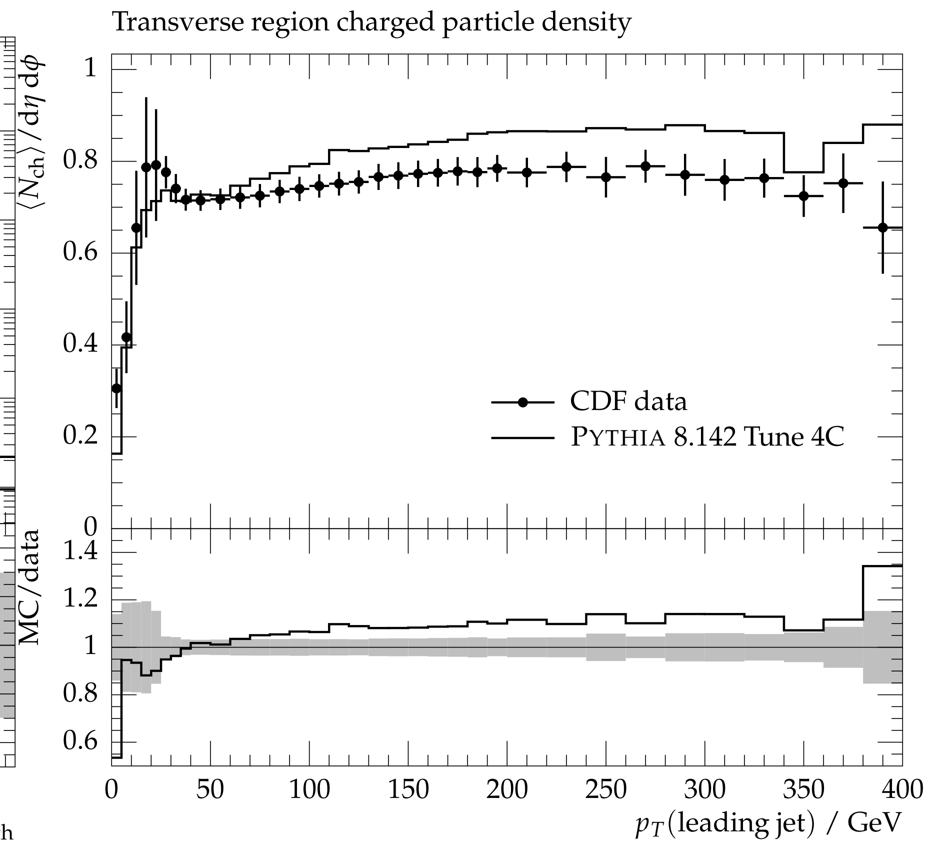
R. Corke & TS, JHEP 03 (2011) 032, JHEP 05 (2011) 009

# Tune 4C now describes LHC data reasonably well:



. . . but at the expense of Tevatron agreement:



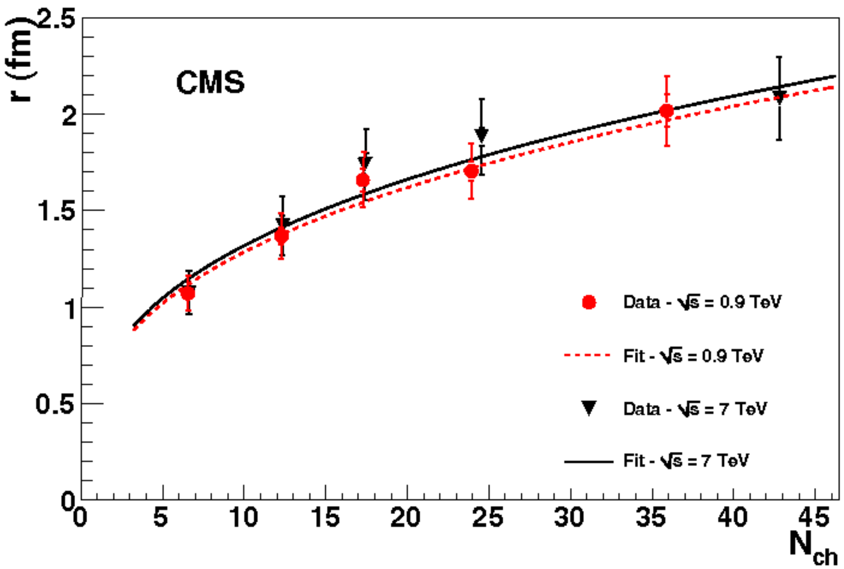



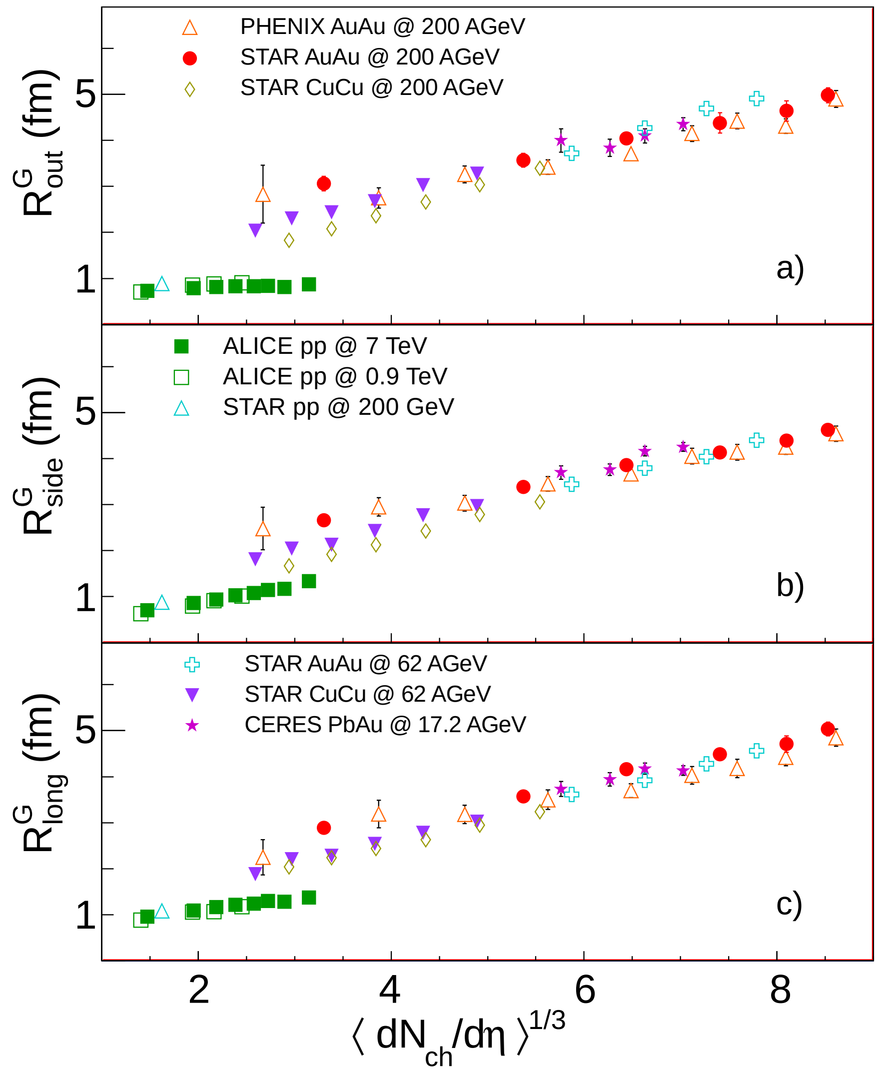

Future:

- better understanding of data?
- official/validated inclusion in Rivet?
- combined tune Tevatron + LHC?

# Further challenges from hadron collider data

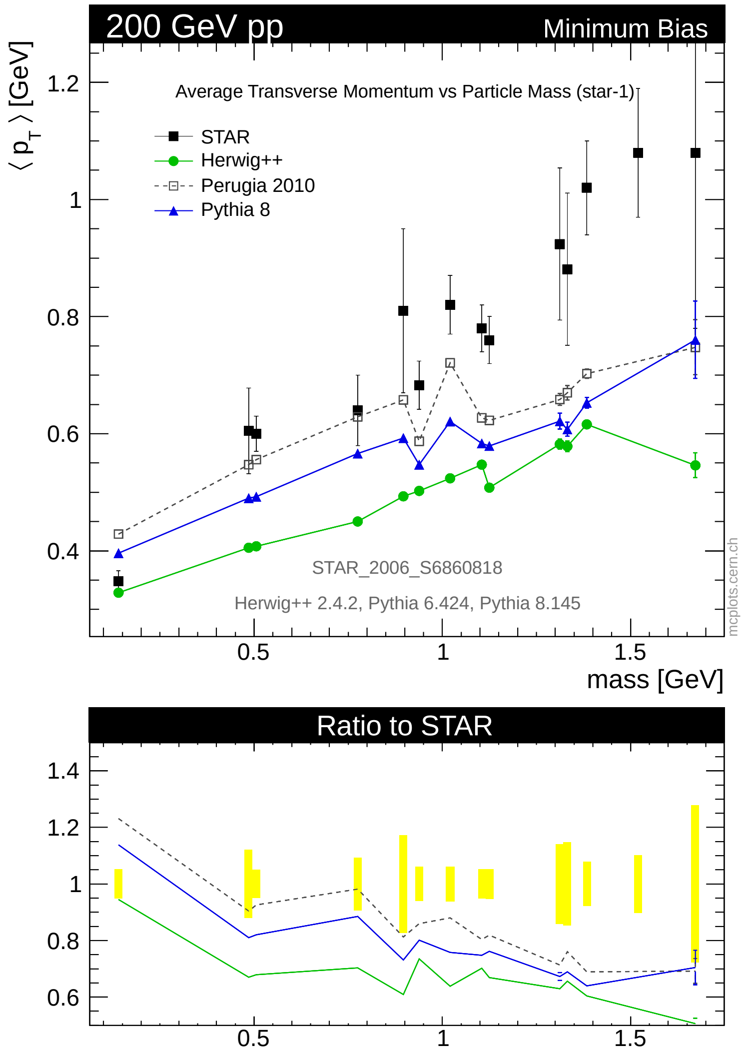
Bose-Einstein $r(N_{\text{ch}}) \propto N_{\text{ch}}^{1/3}$
cannot be accommodated
in PYTHIA effective description
that worked at LEP

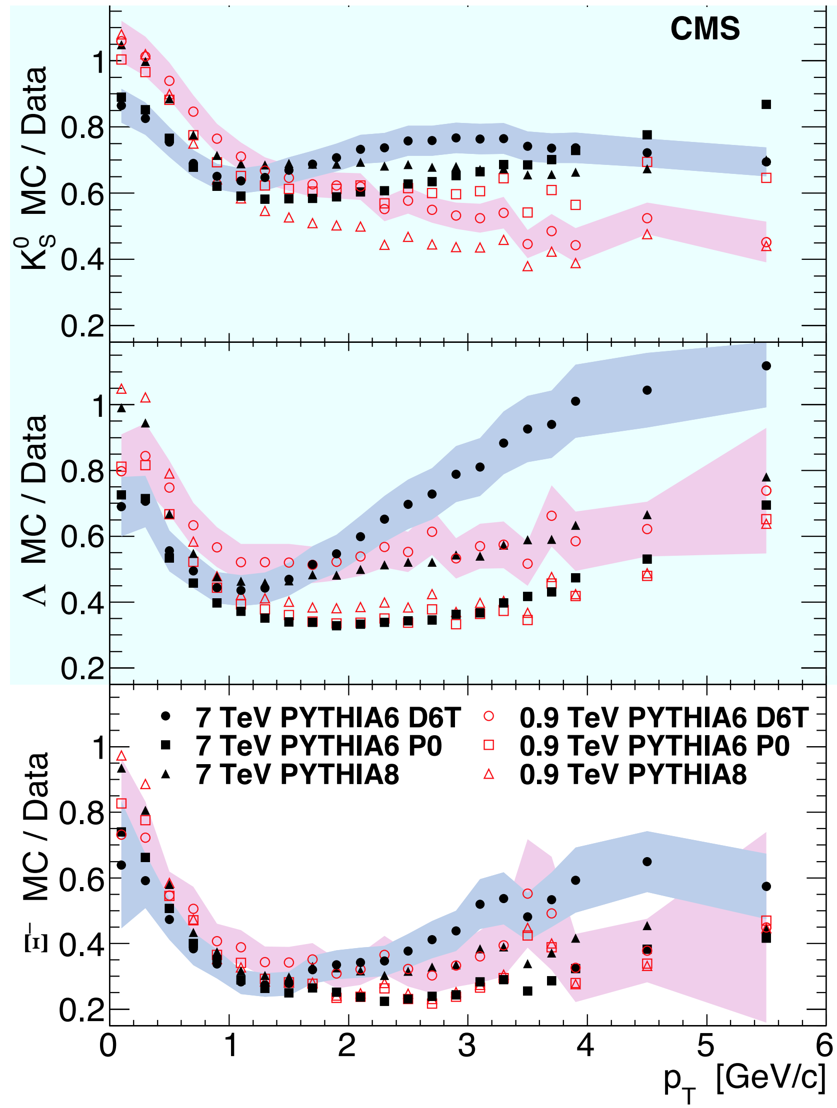



Multiple overlapping fragmenting strings $\Rightarrow$ dense hadron gas!

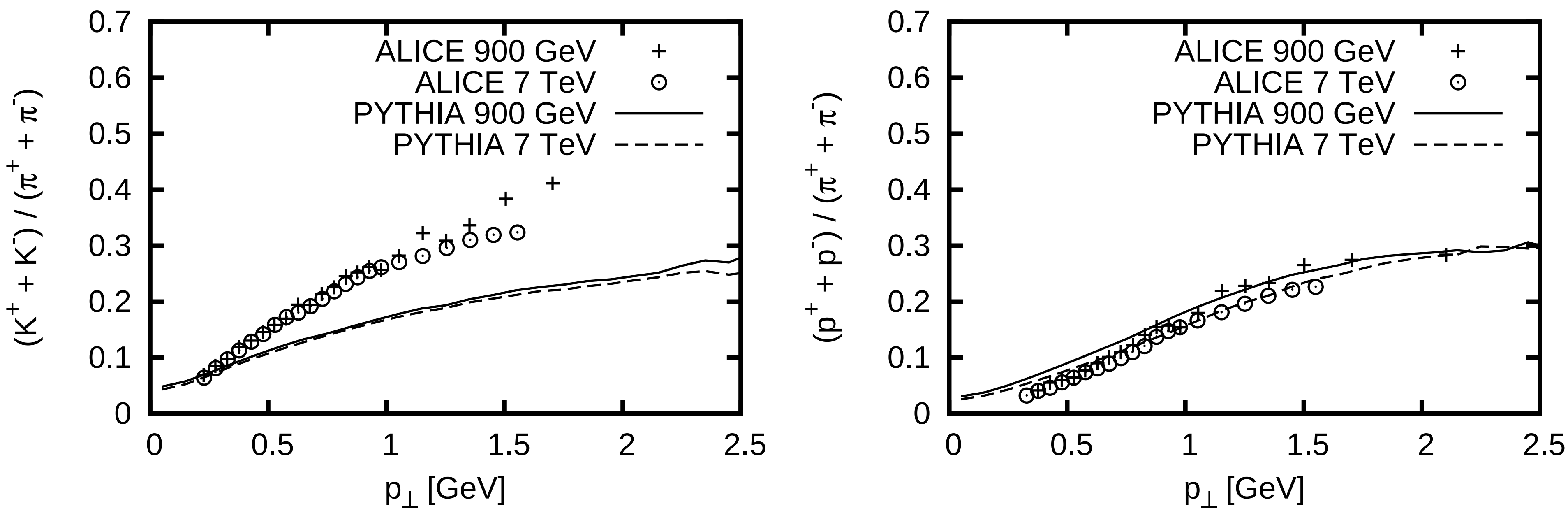
# Need more $p_\perp$ for K, p, Λ, . . . , relative to $\pi^\pm$:



Latest $\pi$/K/p $p_\perp$ spectra from LHC

- ALICE 900 GeV - inelastic events, $|y| < 0.5$, arXiv:1101.4110v3
- ALICE 7 TeV - read from public talks (assume same acceptance as 900 GeV)
- Compare against PYTHIA 8




Pilot study (R. Corke):

model for rescattering as afterburner.

No need to change particle composition (significantly),

only (partial) collective flow?

# Final-state hadron scattering

Final-state hadron scattering

- Low-energy scatterings in CM frame and boost back
- Higher masses take bigger kick, c.f. collective flow in heavy ion
- Ideally assign production vertices to outgoing hadrons and follow path

Simple model based on distance in $y - \phi$ space

- High-$p_\perp$ hadrons in jets formed at later times and less likely to scatter
- Scattering probability

$$P_{ij} = k \frac{\sigma_{ij}^{\text{el}}(s)}{\sigma_{\text{max}}^{\text{el}}} \max\left(0,\ 1 - \frac{\Delta R_{ij}^2}{R_{\text{max}}^2}\right)$$

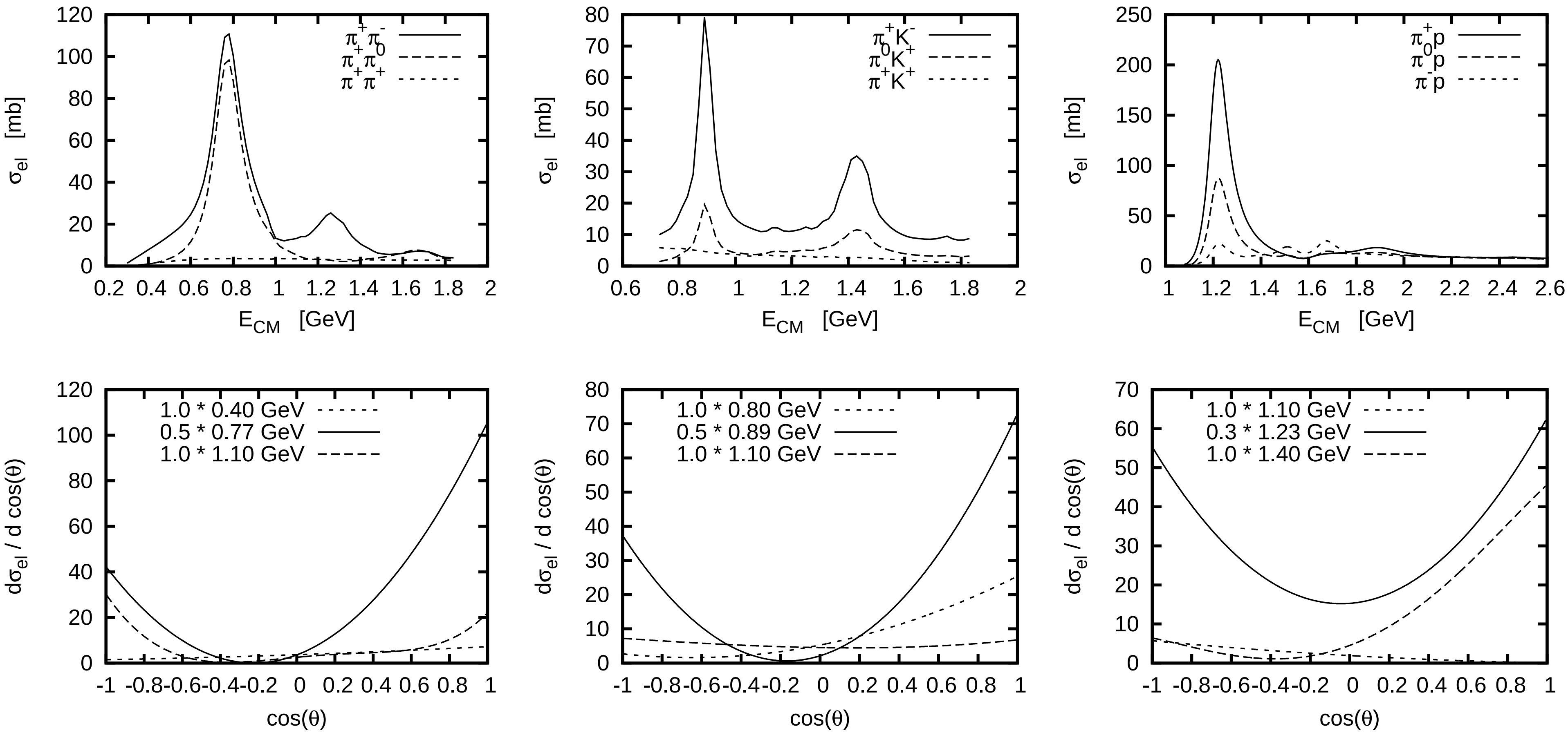
- Order scatterings based on “relative transverse velocity”

$$v_{\perp ij} = \left| \frac{p_{\perp i}}{m_{\perp i}} - \frac{p_{\perp j}}{m_{\perp j}} \right|$$

- Dominated by pions, so focus on $\pi\pi$, $\pi$K and $\pi$p ($\sigma_{\text{max}}^{\text{el}} \sim 200$ mb)
- Most scatterings close to threshold

# Elastic scattering cross sections

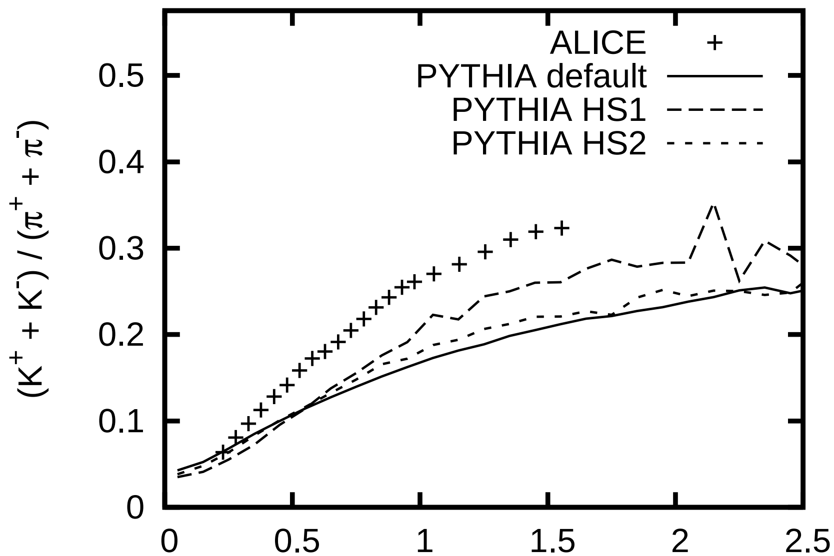
Isospin partial-wave parameterisations (no Columb corrections)

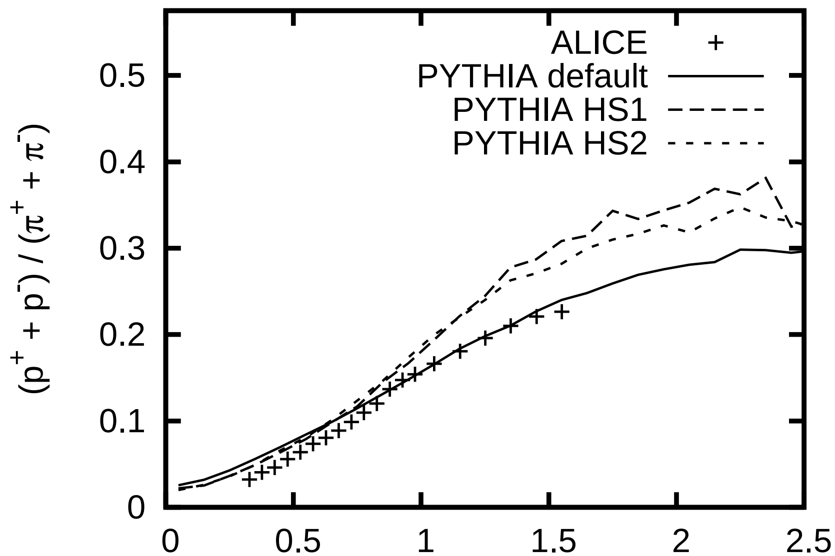


# Preliminary results

Results from simple model (7 TeV)

- Perform scattering after first round of hadron decays ($\eta \to \pi$)
- HS1 - $j = 0.1$, $R_{\text{max}} = 1.0$, flat cross section + isotropic scatterings
- HS2 - $j = 1.0$, $R_{\text{max}} = 1.0$, parameterised cross sections



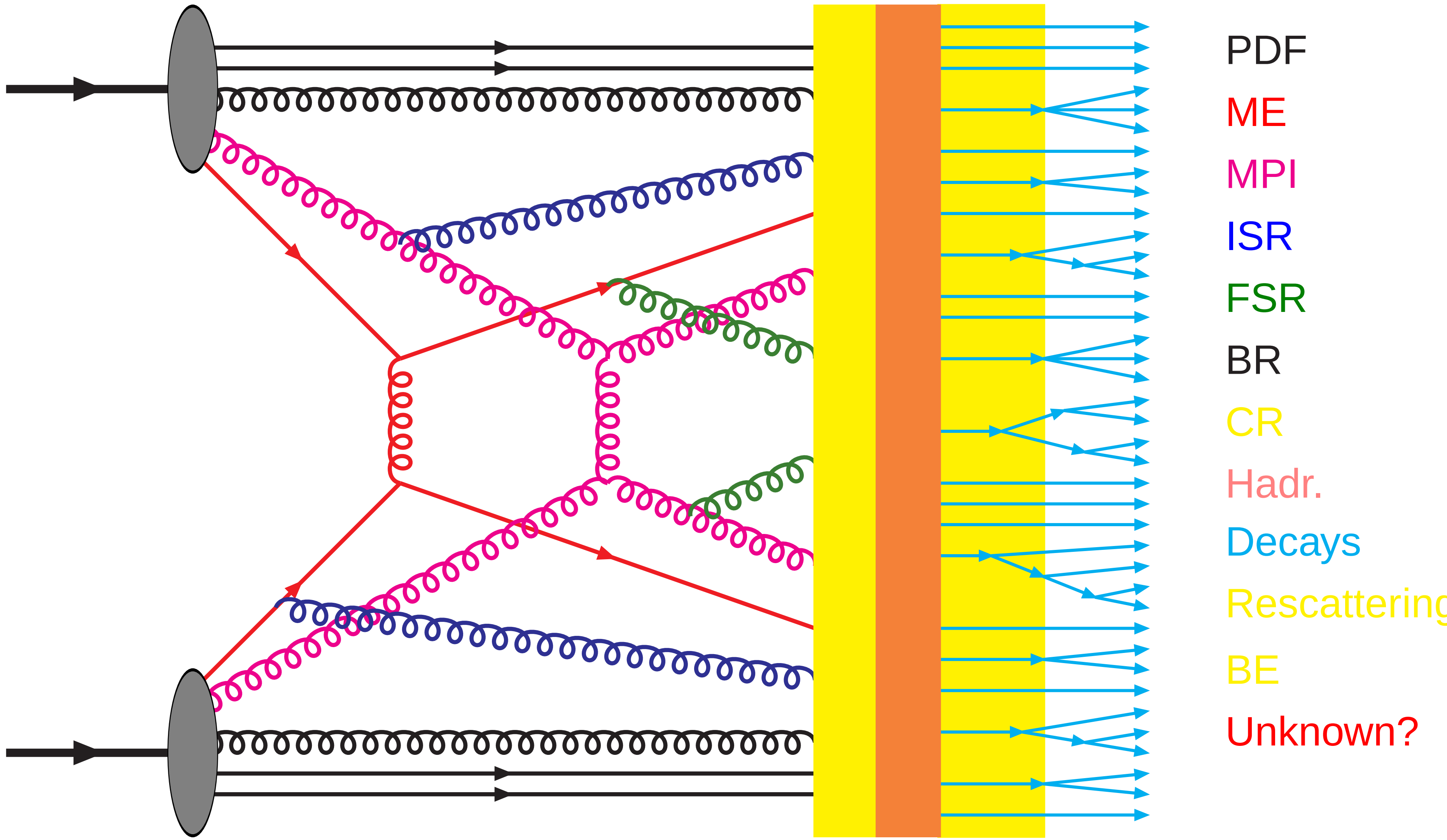




Cross sections are small, limiting effect

- Less multiplicity at 900 GeV $\longrightarrow$ less effect
- Increase in $R_{\text{max}}$ tends to give pairs with higher invariant mass, but cross sections still low in this region

Preliminary results not so promising. . .

. . . but, whatever the correct explanation,
the road leads to more complexity, not less:




Another example: matching to NLO, NNLO

The need for experiment ↔ generators ↔ theory remains.

# Summary and Outlook

- PYTHIA 6 is winding down:
  - ⋆ is supported but not developed;
  - ⋆ still main option for current run (sigh),
  - ⋆ *but not after long shutdown 2013*!
- PYTHIA 8 is the natural successor,
  - ⋆ is (sadly!) not yet quite up to speed in *all* respects,
  - ⋆ but for much already better than PYTHIA6 ,
  - ⋆ is starting to have competitive tunes,
  - ⋆ and will continue to move ahead.
  - ⋆ Currently version 8.150 with the features mentioned here.
- Advise to experimentalists:
  - ⋆ step up PYTHIA 8 usage to gain experience;
  - ⋆ if you want new features then be prepared to use PYTHIA 8;
  - ⋆ provide feedback, both what works and what does not;
  - ⋆ do your own tunes to data and tell outcome.

News list: `http://www.hepforge.org/lists/listinfo/pythia8-announce`

**The work is never done!**